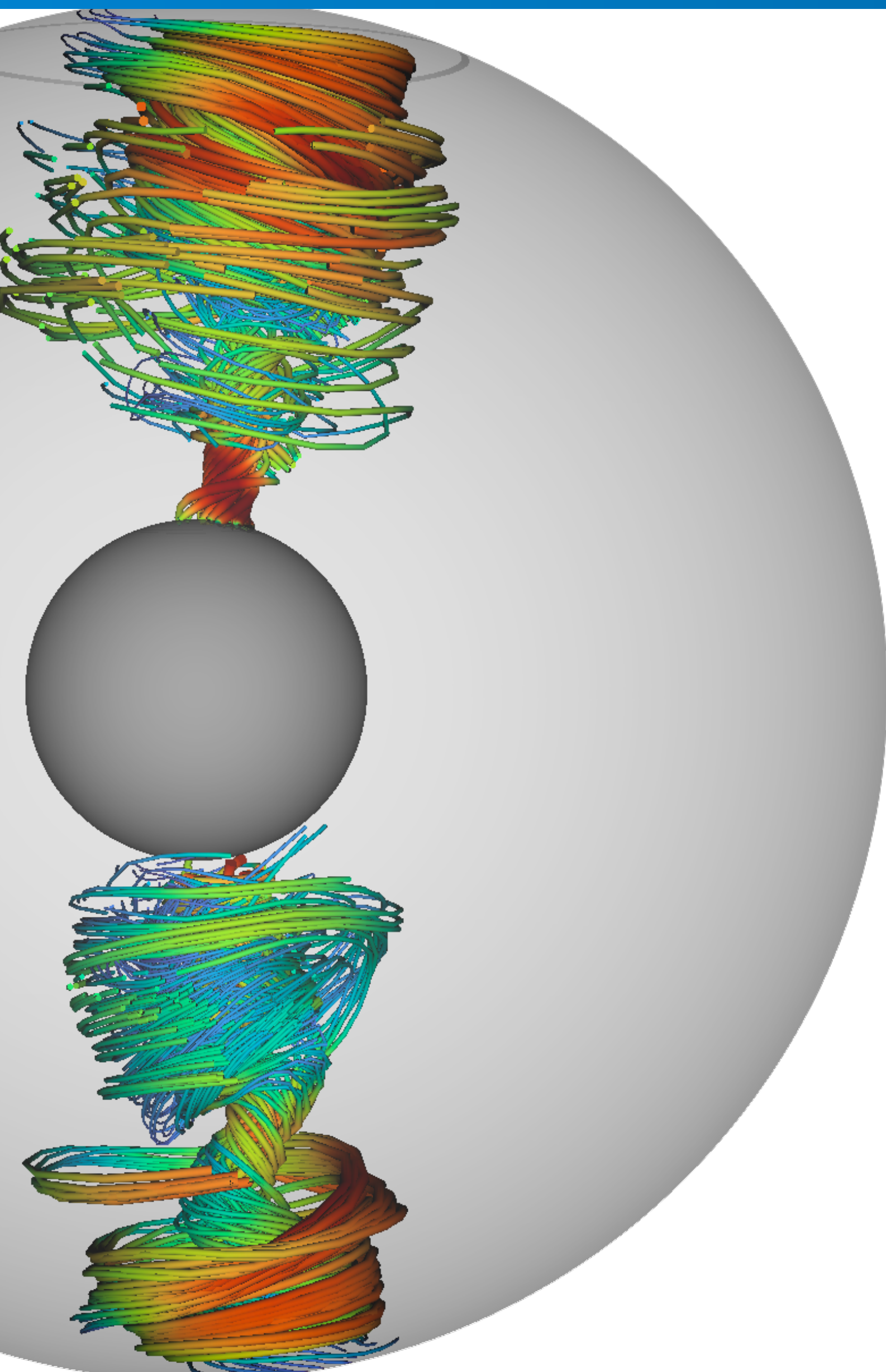




New branch of the Tayler instability-driven dynamo in 3D numerical simulations of stellar radiative zones

Paul Barrère
(Département d'Astronomie,
Université de Genève)

Collaborators on this work:

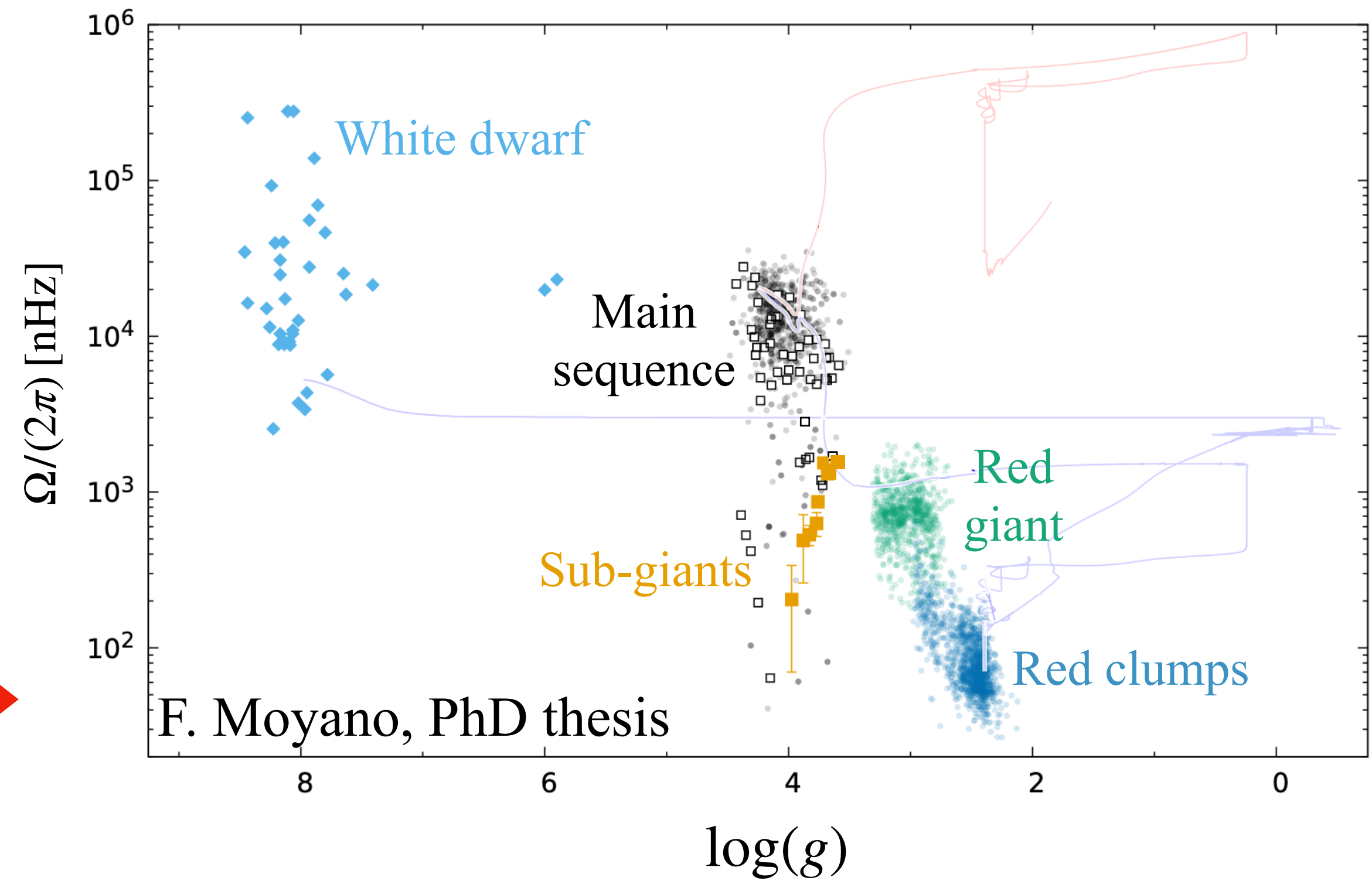
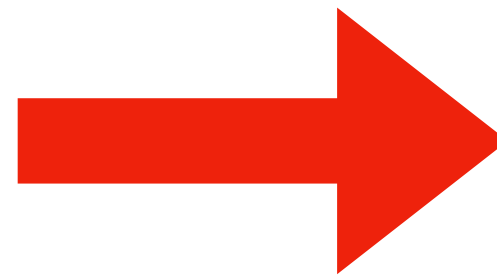
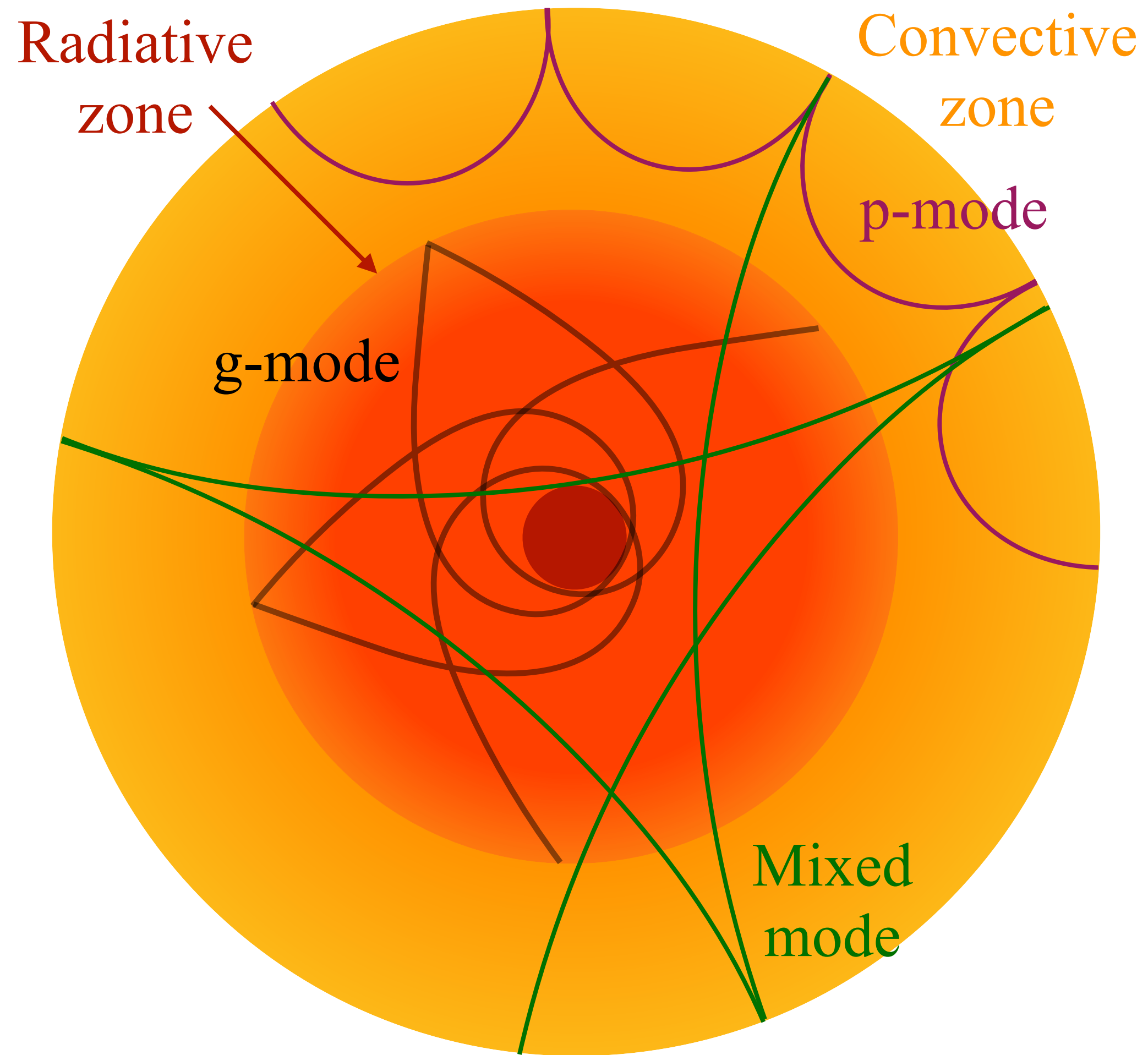
A. Reboul-Salze (AEI, Potsdam),
P. Eggenberger (Univ. Genève)
S. Deheuvels (IRAP, Toulouse)

17th December 2025

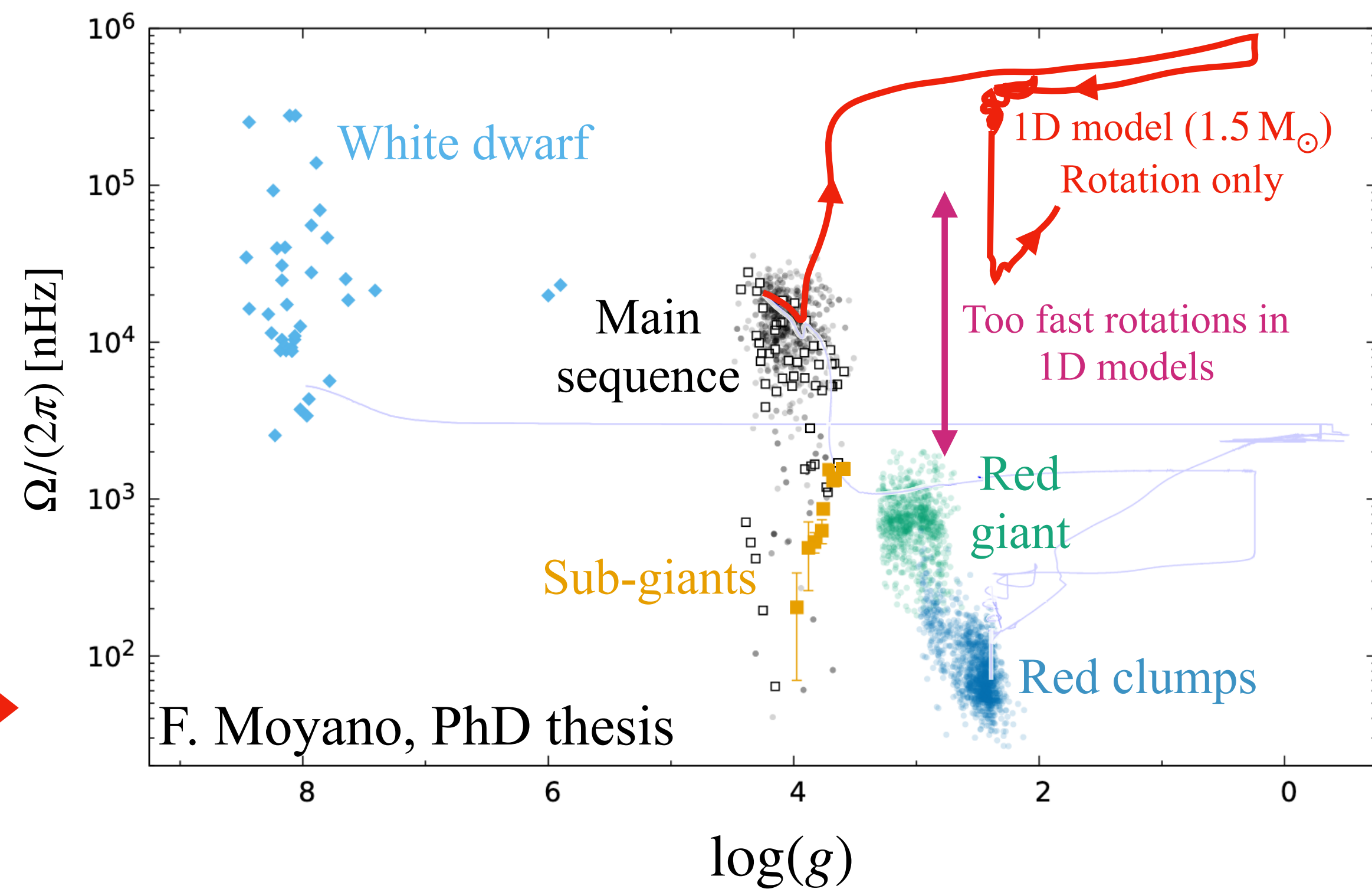
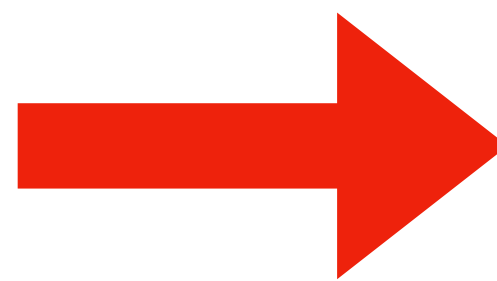
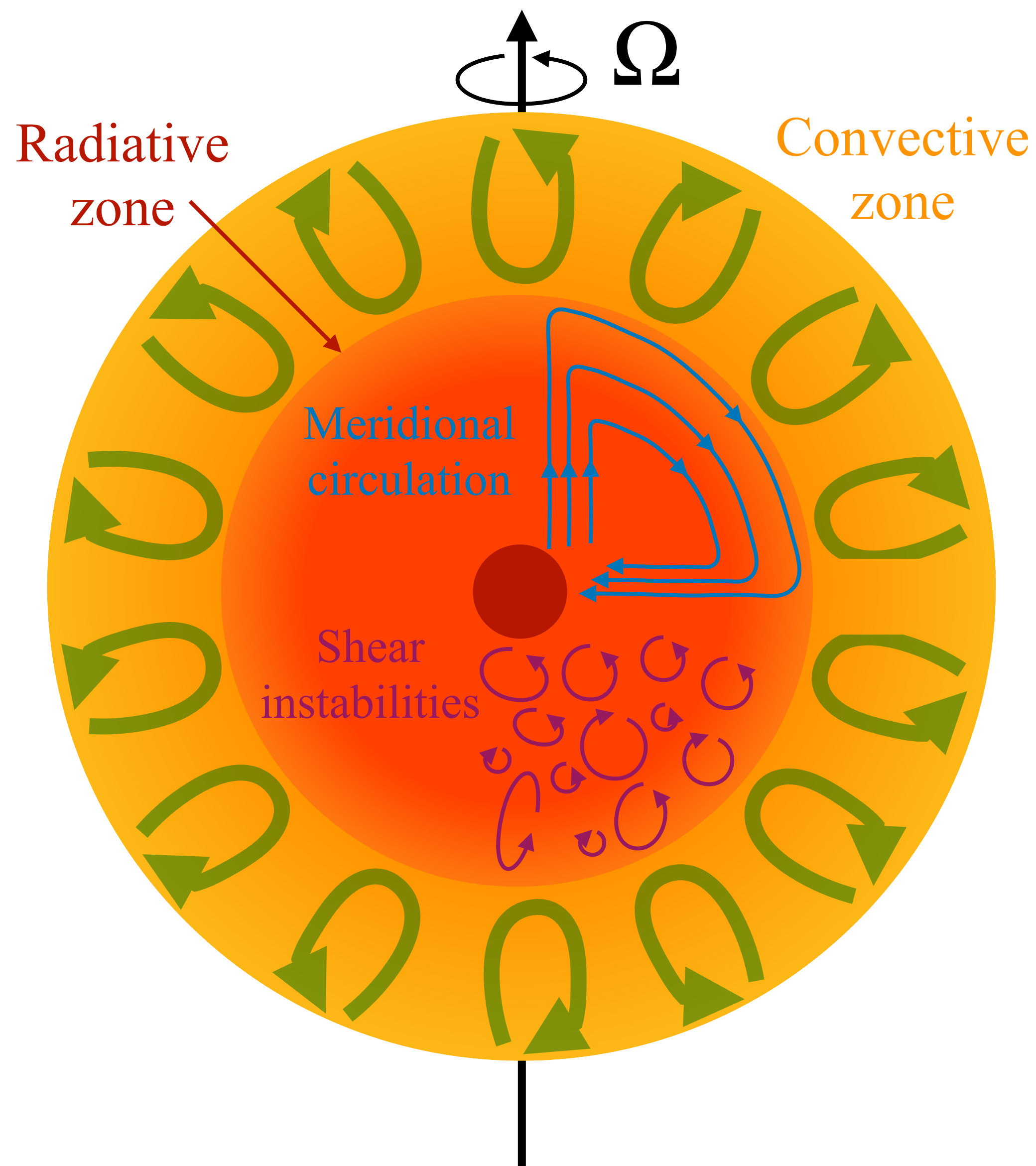


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DE GENÈVE

Rotation of stellar interior through asteroseismology

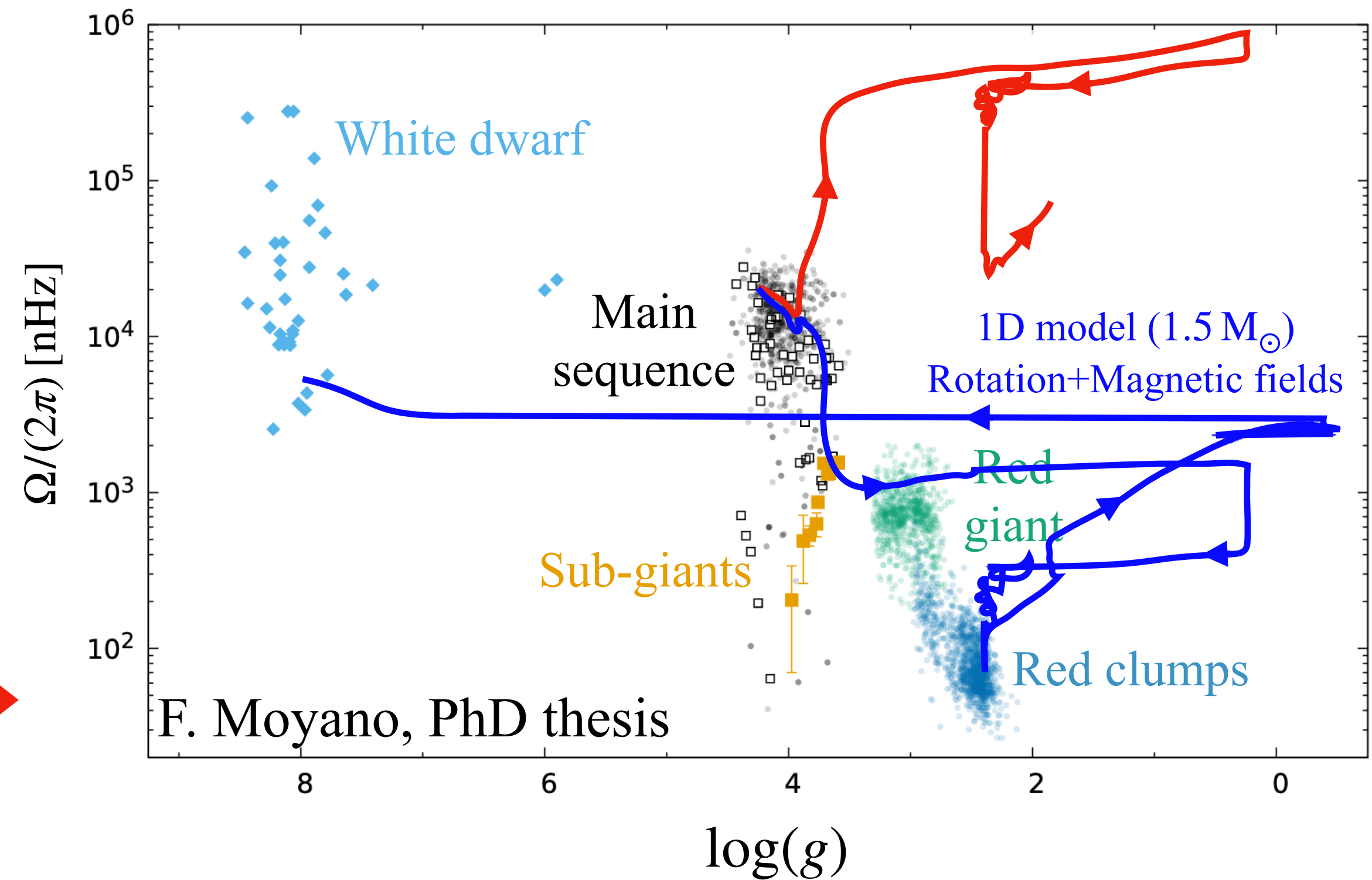
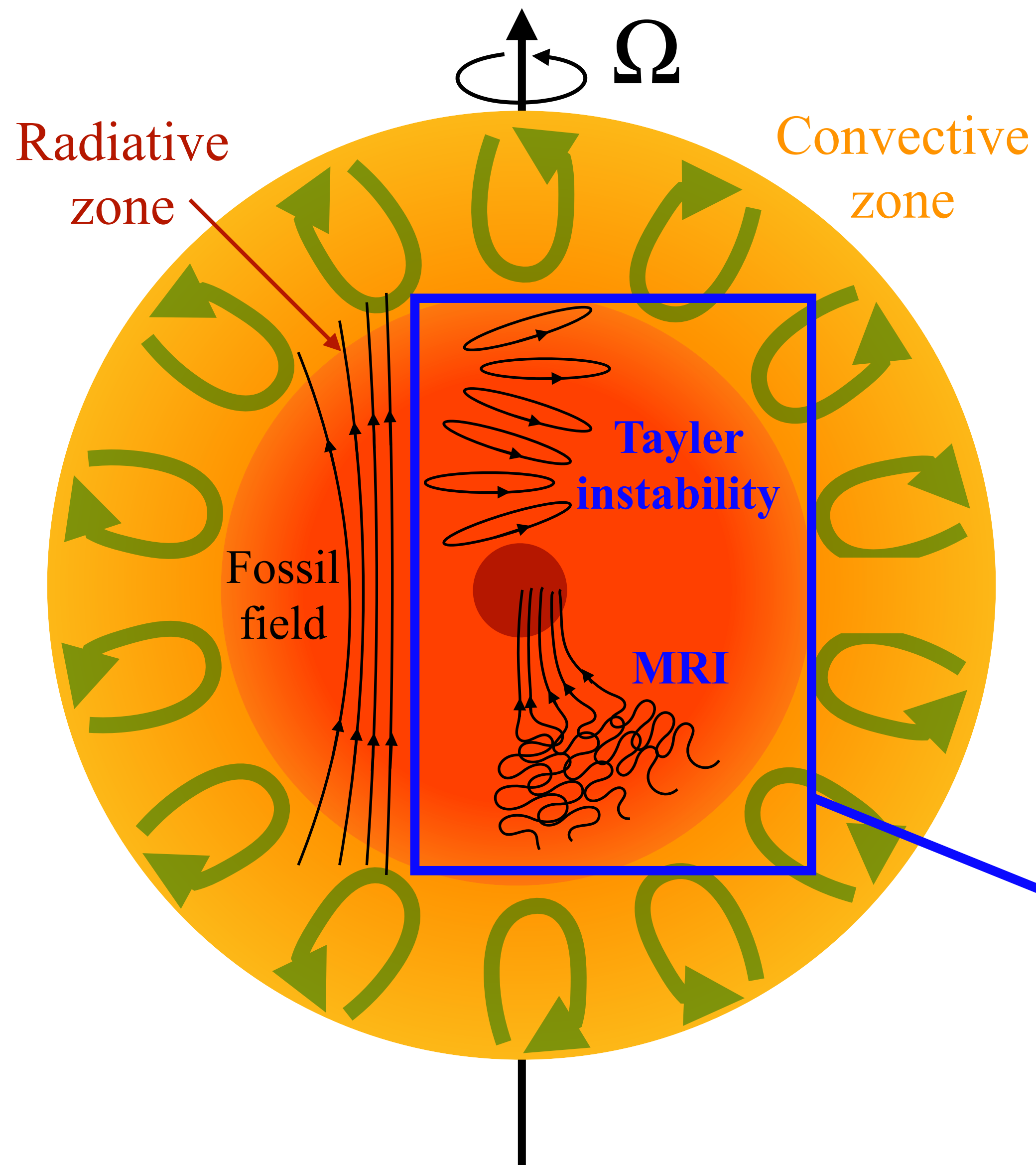


Discrepancy with 1D rotating stellar evolution models



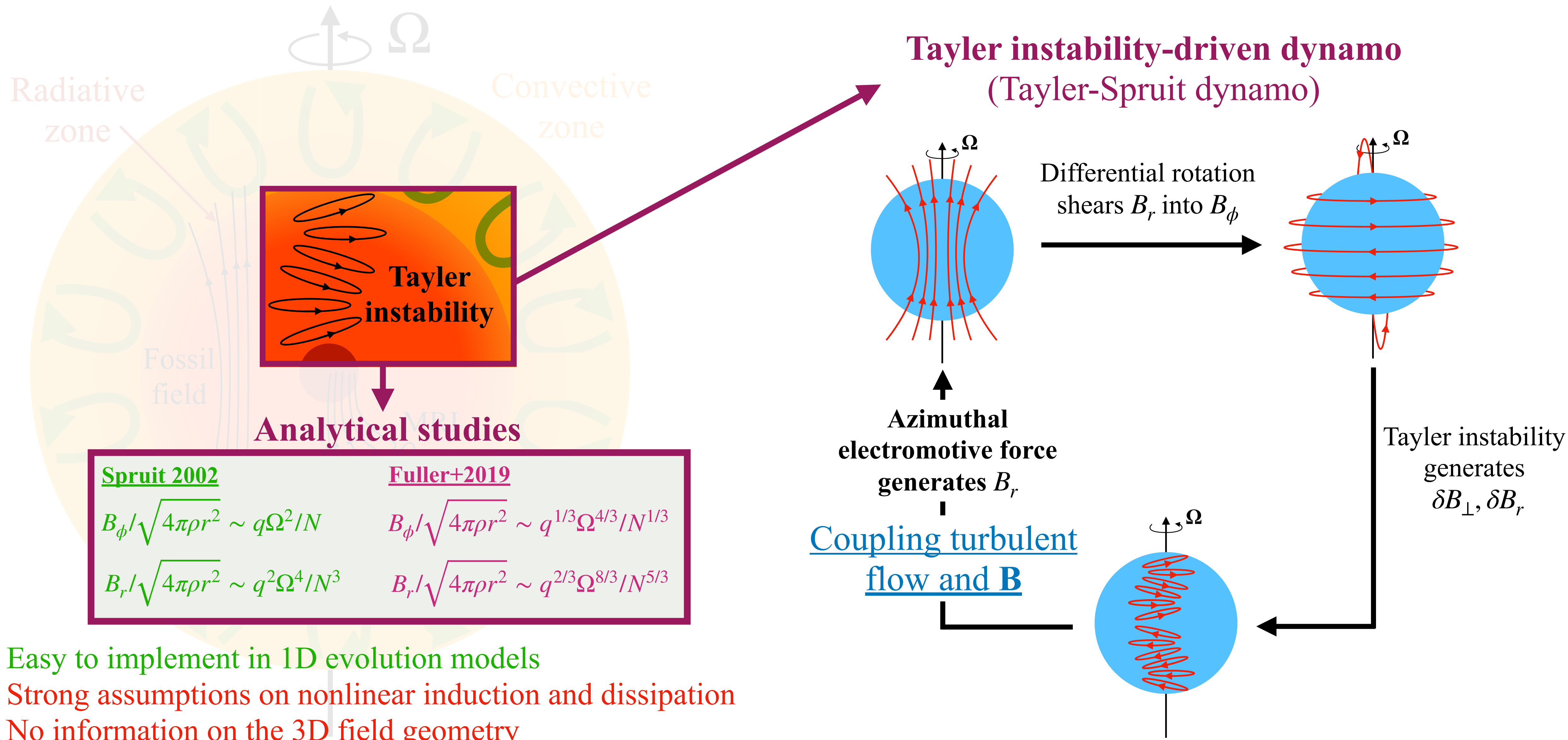
⇒ Other efficient AM transport mechanisms needed

Transport by magnetic fields



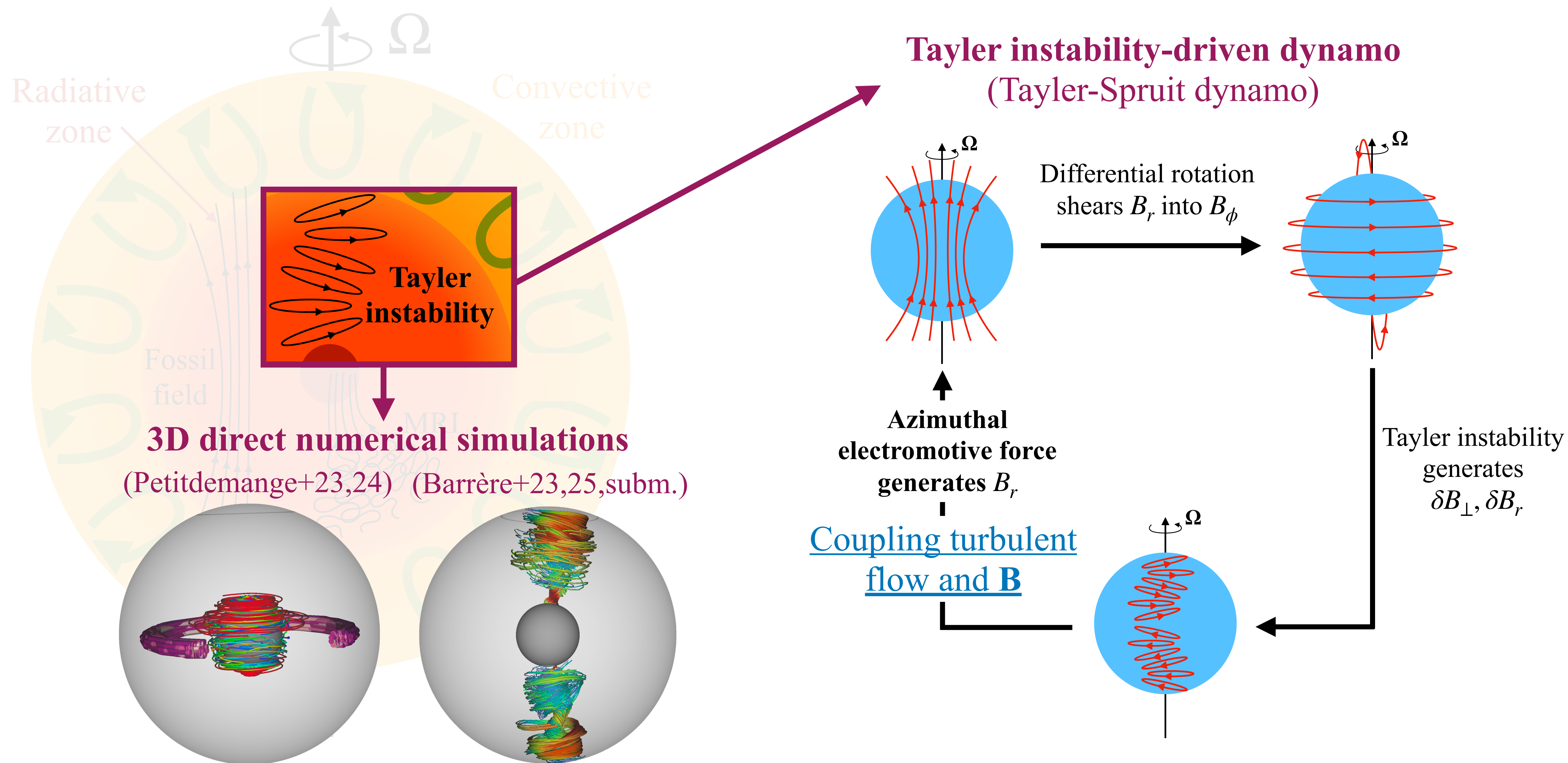
Large-scale magnetic fields produced by
subcritical dynamos
driven by MHD instabilities

Taylor instability-driven dynamo



- Easy to implement in 1D evolution models
- Strong assumptions on nonlinear induction and dissipation
- No information on the 3D field geometry

Necessity of 3D models



Open source code: MagIC

- **Pseudo-spectral code**
 r : Chebyshev polynomial
 θ and ϕ : Sphericals harmonics
- **Parallelisation**
MPI
OpenMP
- **Open source !**
magic-sph.github.io

Home Get it/Run it Contribute! Numerical methods Contents

MagIC

Table of Contents

- Welcome
 - Quickly starting using MagIC
 - Documentation
 - Contributing to the code
 - Giving credit

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Welcome

Magneto-rotational instability in spherical shell
Jouve, L. et al., A&A, 2015

Setup

$\vec{\Omega}$

r_o , r_i , r , θ , ϕ , d

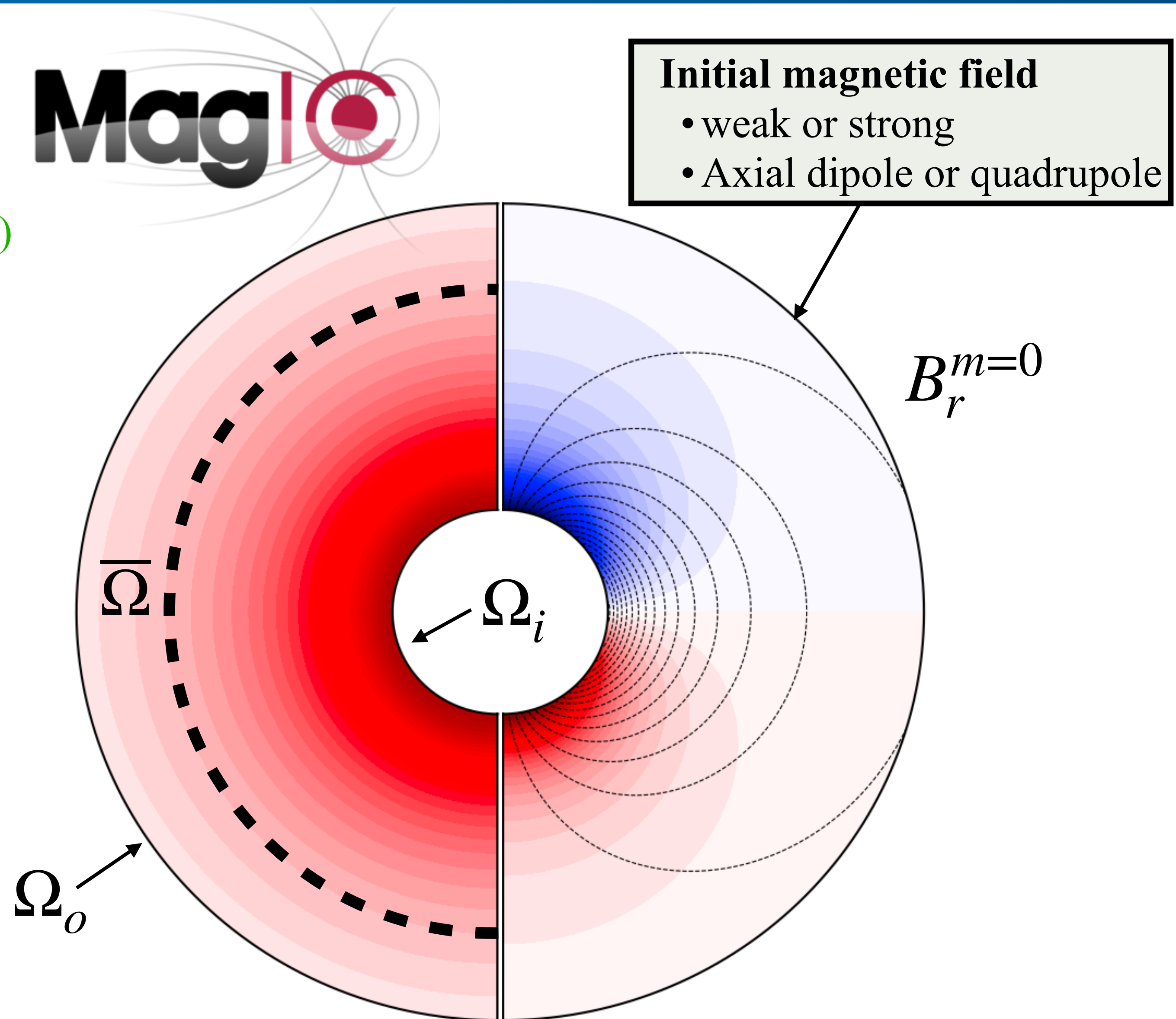
MagIC is a numerical code that can simulate fluid dynamics in a spherical shell. MagIC solves for the Navier-Stokes equation including Coriolis force, optionally coupled with an induction equation for Magneto-Hydro Dynamics (MHD), a temperature (or entropy) equation and an equation for chemical composition under both the anelastic and the Boussinesq approximations.

(Wicht 2002, Gastine & Wicht 2012, Schaeffer 2013)

Numerical setup

Boussinesq approximation:

- **Uniform** density profile
 - Density perturbations **ONLY** in buoyancy term
- Less costly in CPU hours (**~100 kCPU hours per DNS**)
⇒ Easier to produce parametric studies
- Unrealistic in stars



Numerical setup

Boussinesq approximation:

- **Uniform** density profile
- Density perturbations **ONLY** in buoyancy term

- **Explicit diffusivities**

$$Pr = \nu/\kappa = 0.1 \gg 10^{-5}$$

$$Pm = \nu/\eta = 4$$

- **Stably stratified**

$$N/\bar{\Omega} \in [2, 130]$$

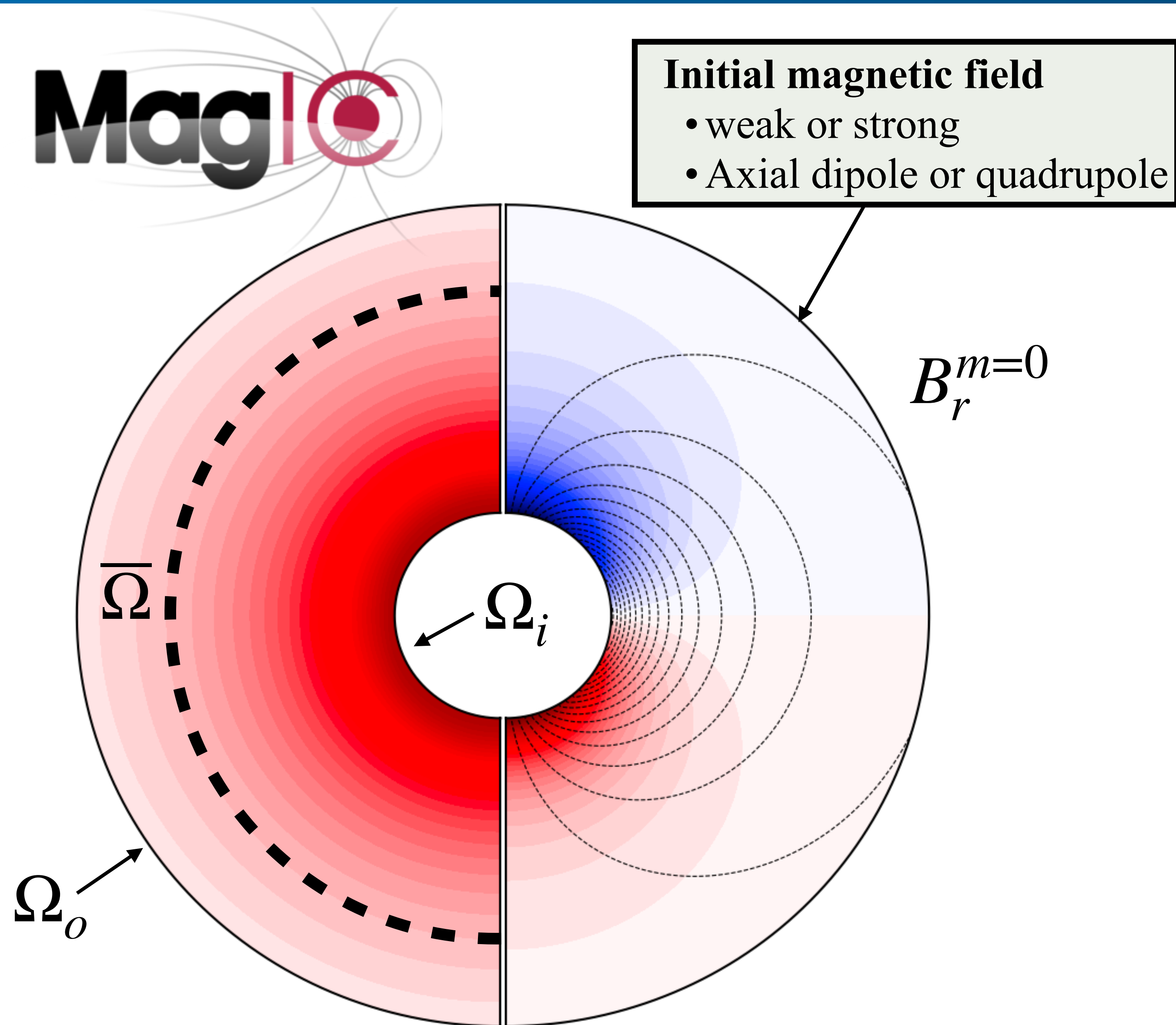
- **Volumetric forcing of differential rotation**

$$\text{Negative shear } \Omega_i > \Omega_o$$

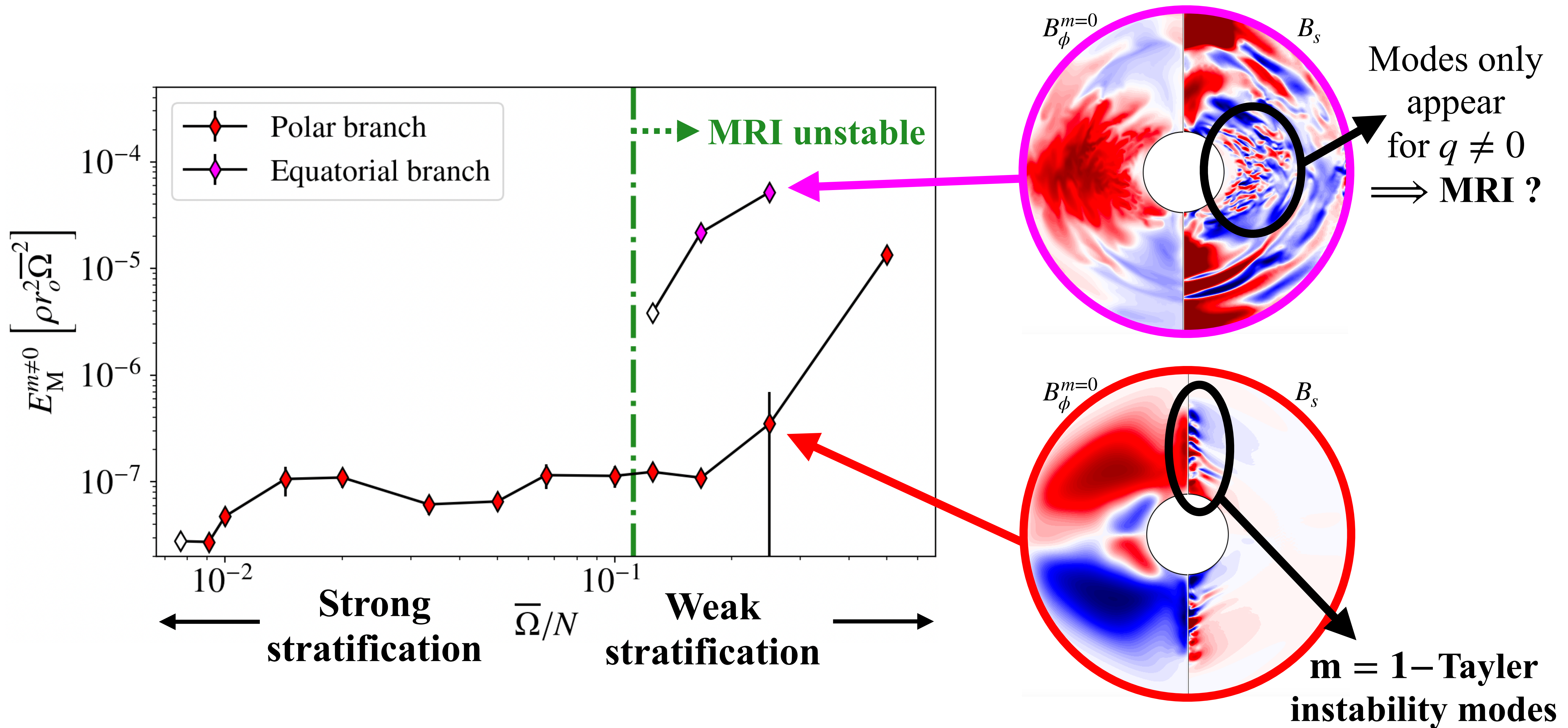
$$\text{Shear rate } q \equiv \frac{r}{\Omega} \frac{\partial \Omega}{\partial r}$$

- **Medium resolution: (257, 256, 512)**

- **High resolution: (384, 384, 768)**

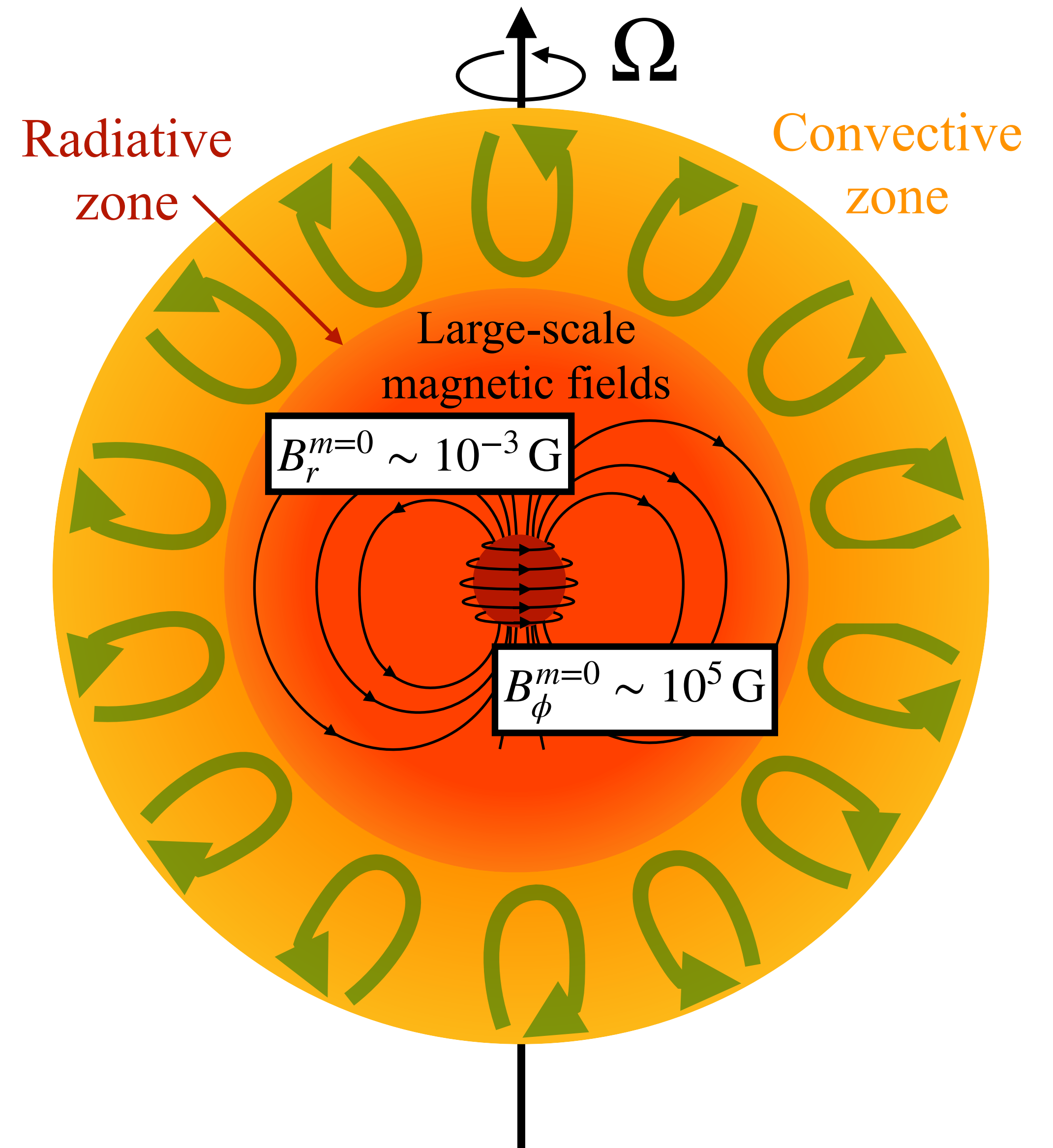
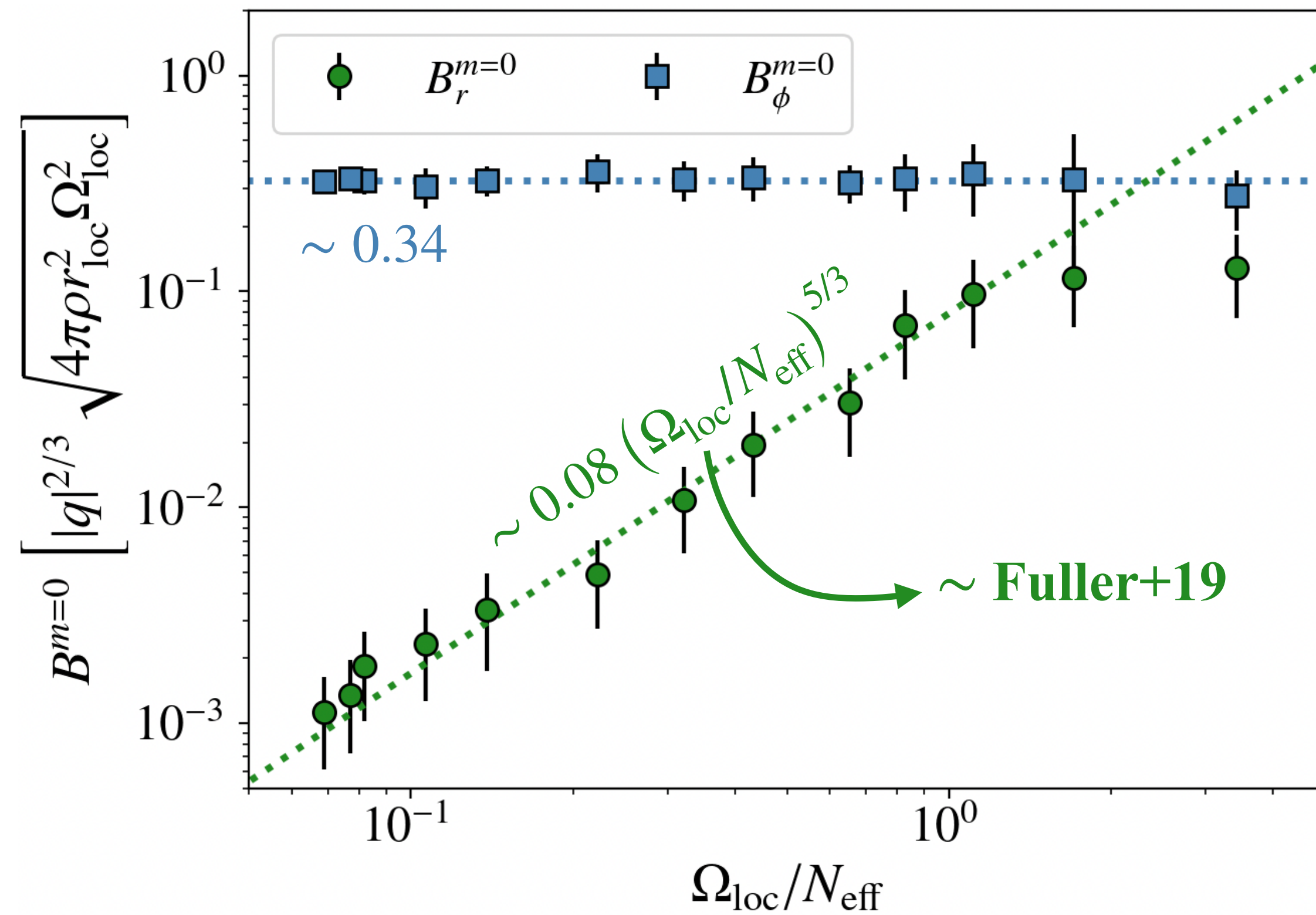


Different regimes of the dynamo



Magnetic field intensity in red giants

Large-scale magnetic fields

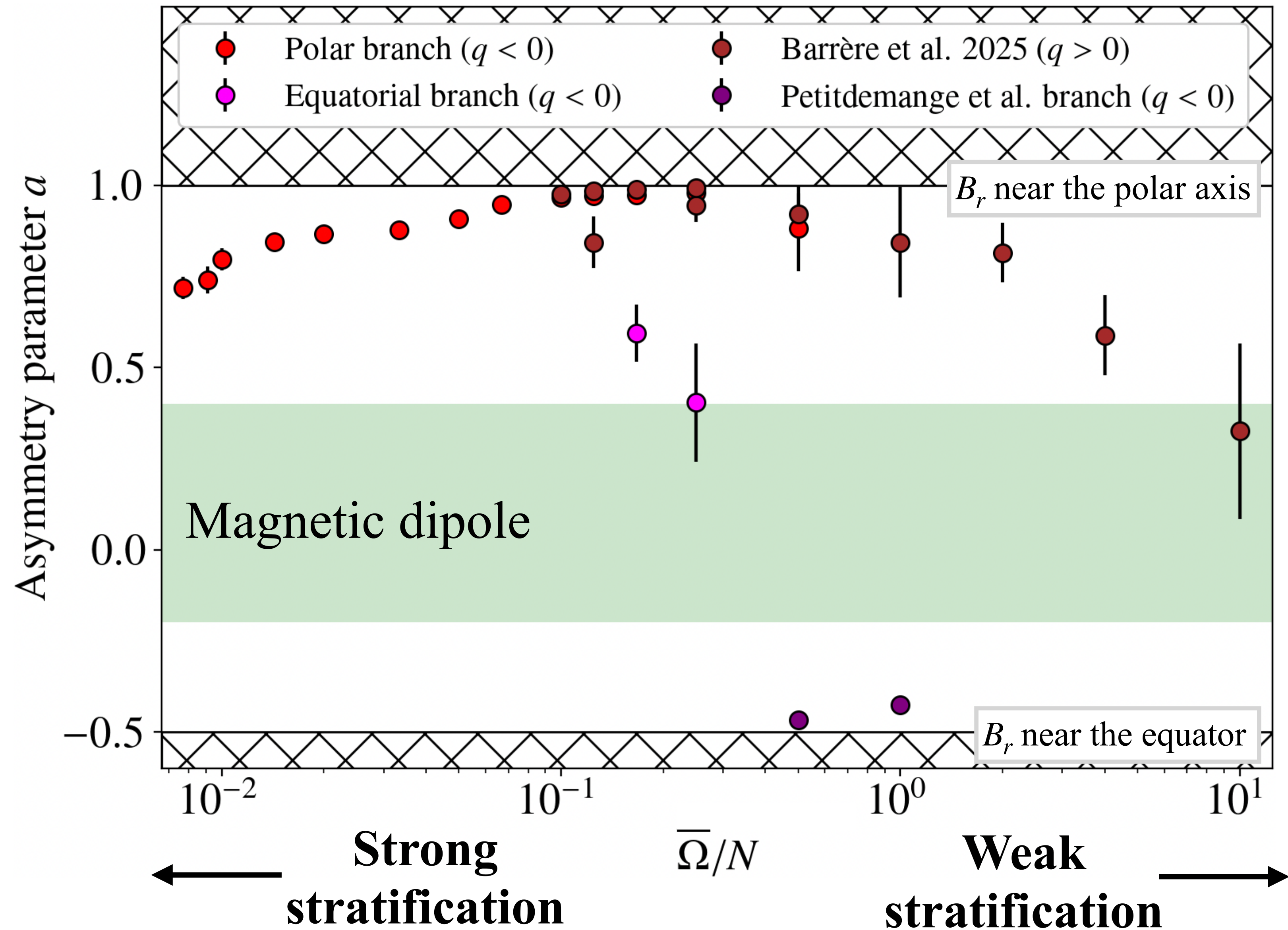


Radial magnetic field topology

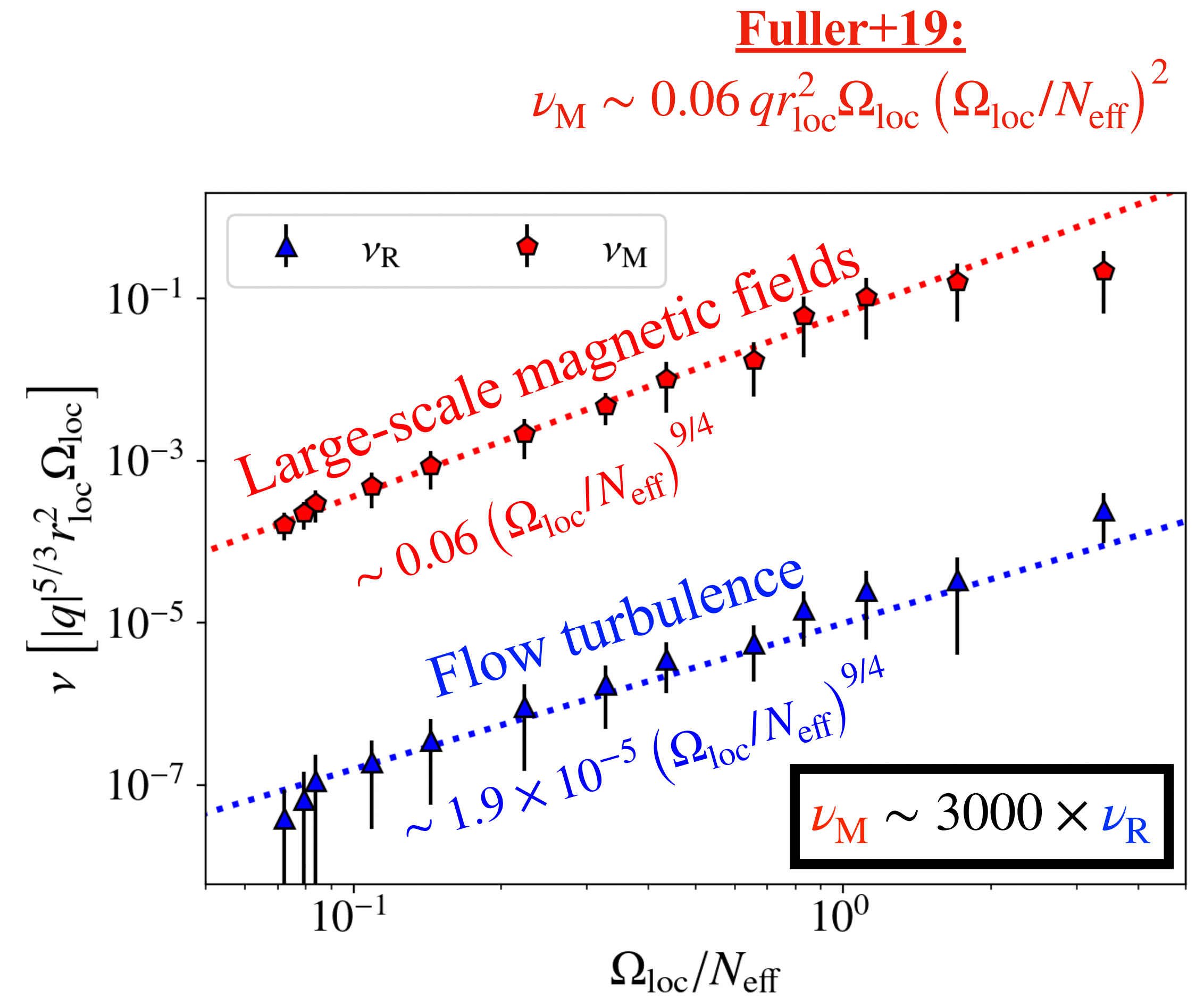
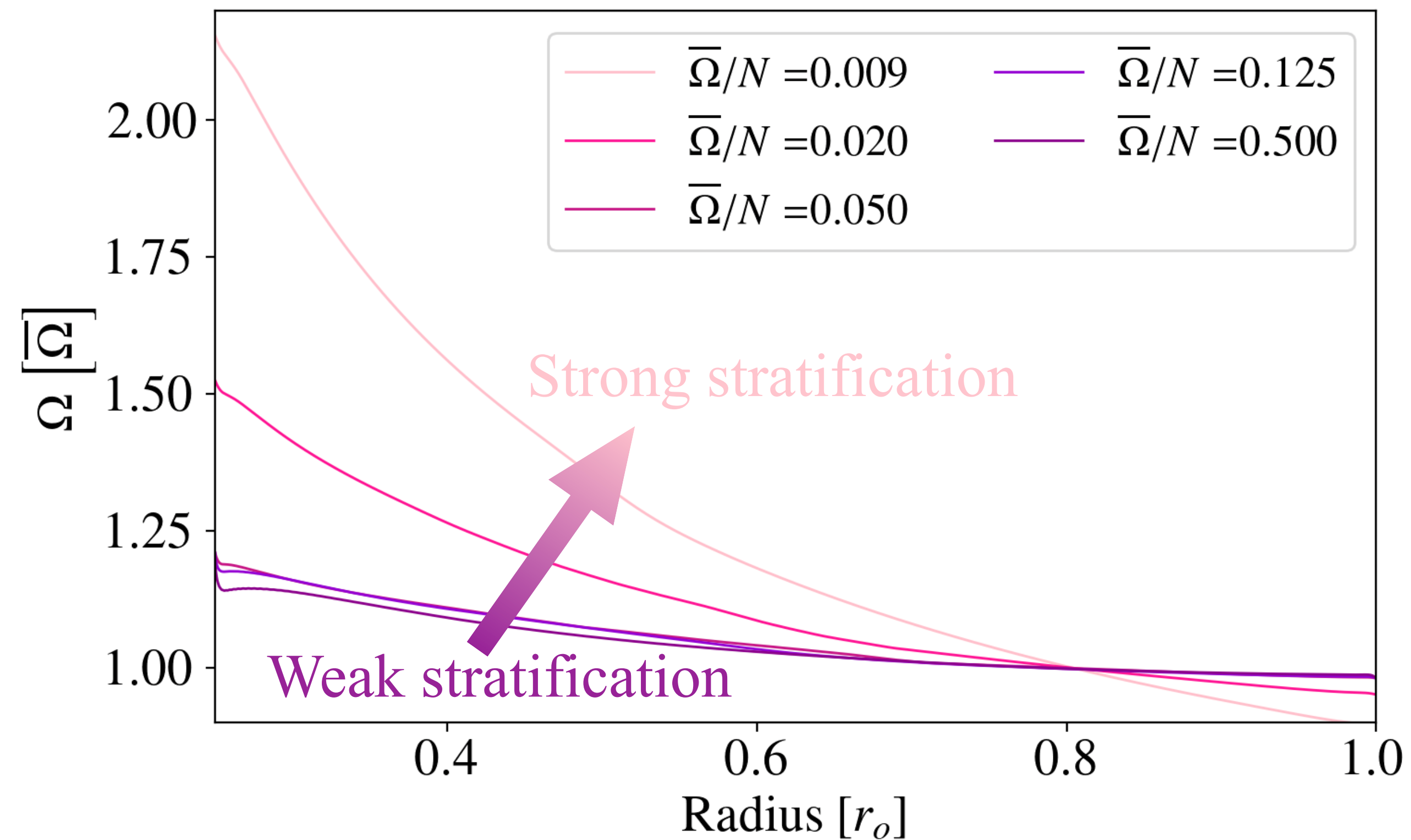
Asymmetry parameter

$$a \equiv \frac{\int_{r_i}^{r_o} K(r) \int \int_S B_r^2 P_2(\cos \theta) d\Omega dr}{\int_{r_i}^{r_o} K(r) \int \int_S B_r^2 d\Omega dr}$$

Weight function
probing g -mode cavity



Angular momentum transport

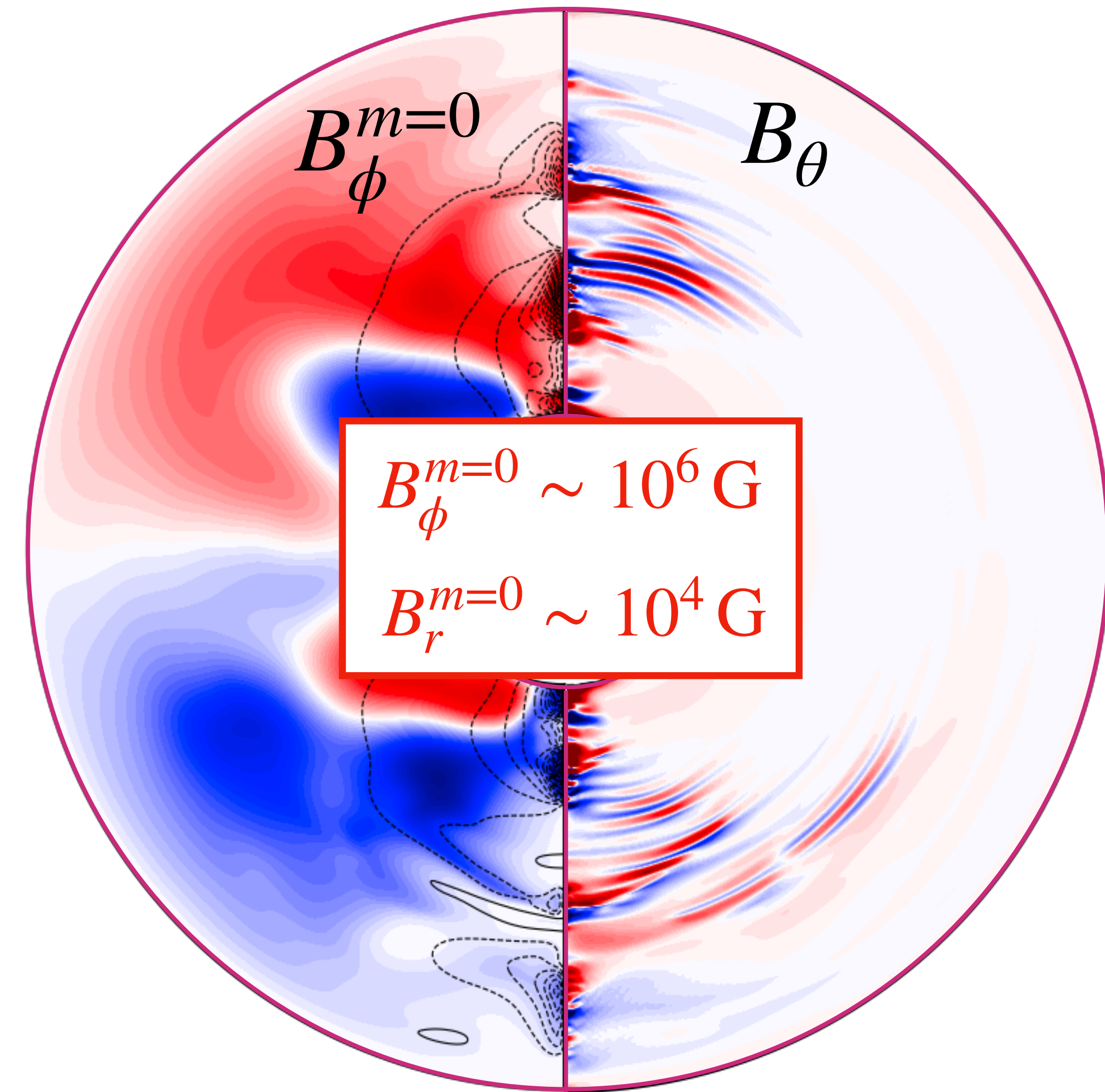
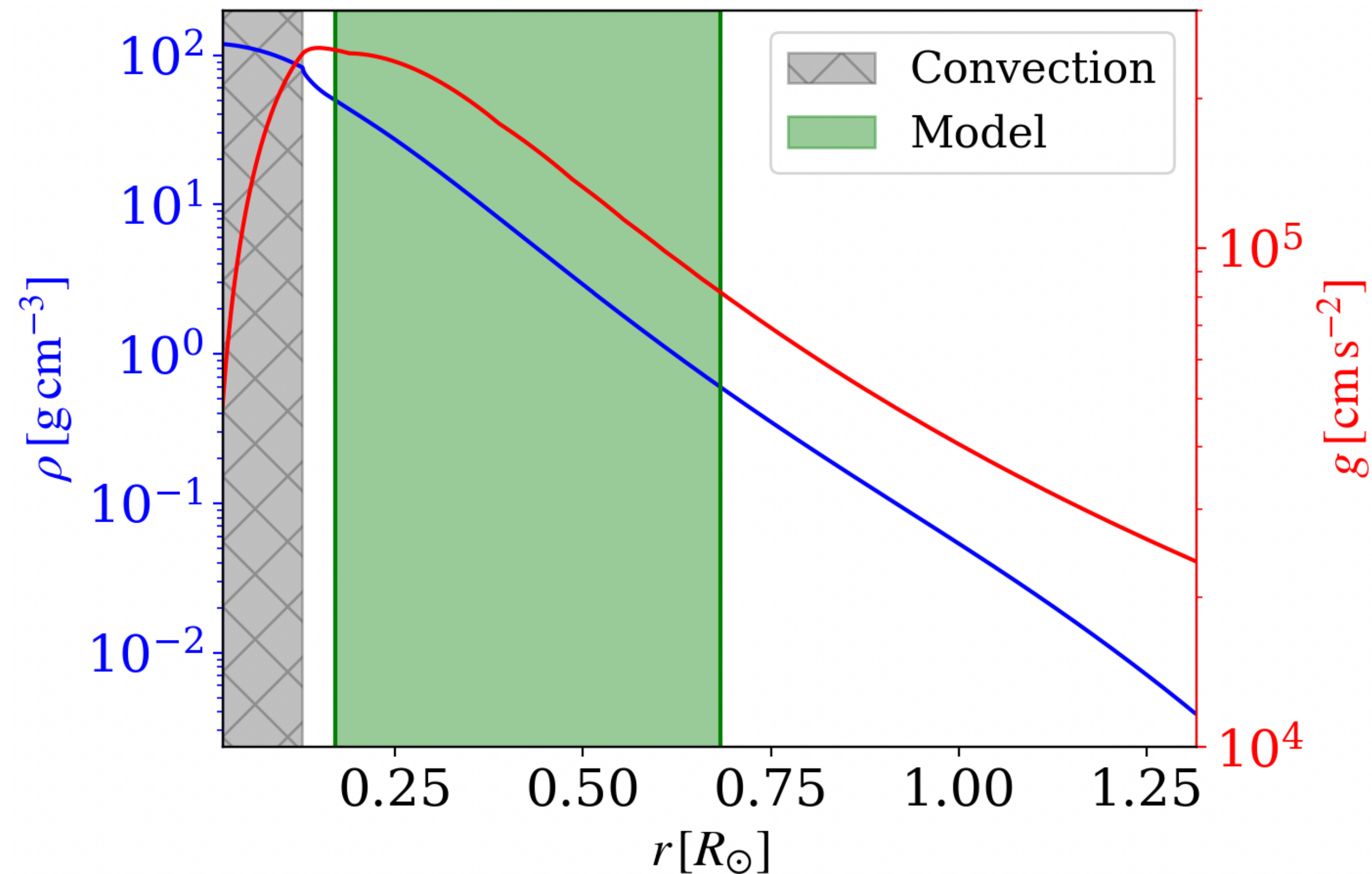


First anelastic simulations

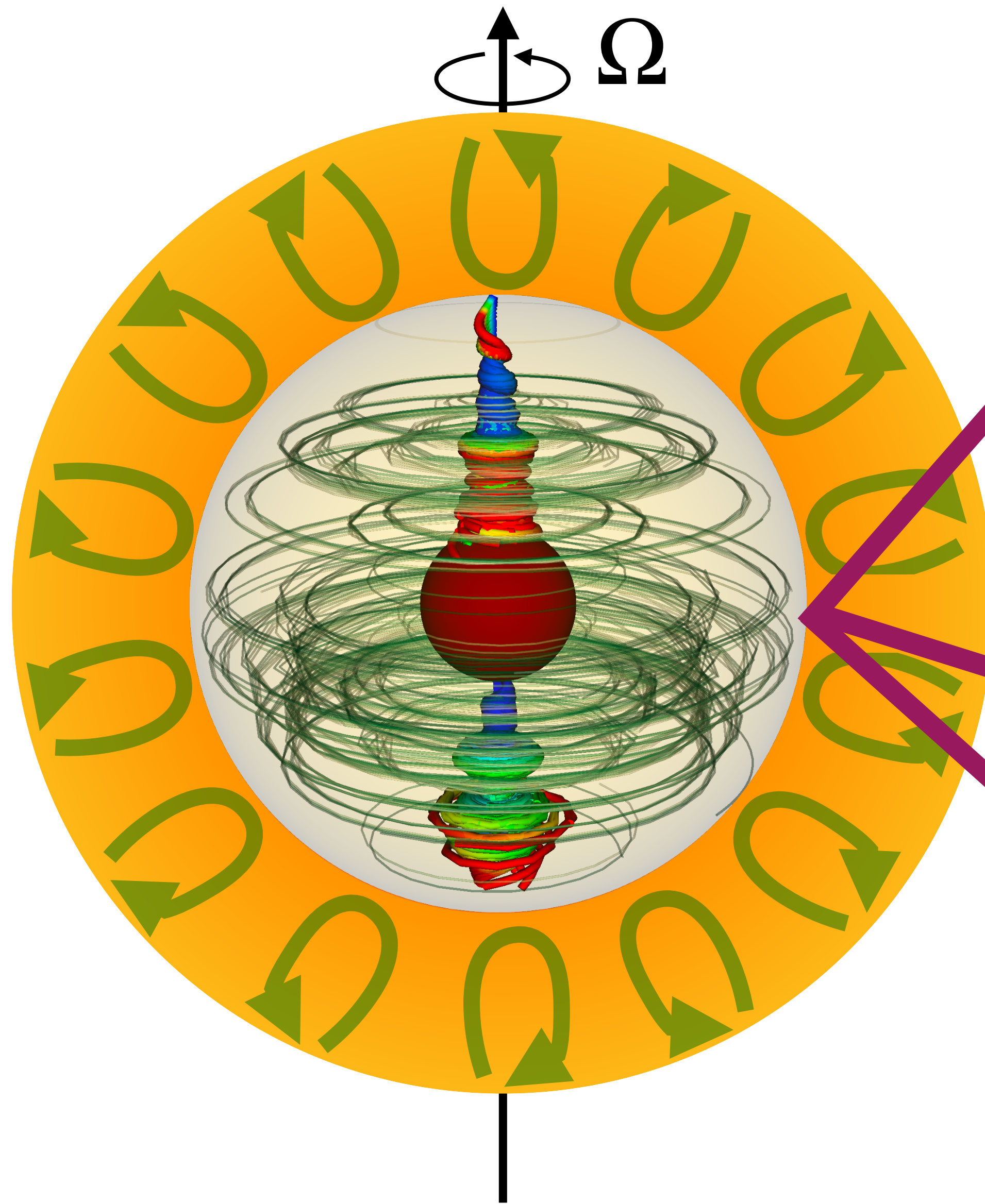
Anelastic approximation:

- **Realistic** density profile
- Flow only produce **perturbations of a reference state**

- Realistic structures
- Numerically costly
- Impact of a density gradient
- Reference state is stationary



Conclusions



New information thanks to 3D models:

- ▶ Tayler instability-driven dynamo active at **large** N/Ω
 $\implies$ Critical $Rm_{\text{turb}} \sim 10$
- ▶ **Toroidal** component ($\ell = 2, m = 0$) dominates
- ▶ Radial field near the **polar axis** ($\ell = 2, m = 0$)

Perspectives

Physics:

- ▶ Anelastic simulations with realistic backgrounds
- ▶ Characterise the nonlinear saturation
- ▶ Study the subcritical transition between both branches

Astrophysics:

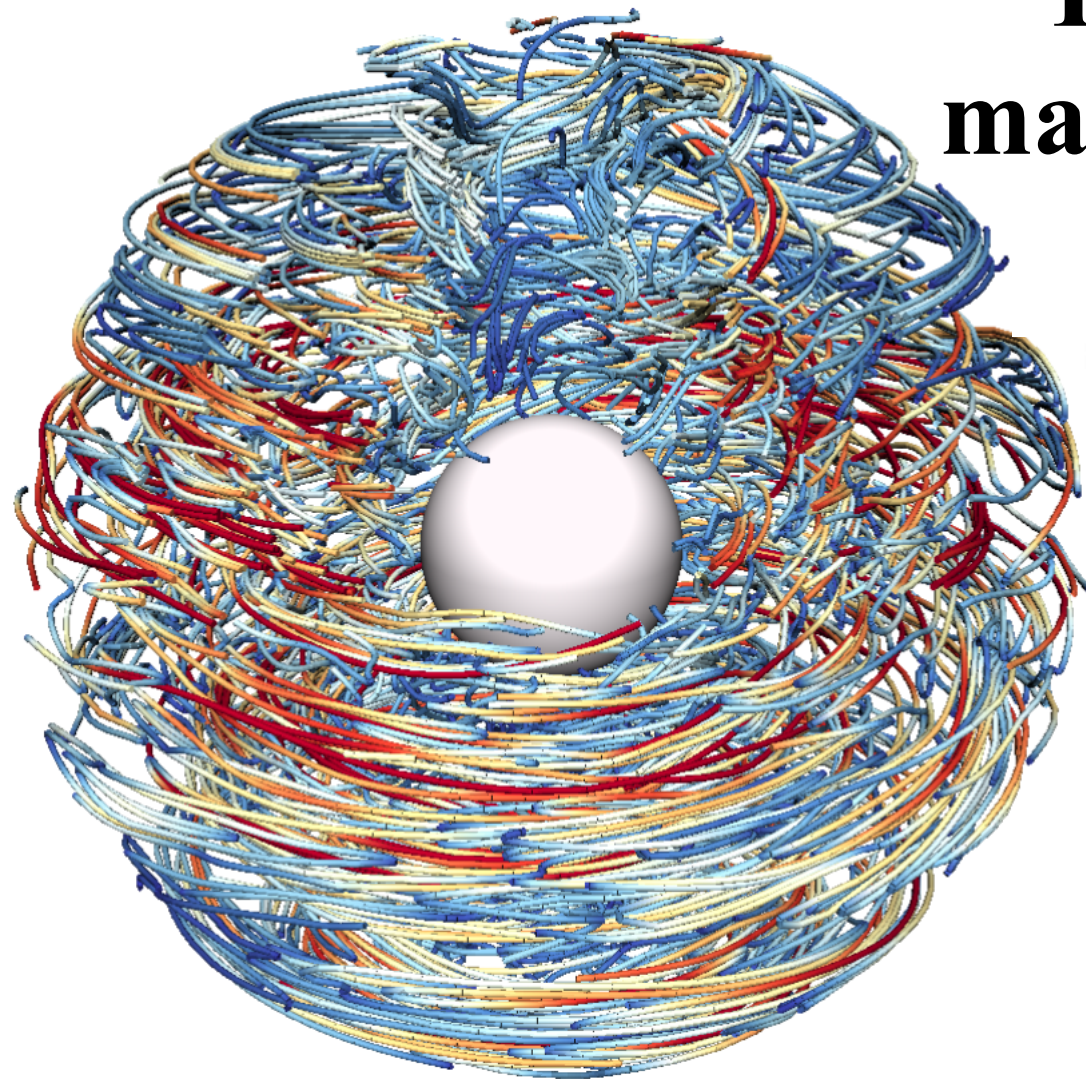
- ▶ Signature in the astero. signal
- ▶ Prescriptions for 1D stellar evolution models

Additional material

Origin of the magnetic field

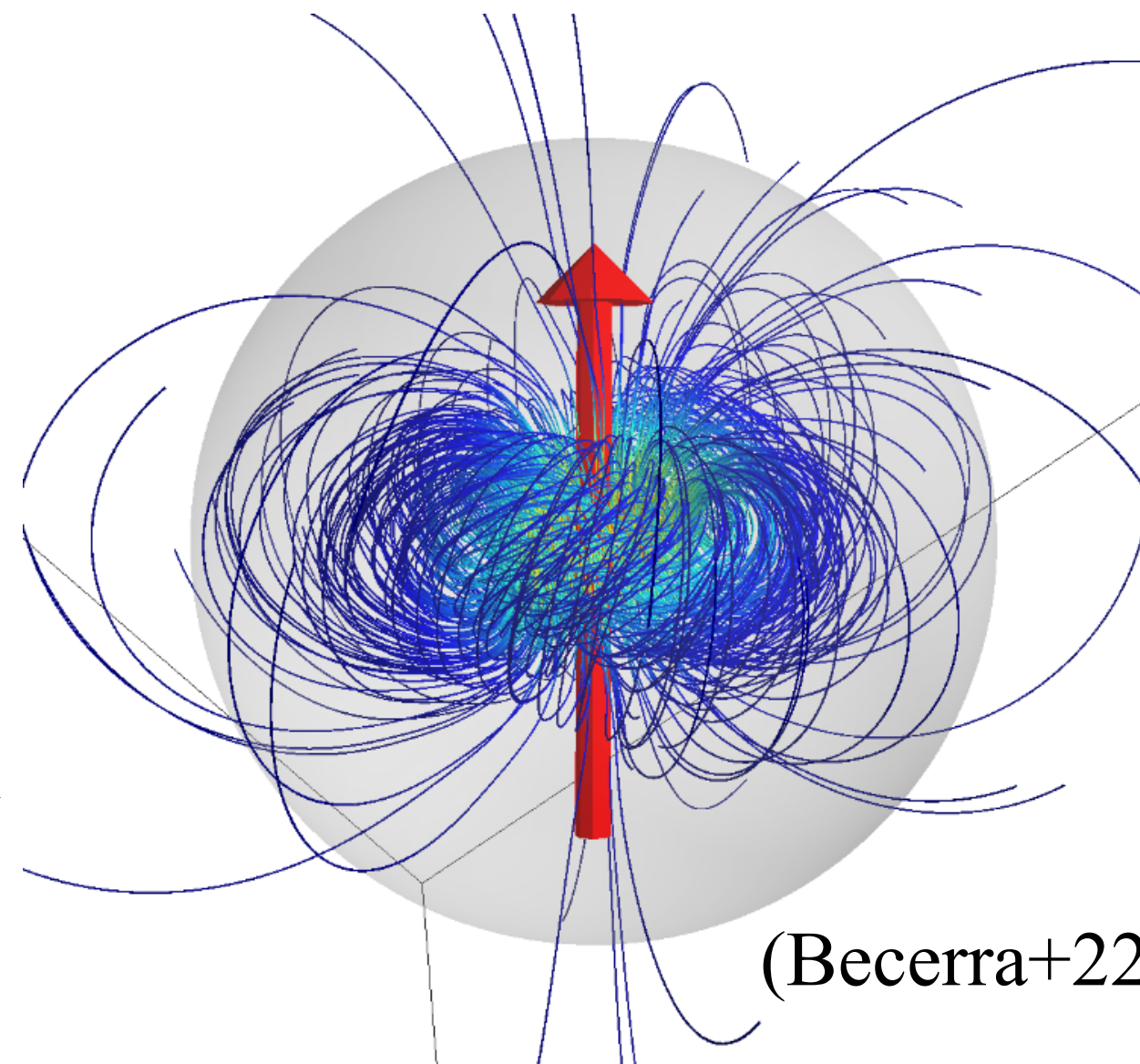
Fossil field origin

**Turbulent
magnetic field**



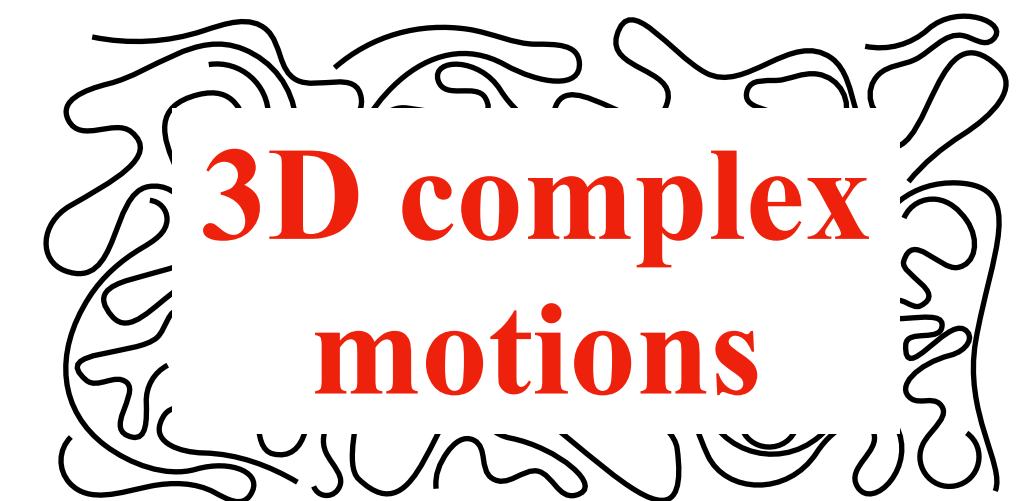
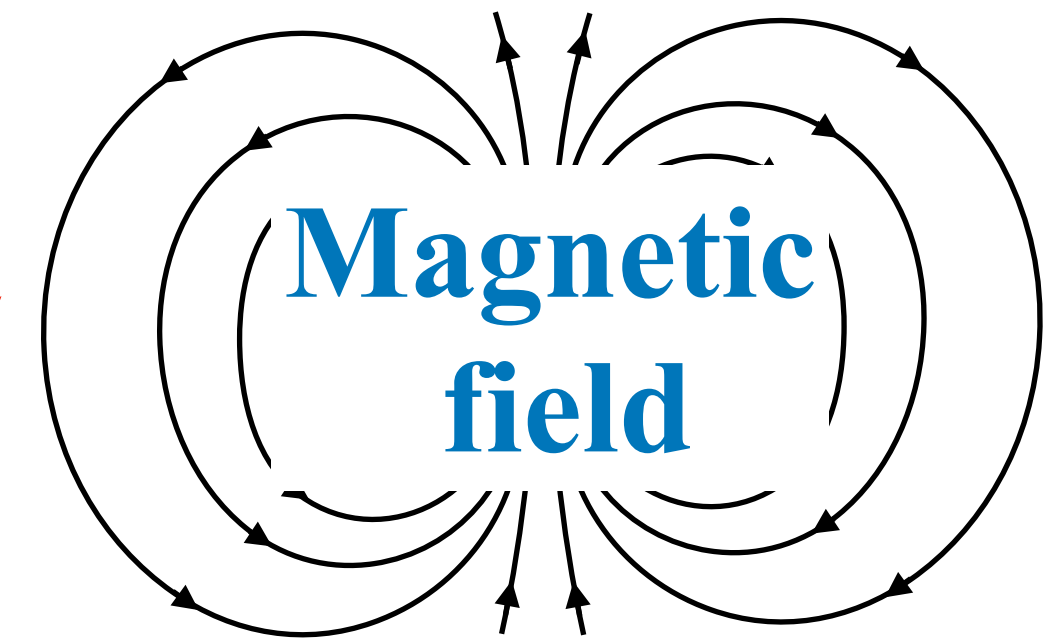
(Reboul-Salze+21)

**Stable
configuration**



(Becerra+22)

Dynamo action



**3D complex
motions**

Diversity of motion sources

=

Different dynamo mechanisms

Magnetic field

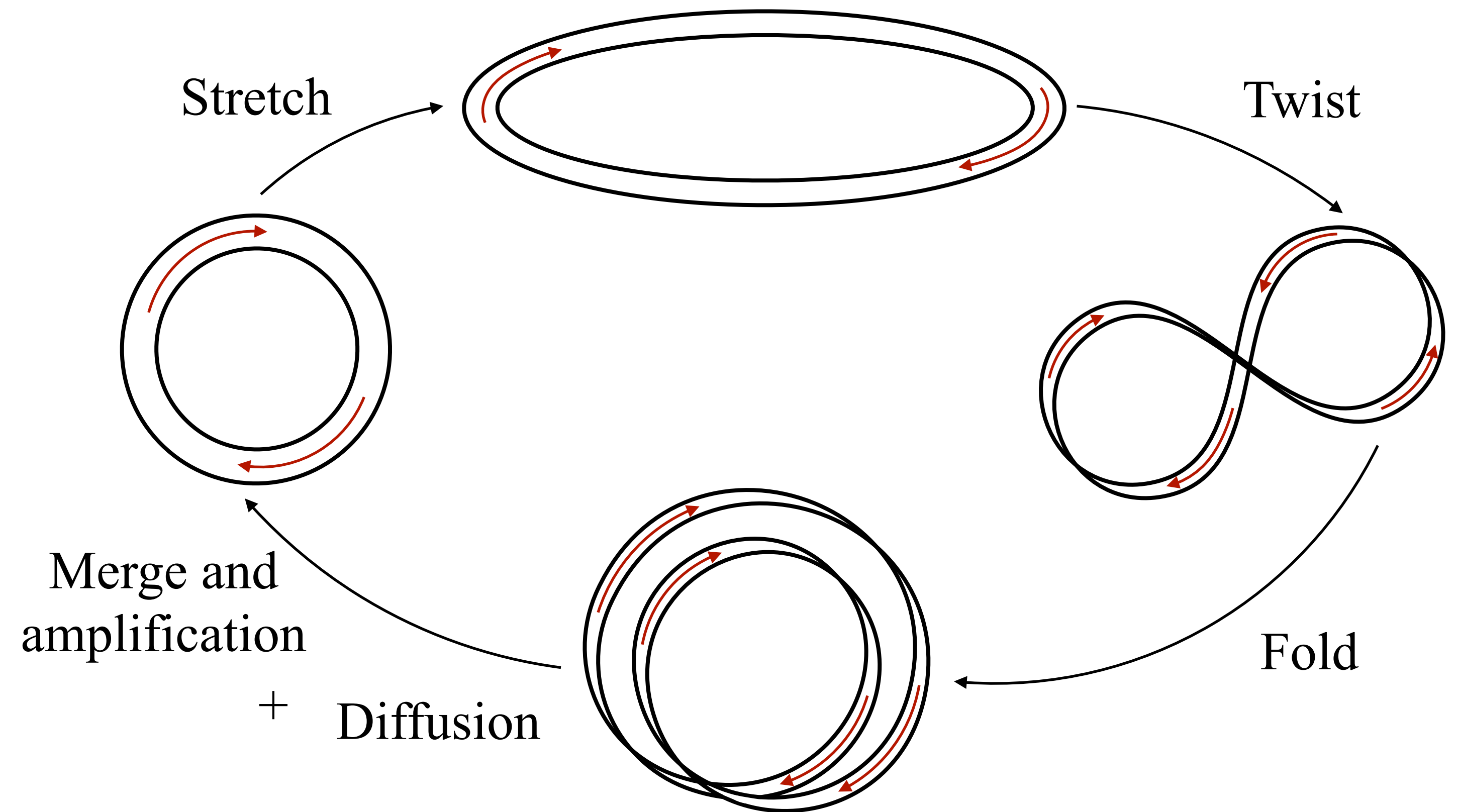
3D complex motions

Diversity of motion sources

=

Different dynamo mechanisms

**Vainshtein & Zel'dovich
fast dynamo process**



Astrophysical dynamos

Magnetic field

« Self-exciting fluid dynamos »
by K. Moffatt & E. Dormy

Reviews by F. Rincon 2019
and S. Tobias 2021

**3D complex
motions**

Diversity of motion sources

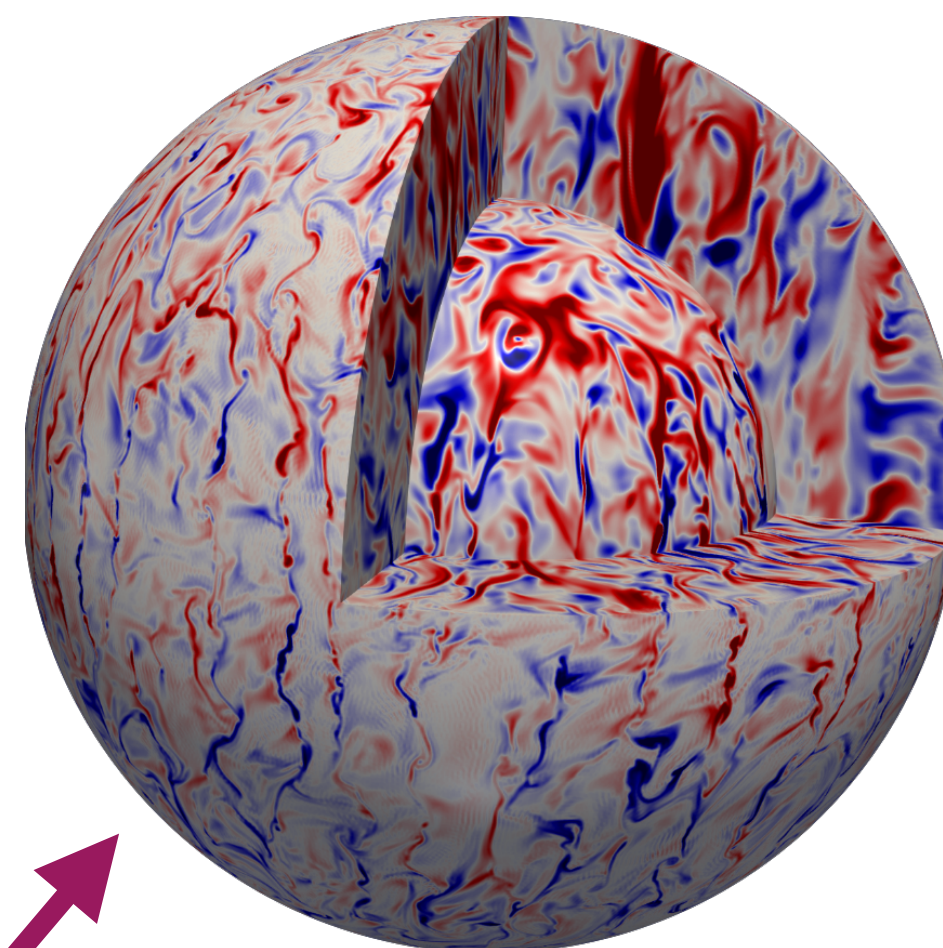
=

Different dynamo mechanisms

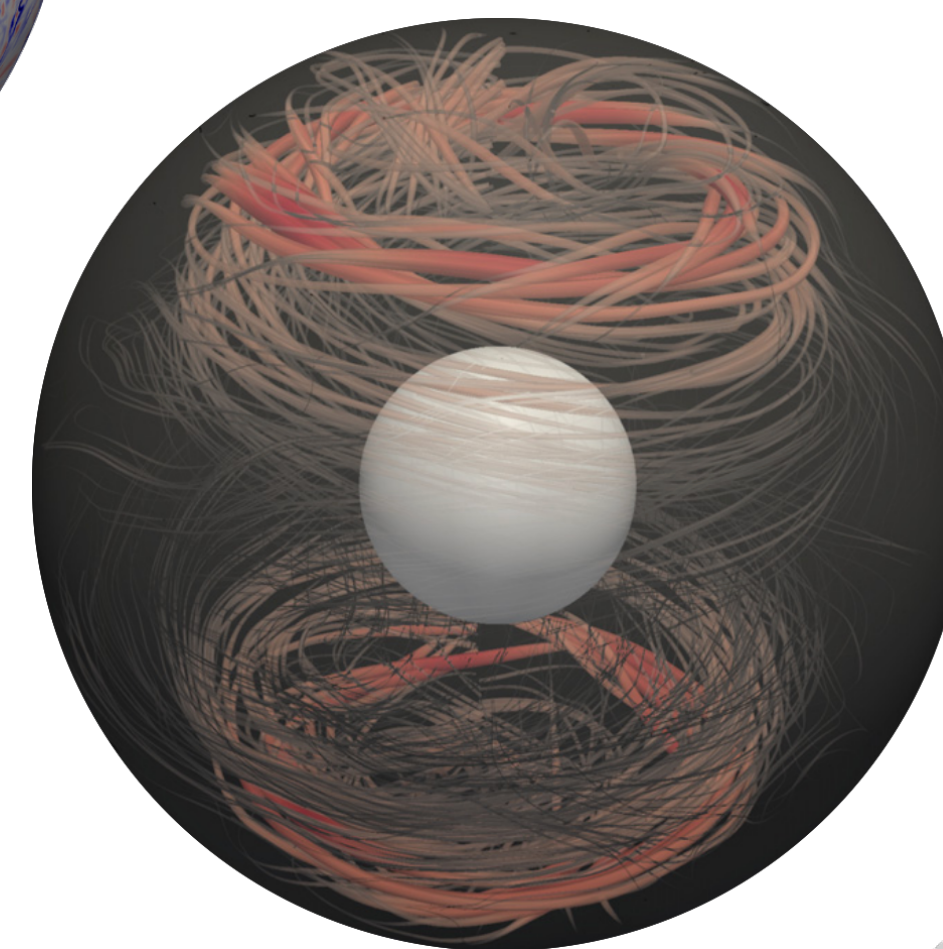
Convection

MRI

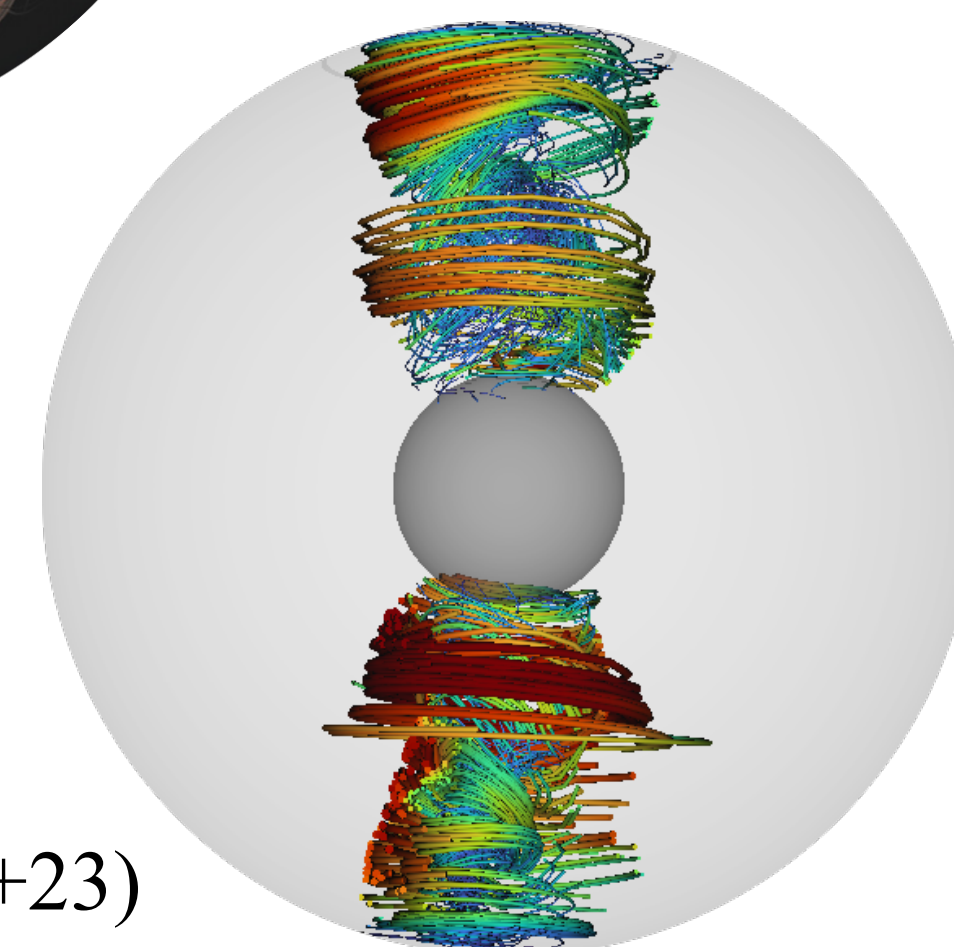
Taylor instability



(Zaire+22)



(Jouve+15)



(Barrère+23)

Main approximations of the modeling

Filter out acoustic waves

Boussinesq approximation:

- **Uniform** density profile
- Density perturbations **ONLY** in buoyancy term

Anelastic approximation:

- **Realistic** density profile
- Flow only produce **perturbations of a reference state**

Dimensionless parameters and stratification

Ekman number $E \equiv \frac{\nu}{d^2 \Omega_o}$

Rayleigh number $Ra \equiv \frac{N^2 d^4}{\nu \kappa}$

Thermal Prandtl number $Pr \equiv \frac{\nu}{\kappa}$

Magnetic Prandtl number $Pm \equiv \frac{\nu}{\eta}$

Same diffusivities for S and Y_e $\nearrow T \rightarrow \Theta$

Proto-neutron stars

$$N^2 \equiv -\frac{g}{\rho} \left(\left. \frac{\partial \rho}{\partial S} \right|_{P, Y_e} \frac{dS}{dr} + \left. \frac{\partial \rho}{\partial Y_e} \right|_{P, S} \frac{dY_e}{dr} \right)$$

Radiative zones

$$N^2 \equiv -\frac{g}{\rho} \left(\left. \frac{\partial \rho}{\partial S} \right|_{P, \mu} \frac{dS}{dr} + \left. \frac{\partial \rho}{\partial \mu} \right|_{P, S} \frac{d\mu}{dr} \right)$$

Volumetric forcing

Forcing of differential rotation

Forcing on the boundaries

- Impose fixed Ω_i and Ω_o
- Spherical-Couette configuration
- **Instabilities due to the set-up**

- $\mathbf{v} = \mathbf{u} + v_f \mathbf{e}_\phi$
with $v_f = r \sin \theta \Omega_f$ stationary field
- **Body force:** $\mathbf{f} = -u_\phi^{m=0}/t_f$
- **Forcing timescale:** $t_f \sim 10 \times t_\Omega$

Volumetric forcing in MagIC

Separation of variables: $\vec{V} = \vec{U} + \vec{V}_f$

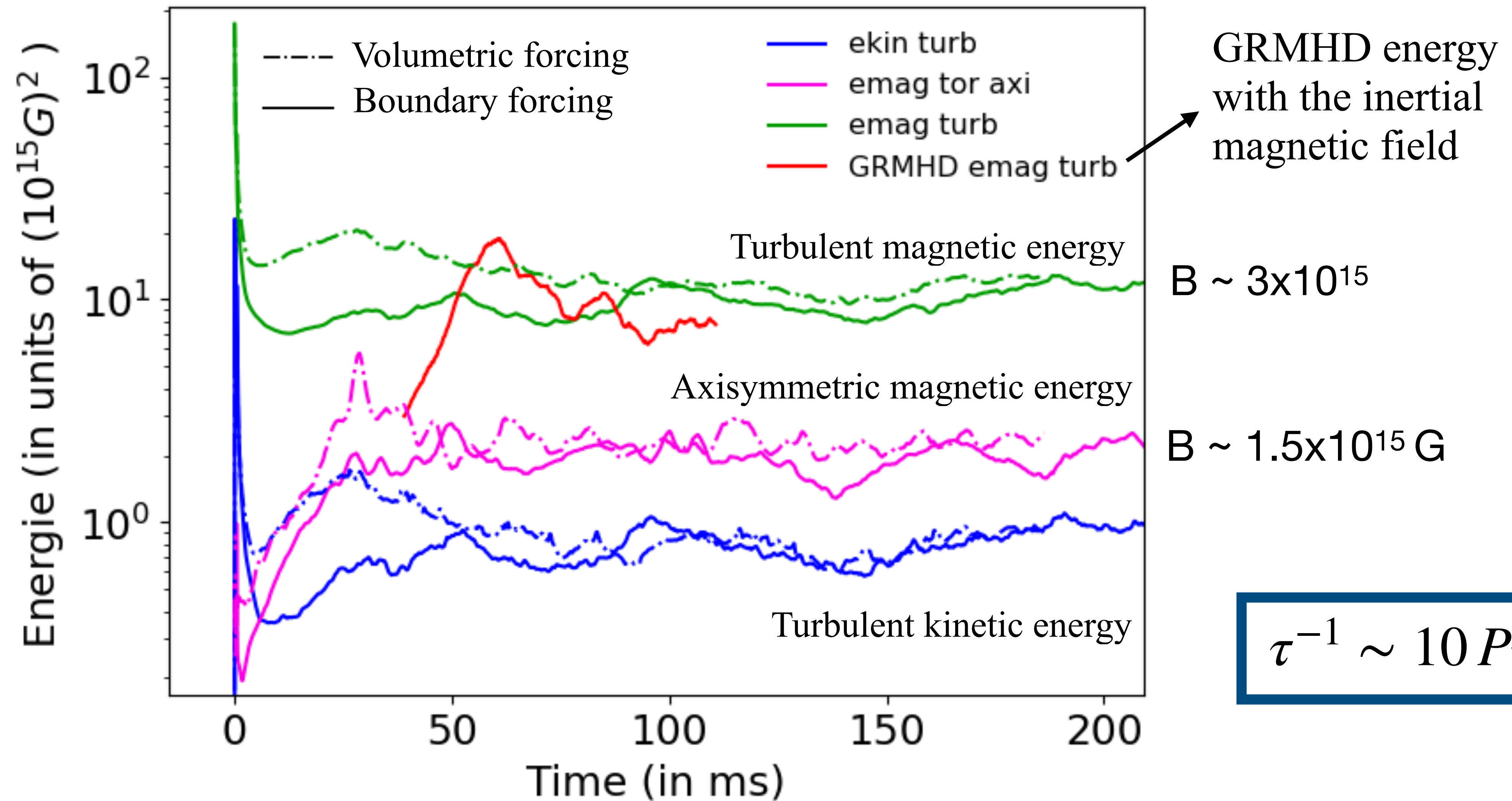
$$\left\{ \begin{array}{l} \vec{\nabla} \cdot (\bar{\rho} \vec{U}) = 0 \quad \text{et} \quad \vec{\nabla} \cdot \vec{B} = 0, \\ \\ \frac{\partial U}{\partial t} + \left((\vec{U} + \vec{V}_f) \cdot \vec{\nabla} \right) (\vec{U} + \vec{V}_f) + 2\Omega_0 \vec{e}_z \times (\vec{U} + \vec{V}_f) = \overset{\text{Coriolis}}{-\vec{\nabla} \left(\frac{\Pi'}{\bar{\rho}} \right)} + \overset{\text{Pressure}}{\frac{g_0 r_0^2}{C_p r^2} S' \vec{e}_r} + \overset{\text{Buoyancy}}{\frac{\nu}{\bar{\rho}} \vec{\nabla} \cdot \vec{\sigma}}, \\ \\ \overset{\text{“Dissipation” Forcing}}{-\tau U_{\phi, m=0}} + \overset{\text{Lorentz}}{\frac{1}{\mu_0 \bar{\rho}} [(\vec{\nabla} \times \vec{B}) \times \vec{B}]} + \overset{\text{Viscosity}}{\frac{\nu}{\bar{\rho}} \vec{\nabla} \cdot \vec{\sigma}}, \\ \\ \frac{\partial \vec{B}}{\partial t} = \overset{\text{Induction}}{\vec{\nabla} \times [(\vec{U} + \vec{V}_f) \times \vec{B}]} + \overset{\text{Resistivity}}{\eta \vec{\nabla}^2 \vec{B}}, \\ \\ \bar{\rho} \bar{T} \left(\overset{\text{Advection}}{\frac{\partial S'}{\partial t} + ((\vec{U} + \vec{V}_f) \cdot \vec{\nabla}) S'} + \overset{\text{Heat flux}}{(U_r + \vec{V}_f) \frac{d\bar{S}}{dr}} \right) = \overset{\text{Heating}}{\kappa \vec{\nabla} \cdot (\bar{\rho} \bar{T} \vec{\nabla} S')} + \nu Q_\nu + \frac{\eta}{\mu_0^2} (\vec{\nabla} \times \vec{B})^2 \end{array} \right.$$

Implemented in MagIC
by A. Reboul-Salze

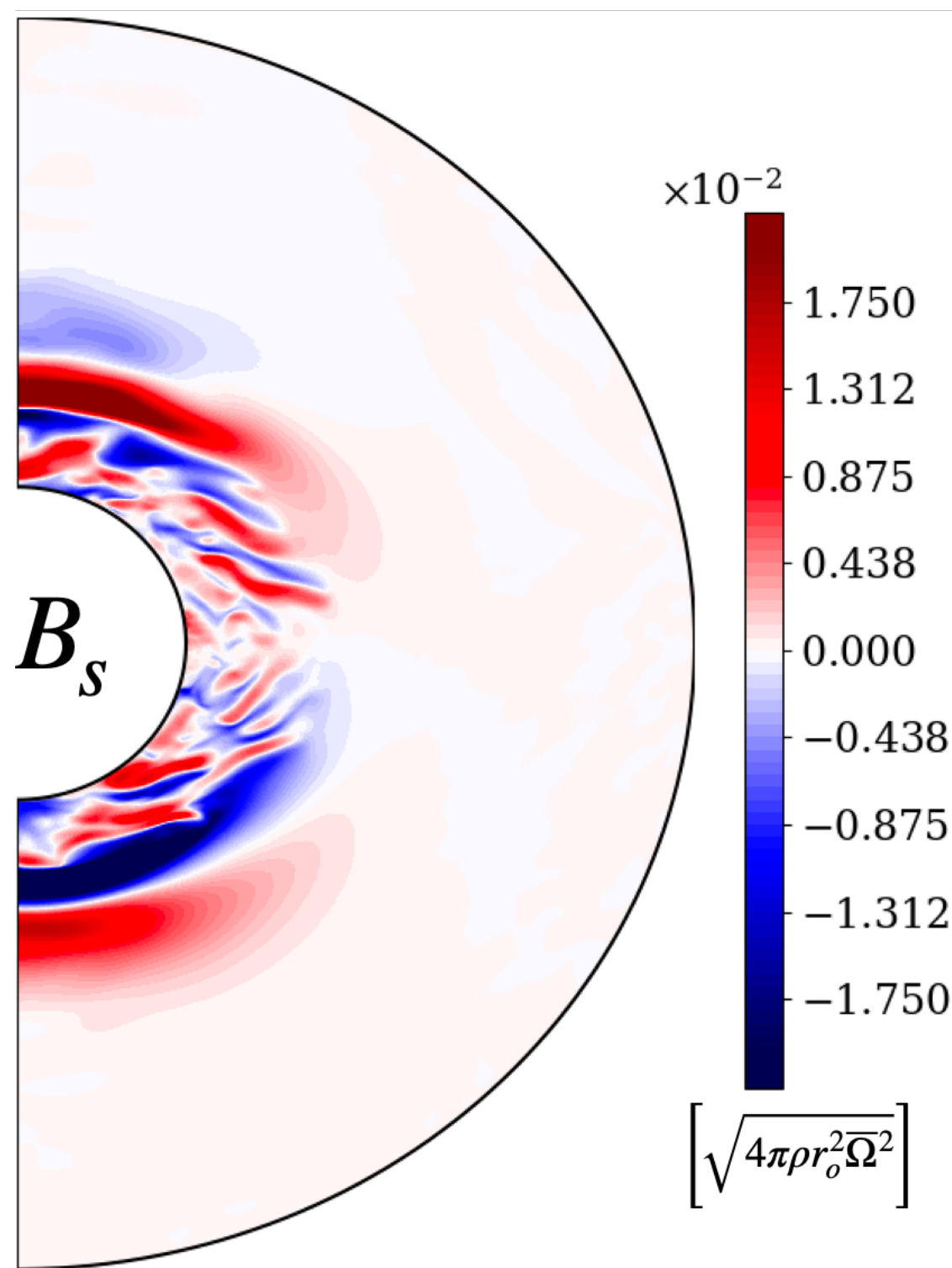
➔ More efficient forcing than just forcing at the outer boundaries

In the following $\tau^{-1} \sim 10 P_{\overline{\Omega}}$

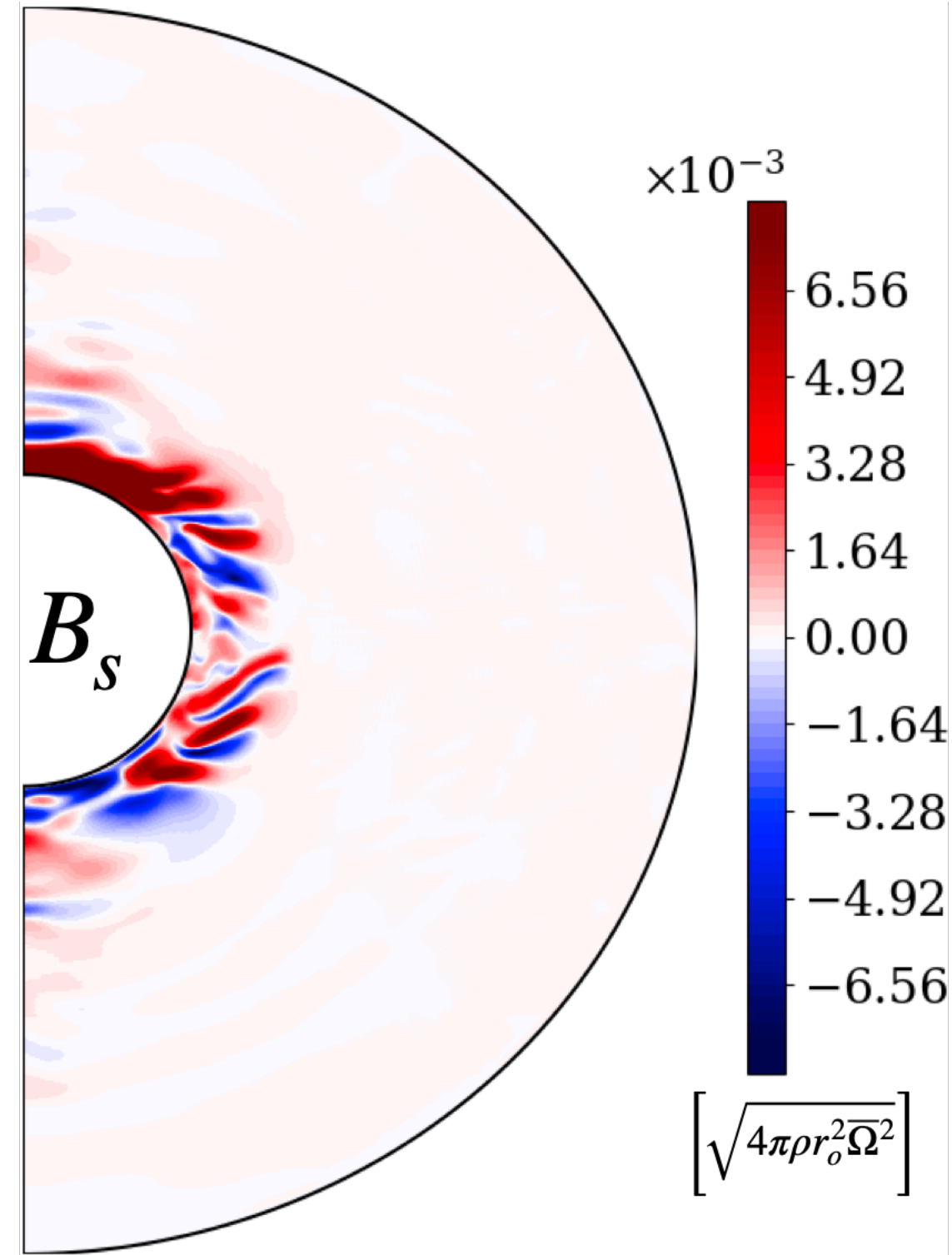
Boundary vs. volumetric forcing: case for MRI



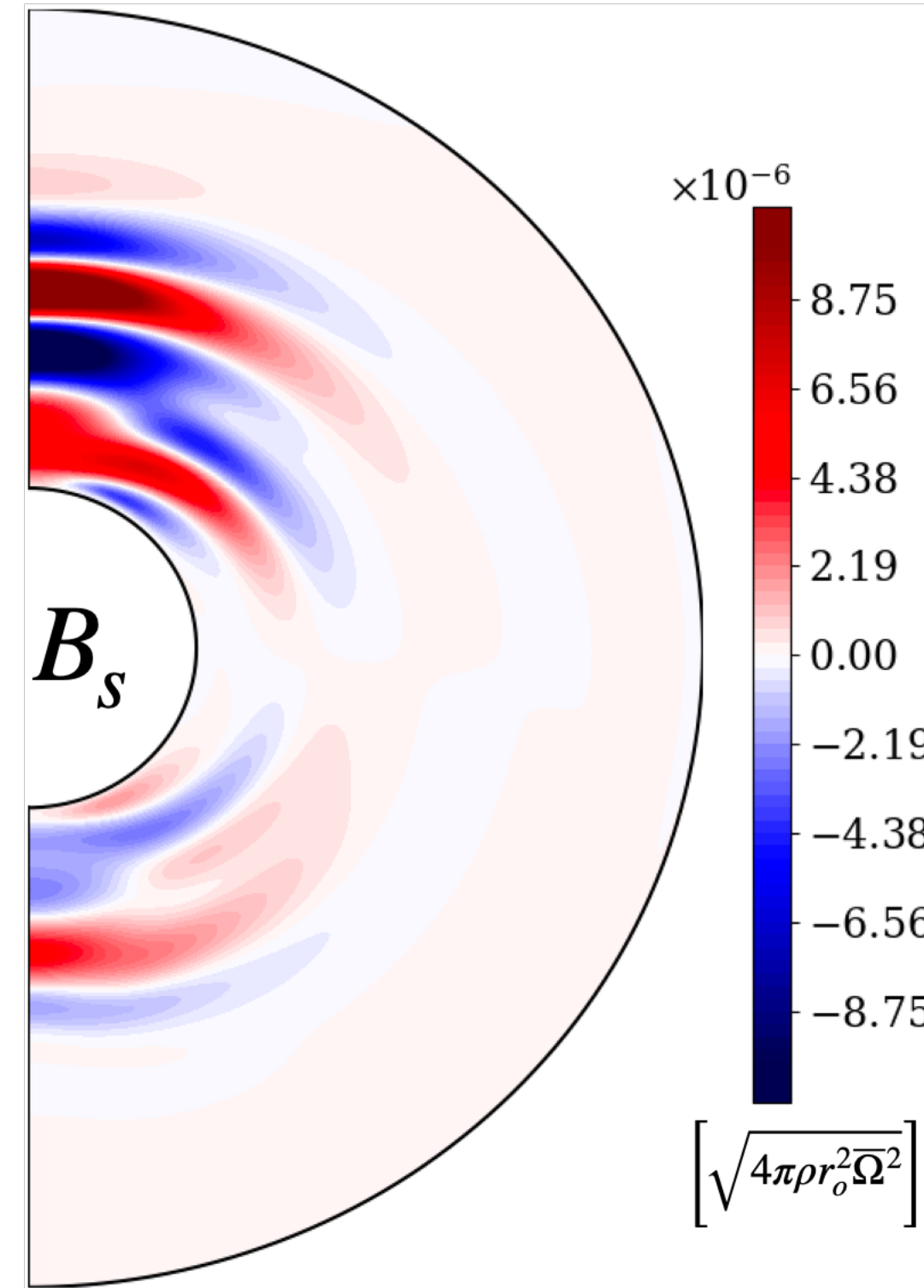
Two different dynamos



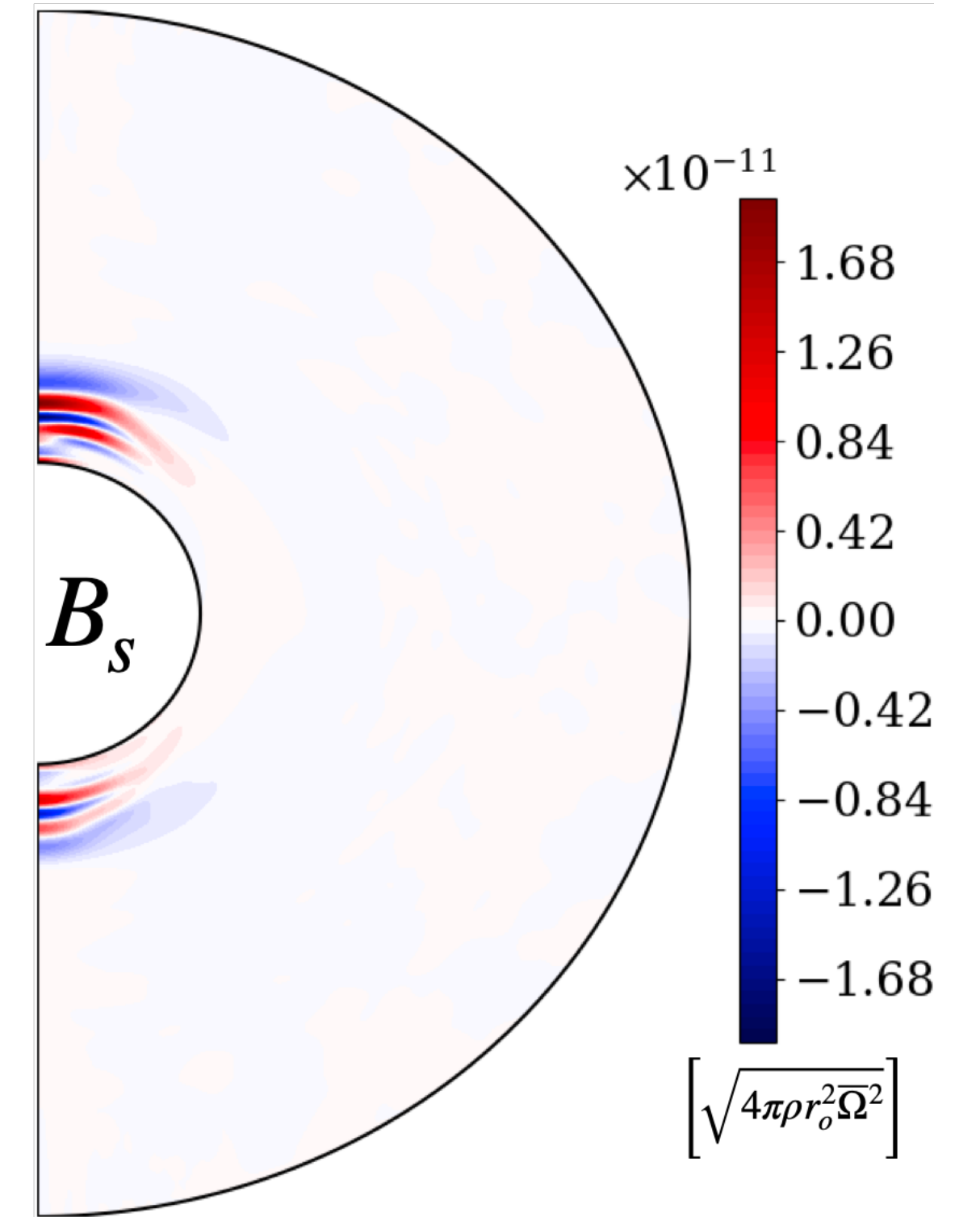
$q = -1$
 $B_\phi^{m=0} \propto 1/r$



$q = -1$
 $B_\phi^{m=0} \propto 1/r^2$

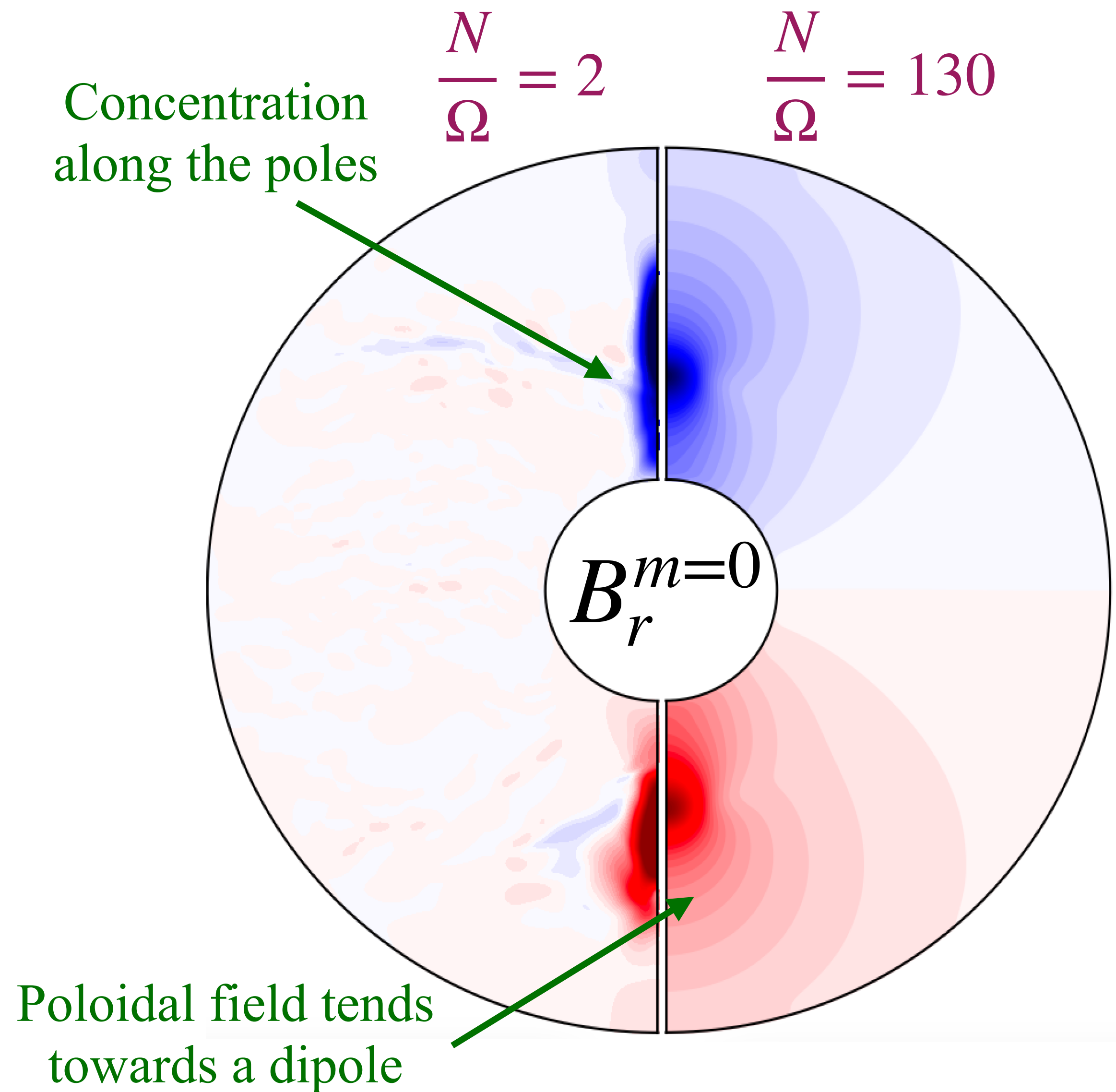
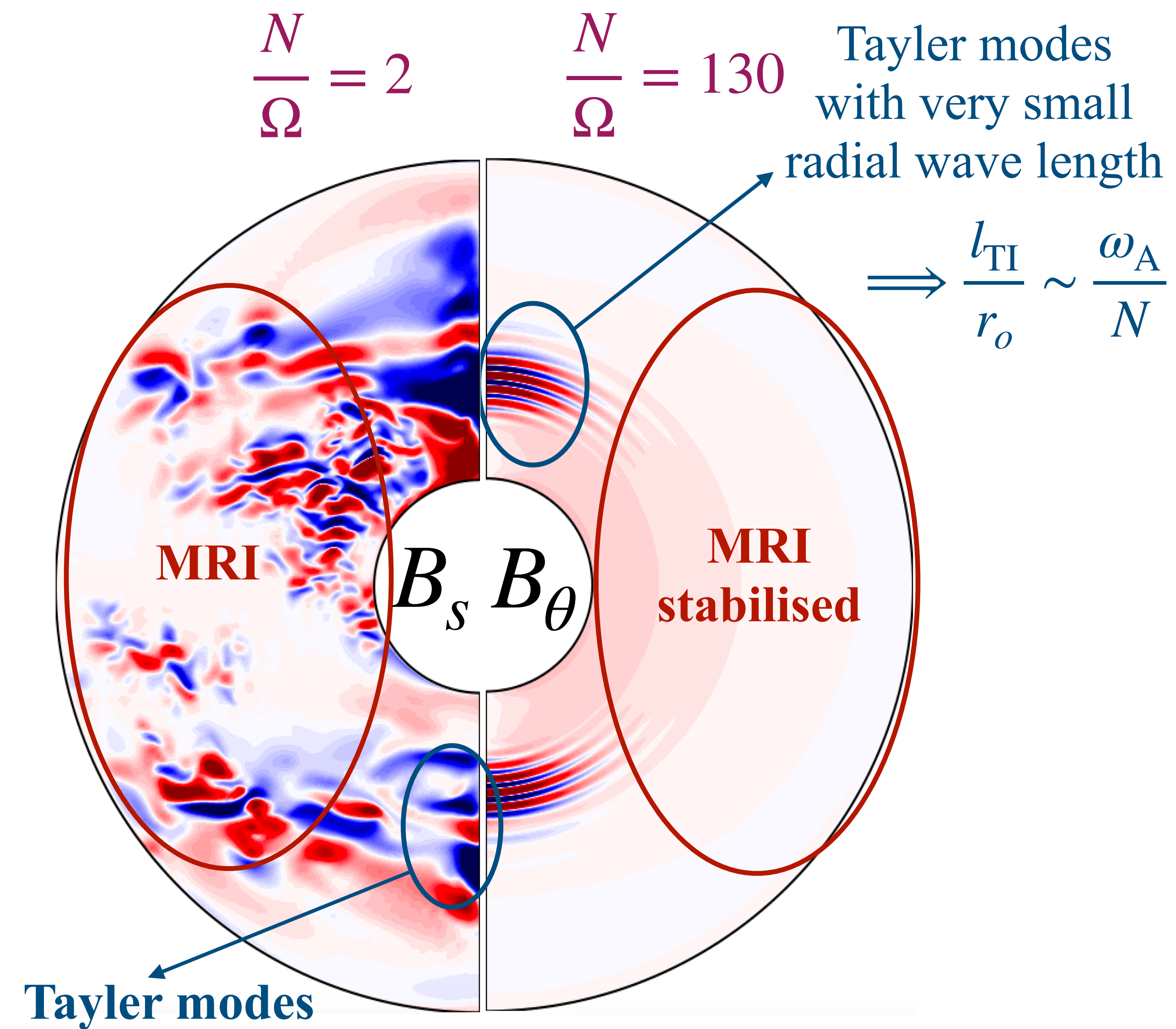


$q = 0$
 $B_\phi^{m=0} \propto 1/r$

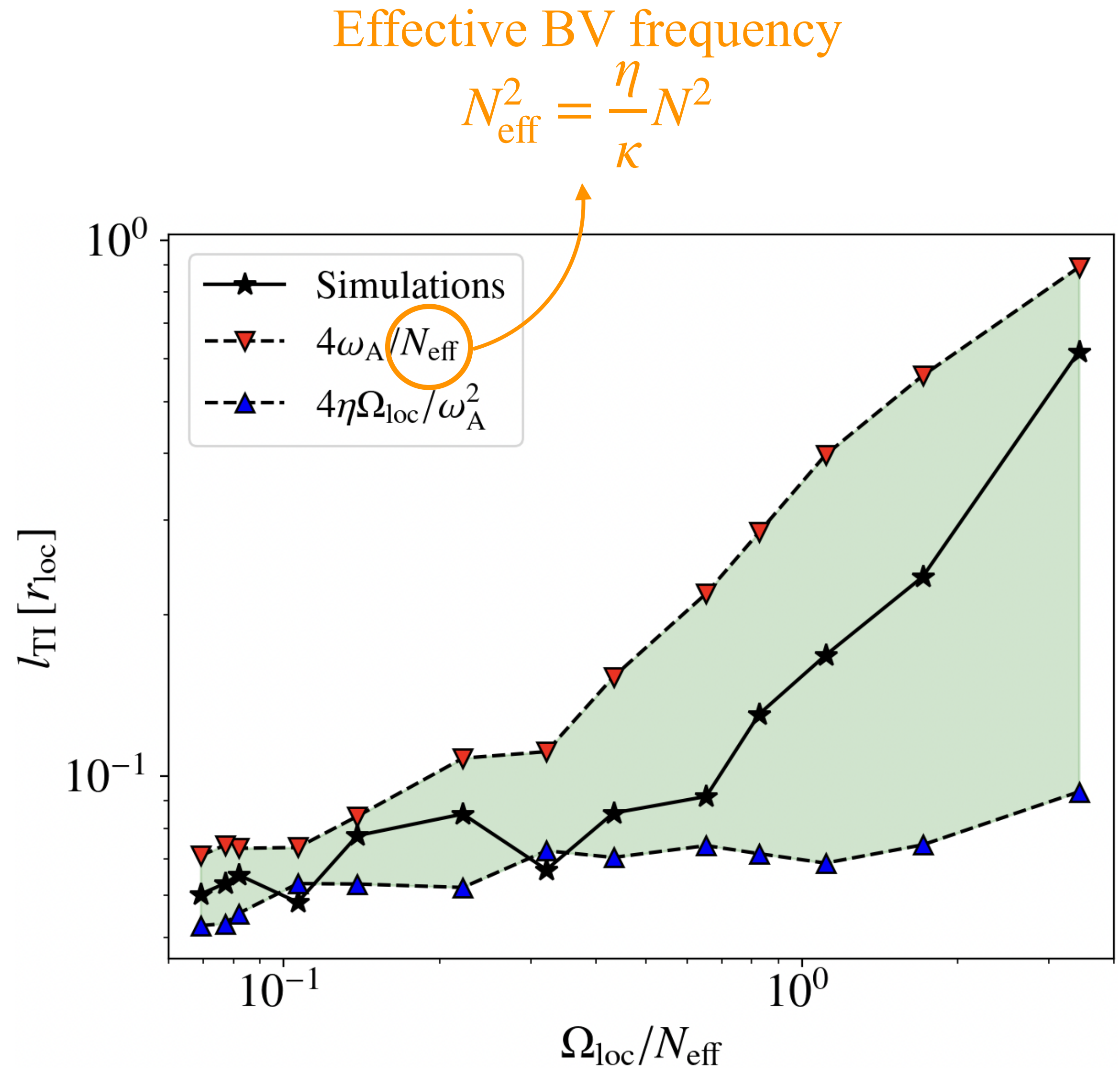
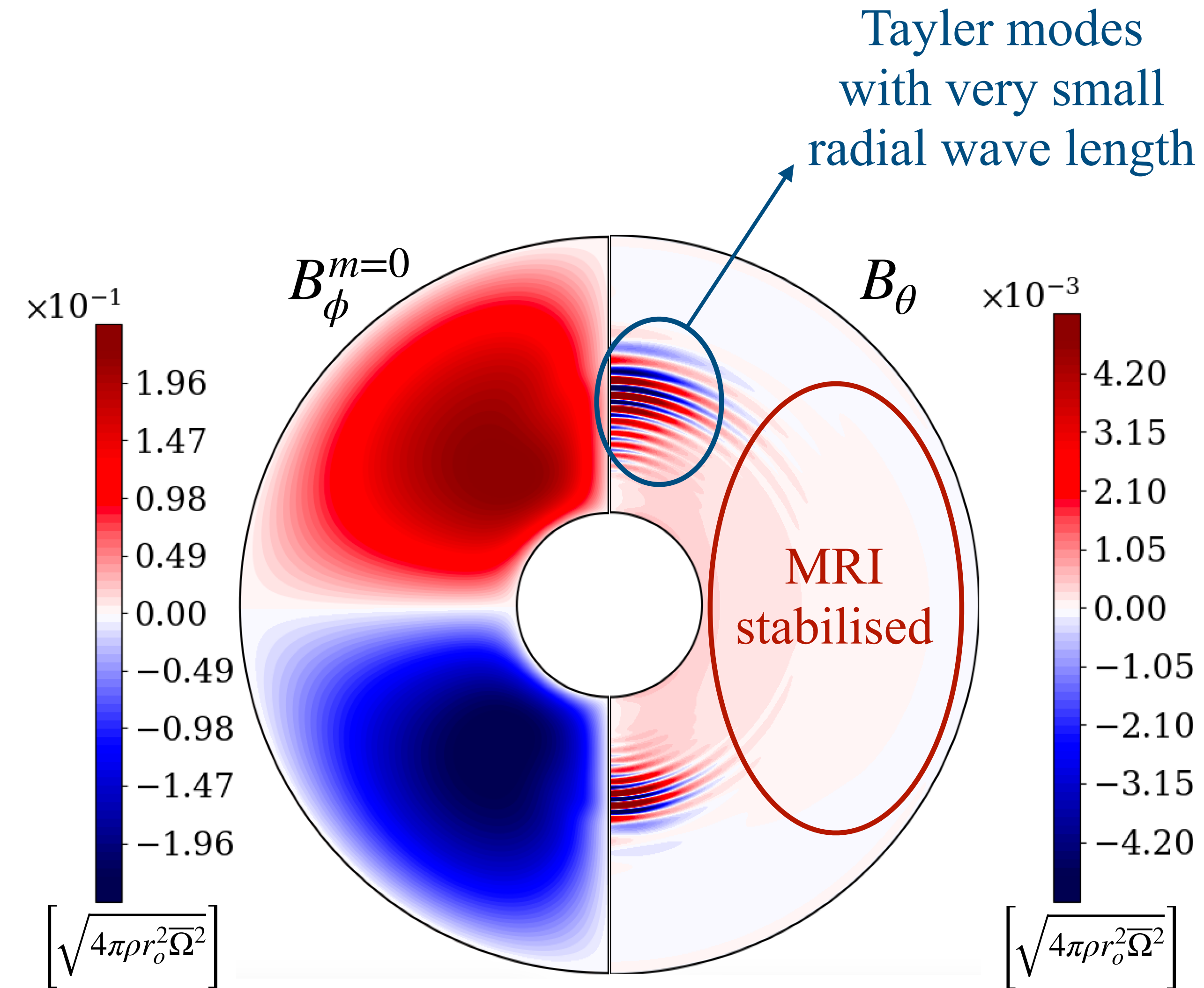


$q = 0$
 $B_\phi^{m=0} \propto 1/r^2$

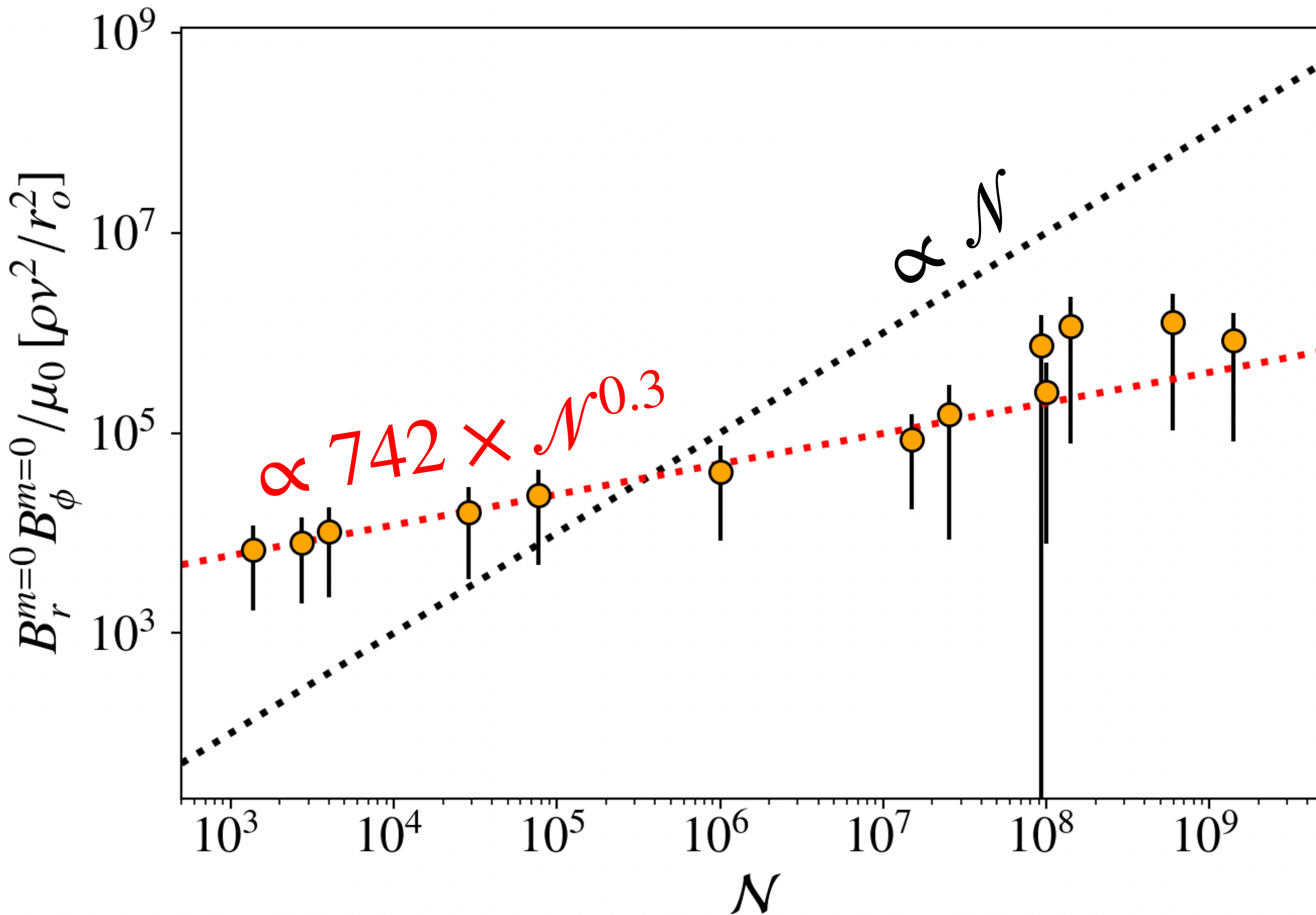
Dipolar branch in a large range of stratifications



Impact on the Tayler mode radial length scale



Scaling Maxwell torque as in Petitdemange+23



$$\mathcal{N} = \beta q^3 \frac{r_o^4 \overline{\Omega}^2}{\nu^2} \left(\frac{\overline{\Omega}}{N} \right)^4$$

with $q = \frac{u_\phi^{\text{loc}}}{l_{\text{TI}} \overline{\Omega}}$ and $\beta = 0.1$

Comparison with Petitdemange et al. 2023, 2024

This study

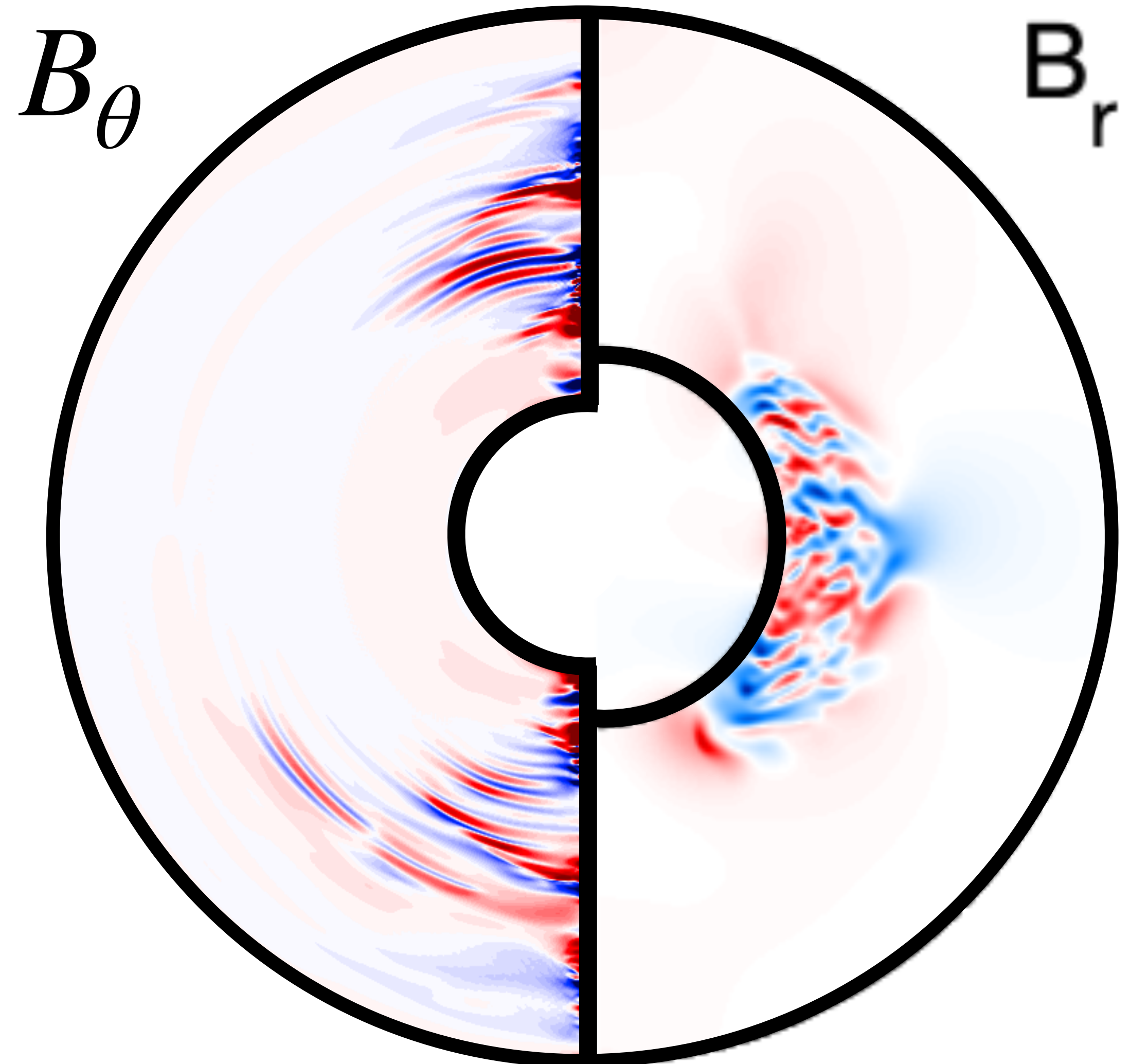
Petitdemange et al.
2023, 2024:

Setup differences:

- | | |
|------------------|--------------------|
| ▶ Volum. forcing | ▶ Boundary forcing |
|------------------|--------------------|

Mag. Field differences:

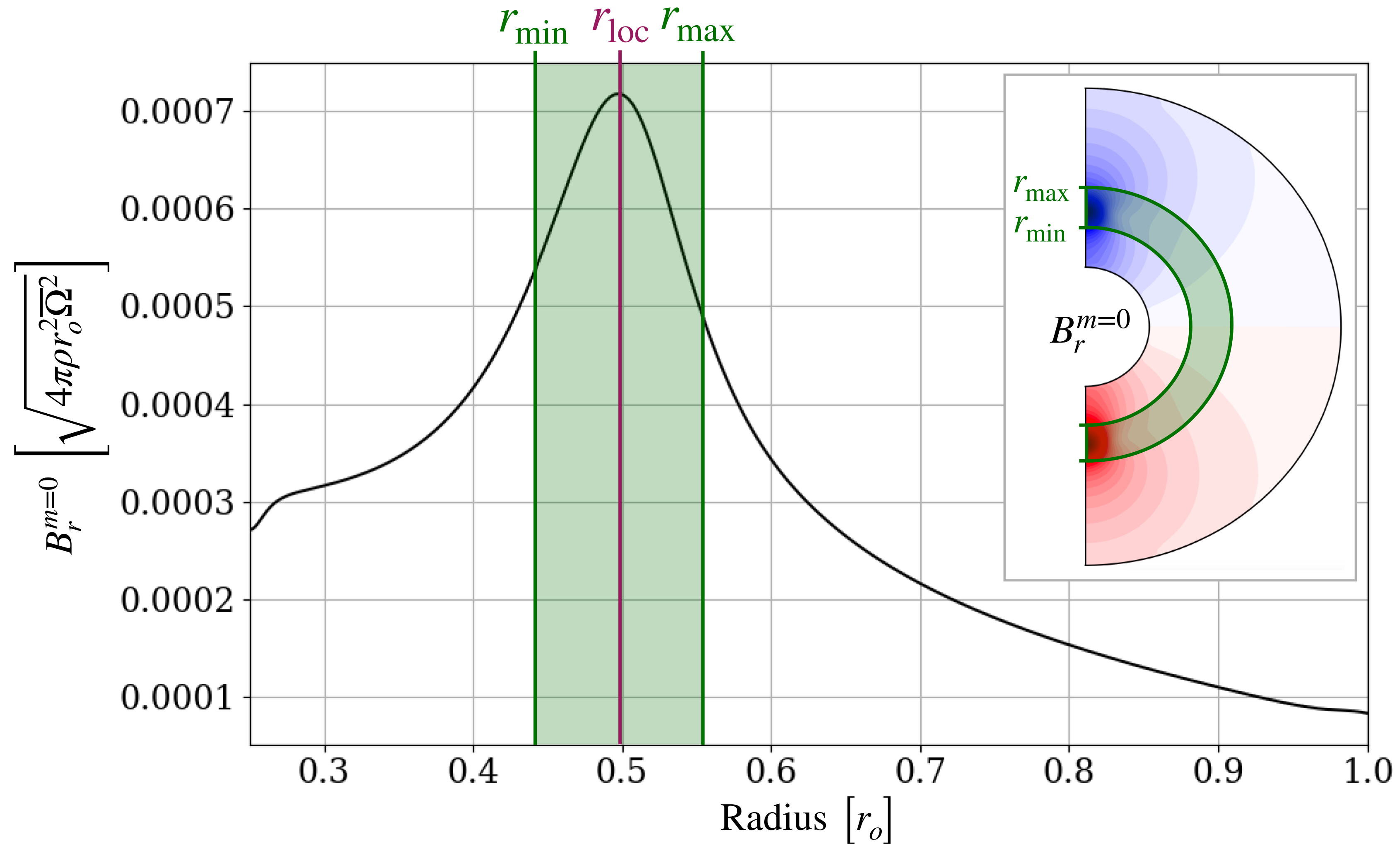
- | | |
|--------------------------------------|---------------------------------------|
| ▶ Near the poles | ▶ Near the equator |
| ▶ More efficient
transport | ▶ Spruit 2002's scaling
law |



This study

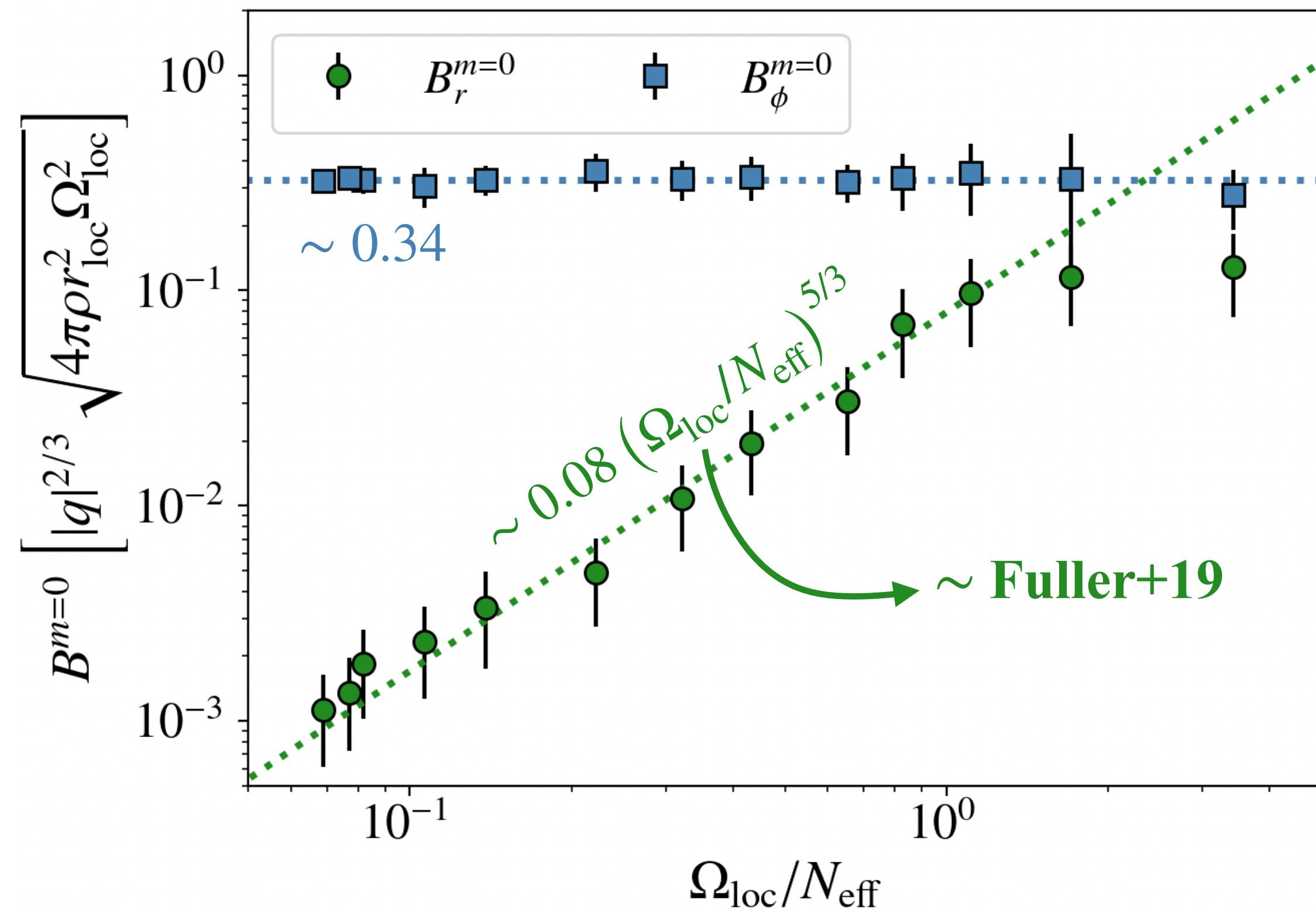
Petitdemange et al. 2024

Radial profile after latitudinal and time average

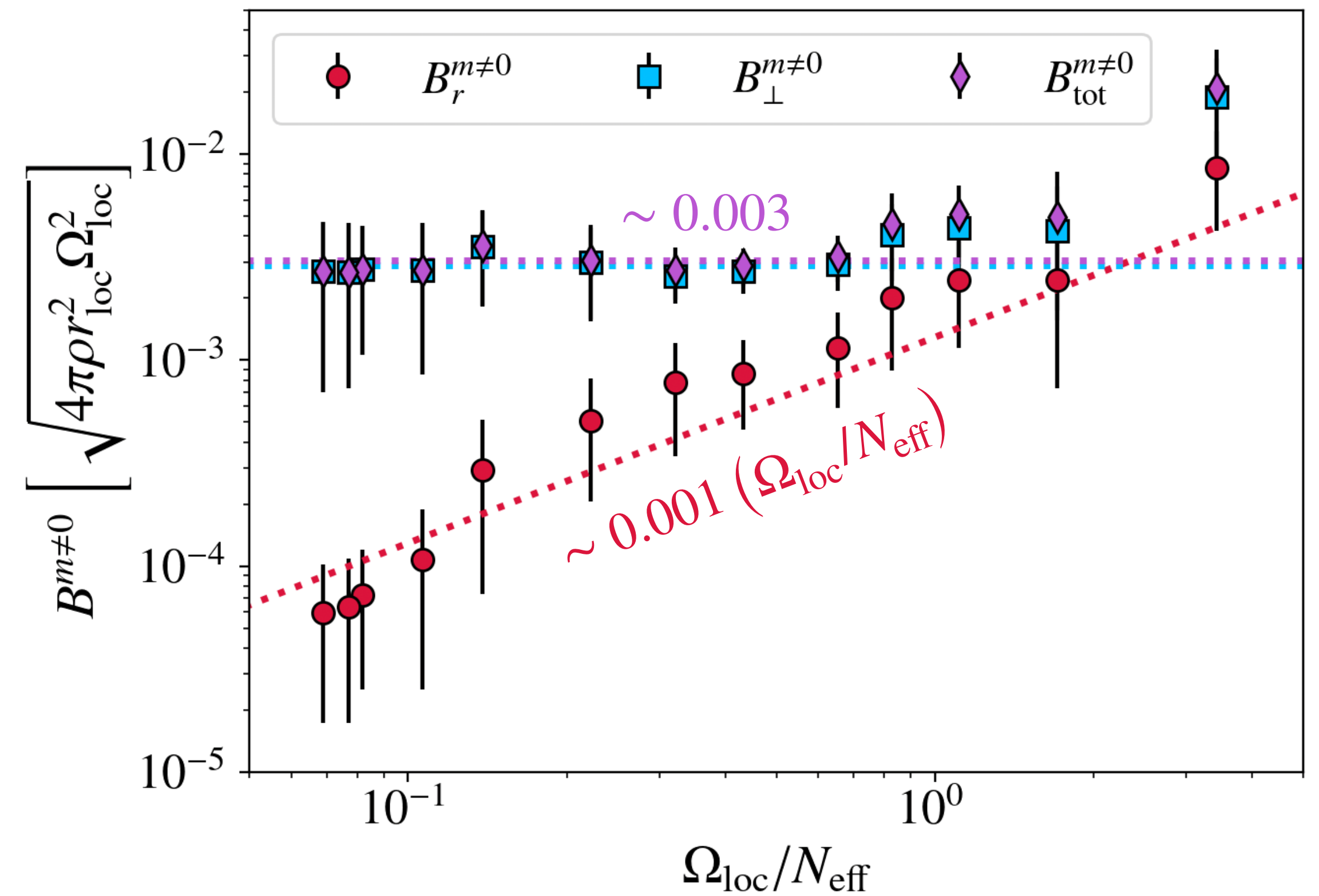


Magnetic field scaling laws

Large-scale magnetic fields



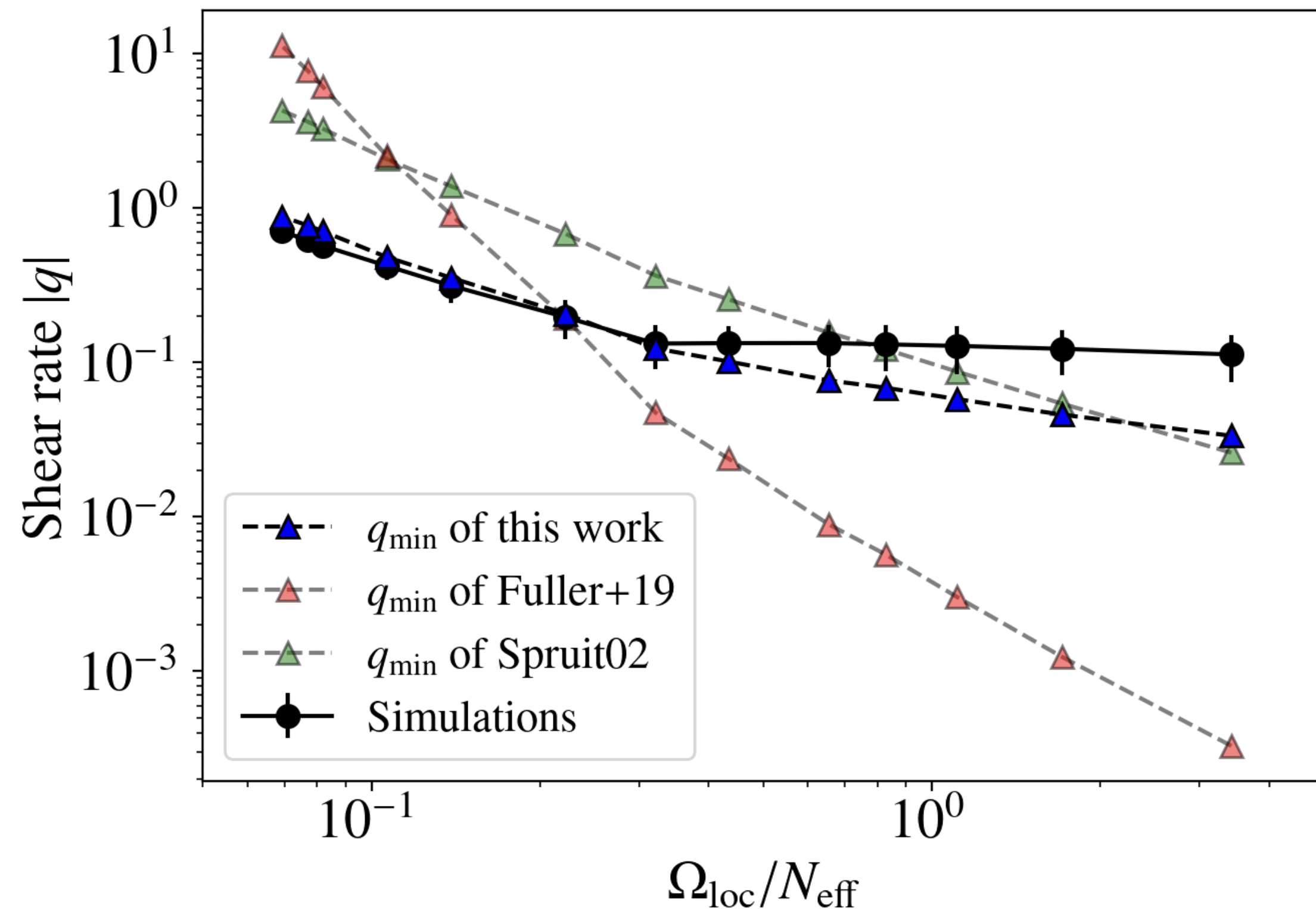
Turbulent magnetic fields



Angular momentum transport

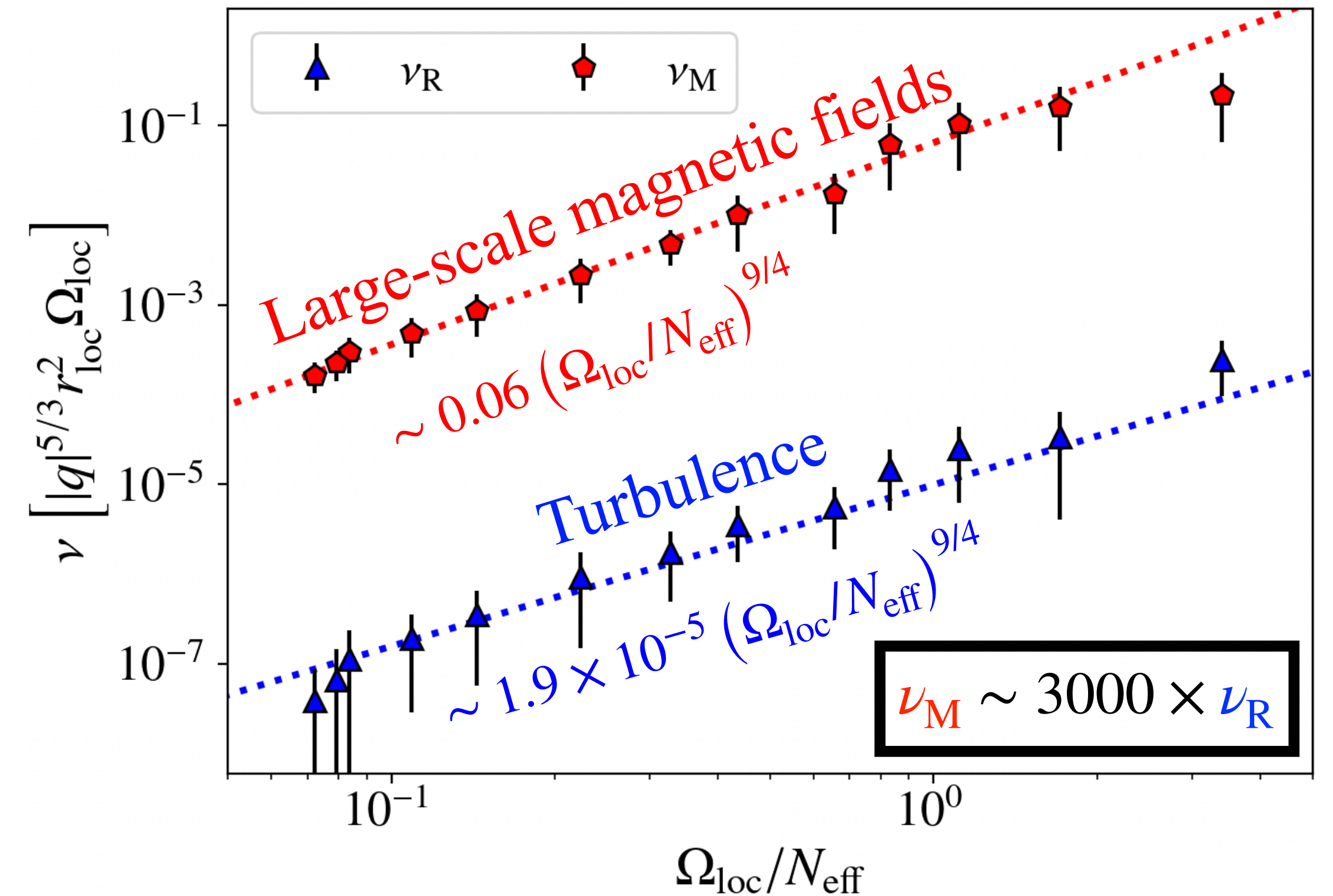
New expression of $q_{\min}$:

$$q_{\min} \sim 5.2 \left(\frac{N_{\text{eff}}}{\Omega_{\text{loc}}} \right)^{3/4} \left(\frac{\eta}{r_{\text{loc}}^2 \Omega_{\text{loc}}} \right)^{3/8}$$



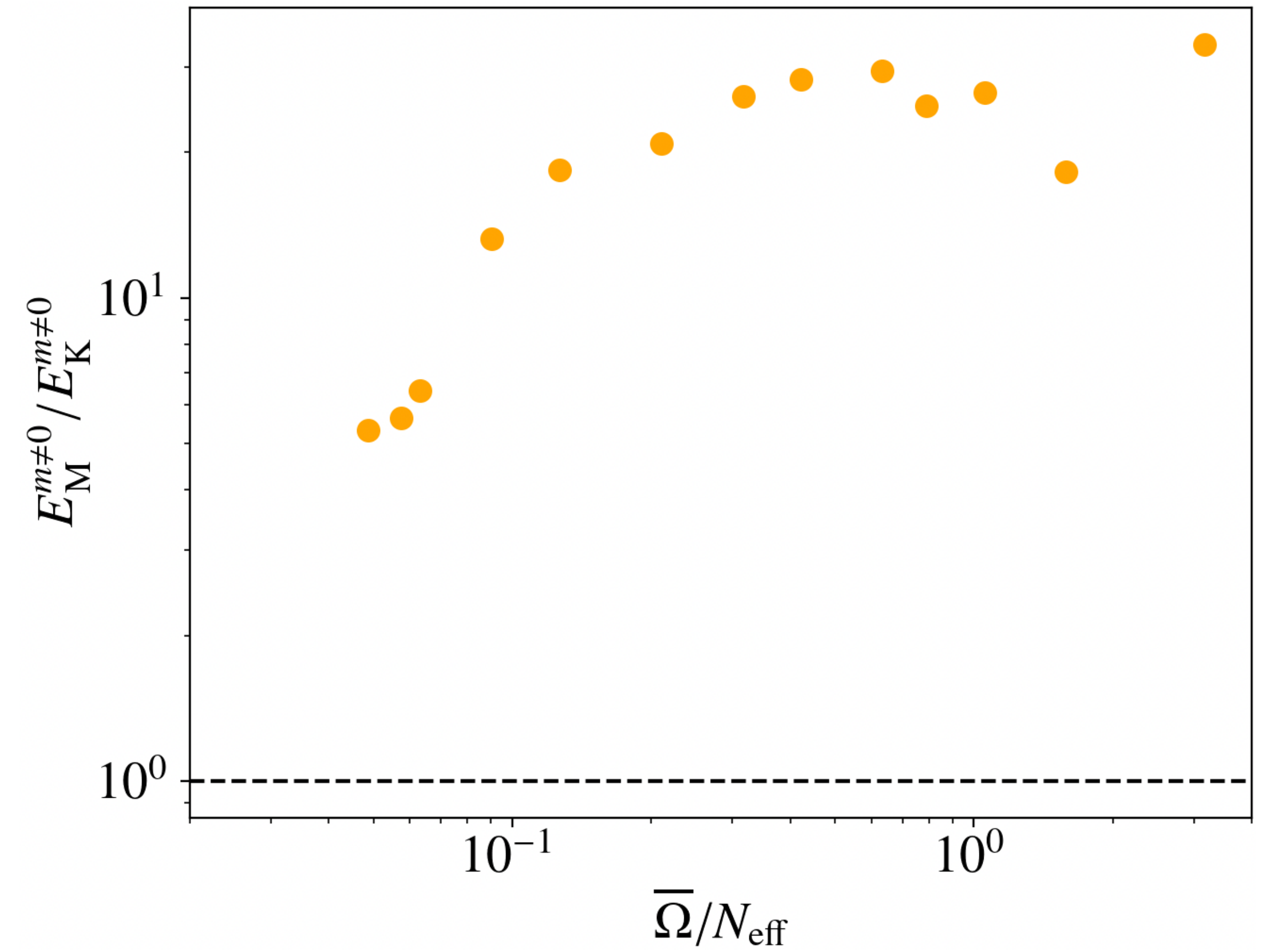
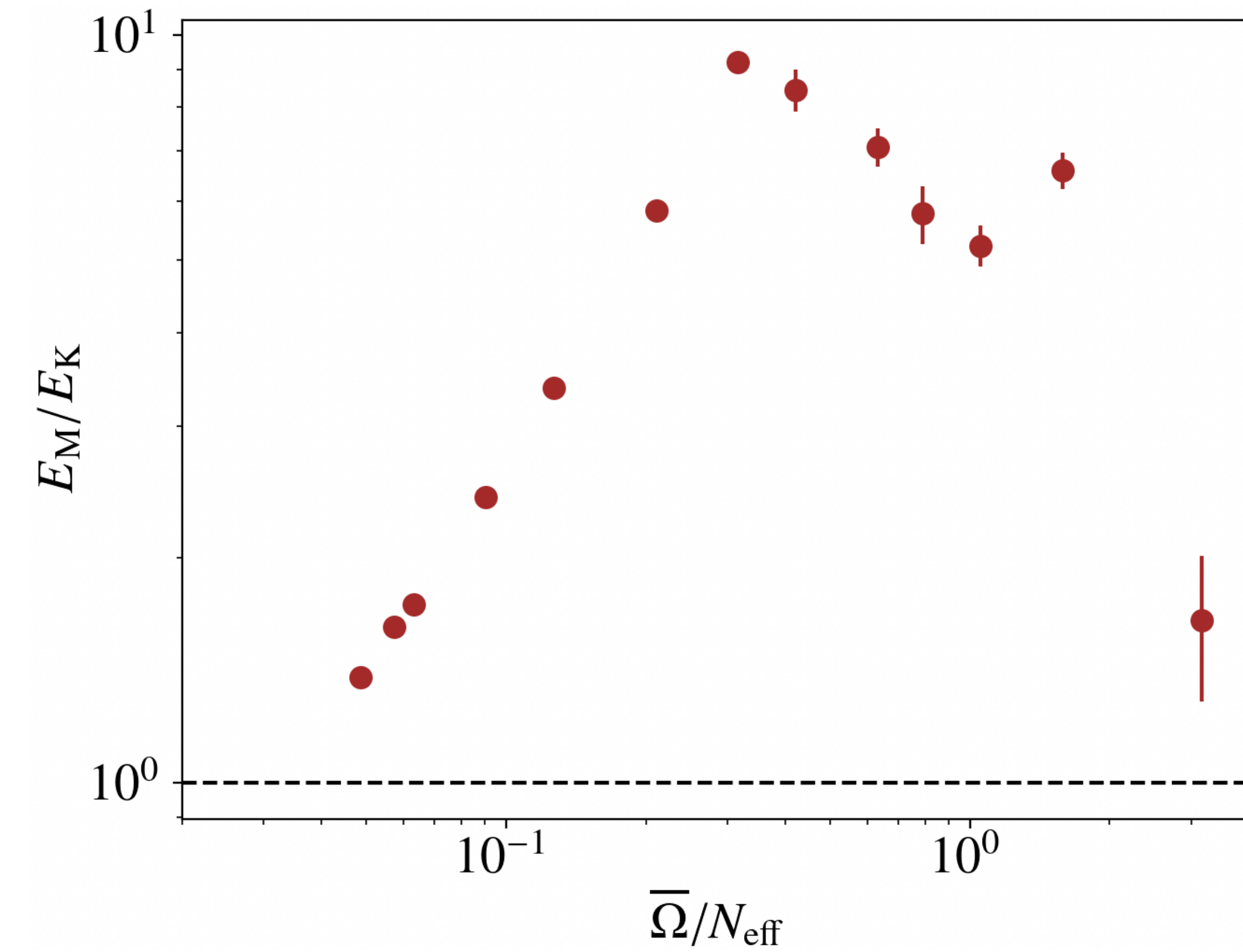
Fuller+19:

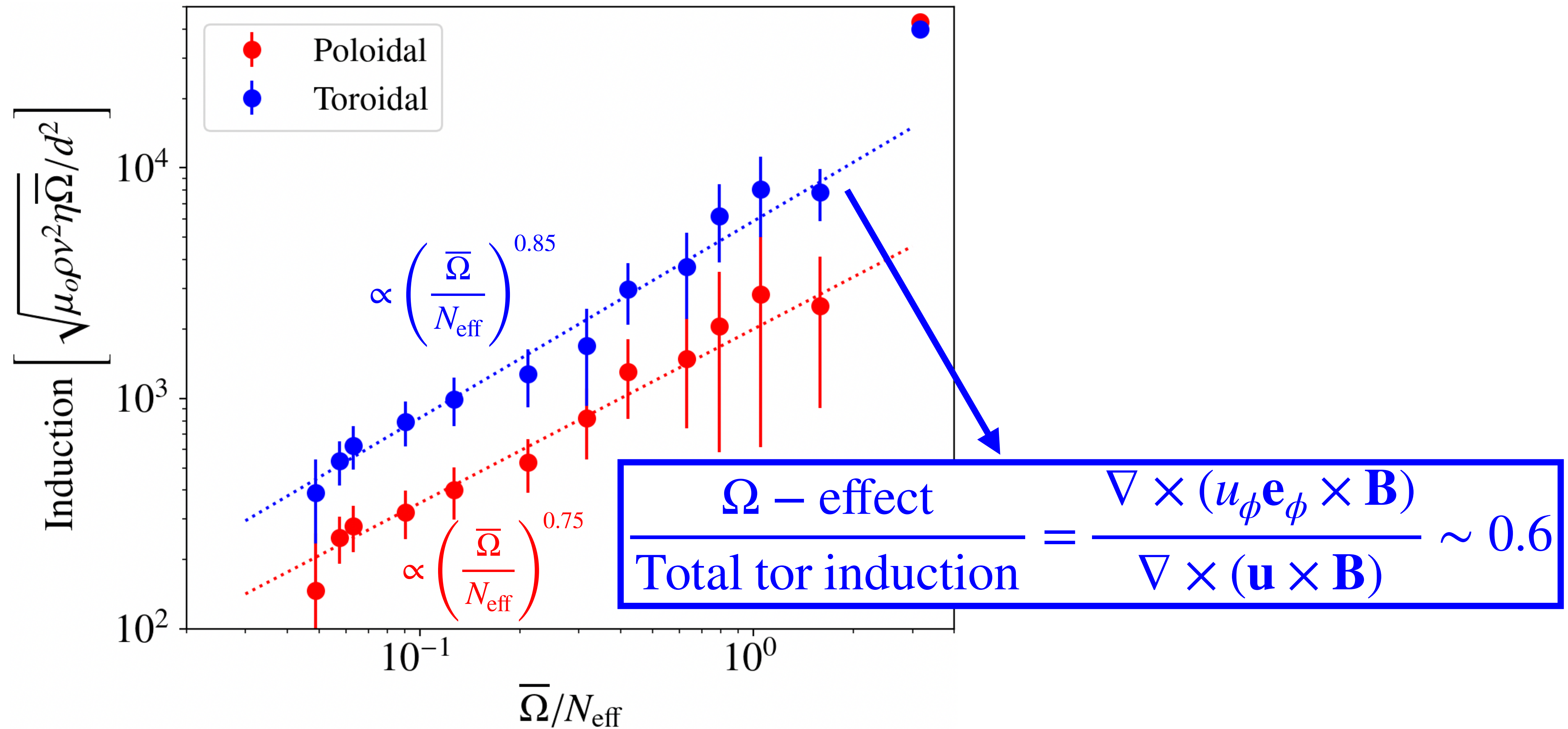
$$\nu_{\text{M}} \sim 0.06 q r_{\text{loc}}^2 \Omega_{\text{loc}} \left(\Omega_{\text{loc}}/N_{\text{eff}} \right)^2$$



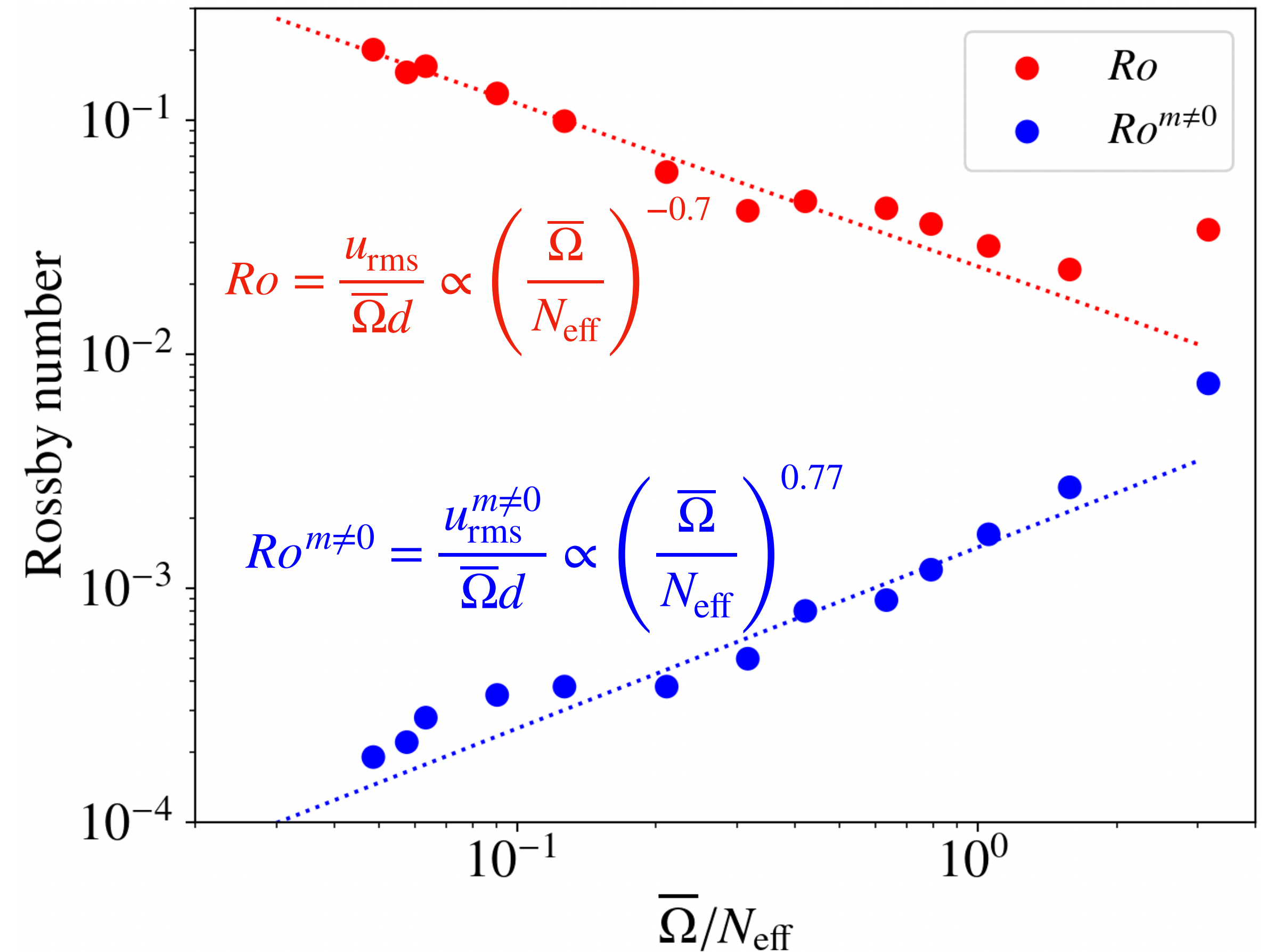
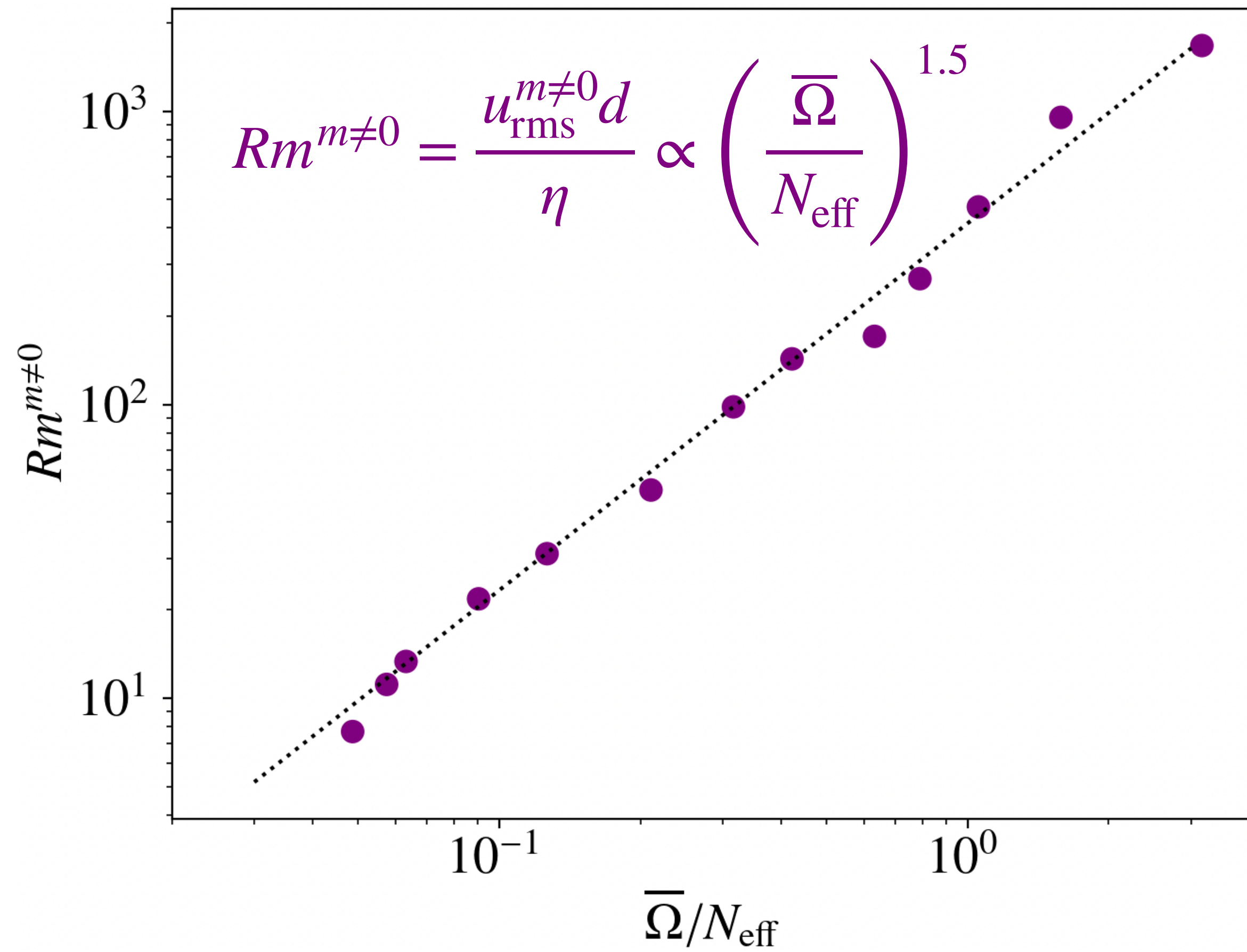
⇒ Slightly less efficient AM transport than Fuller+19,
but much lower $q_{\min}$ in red giants

Super-equipartition

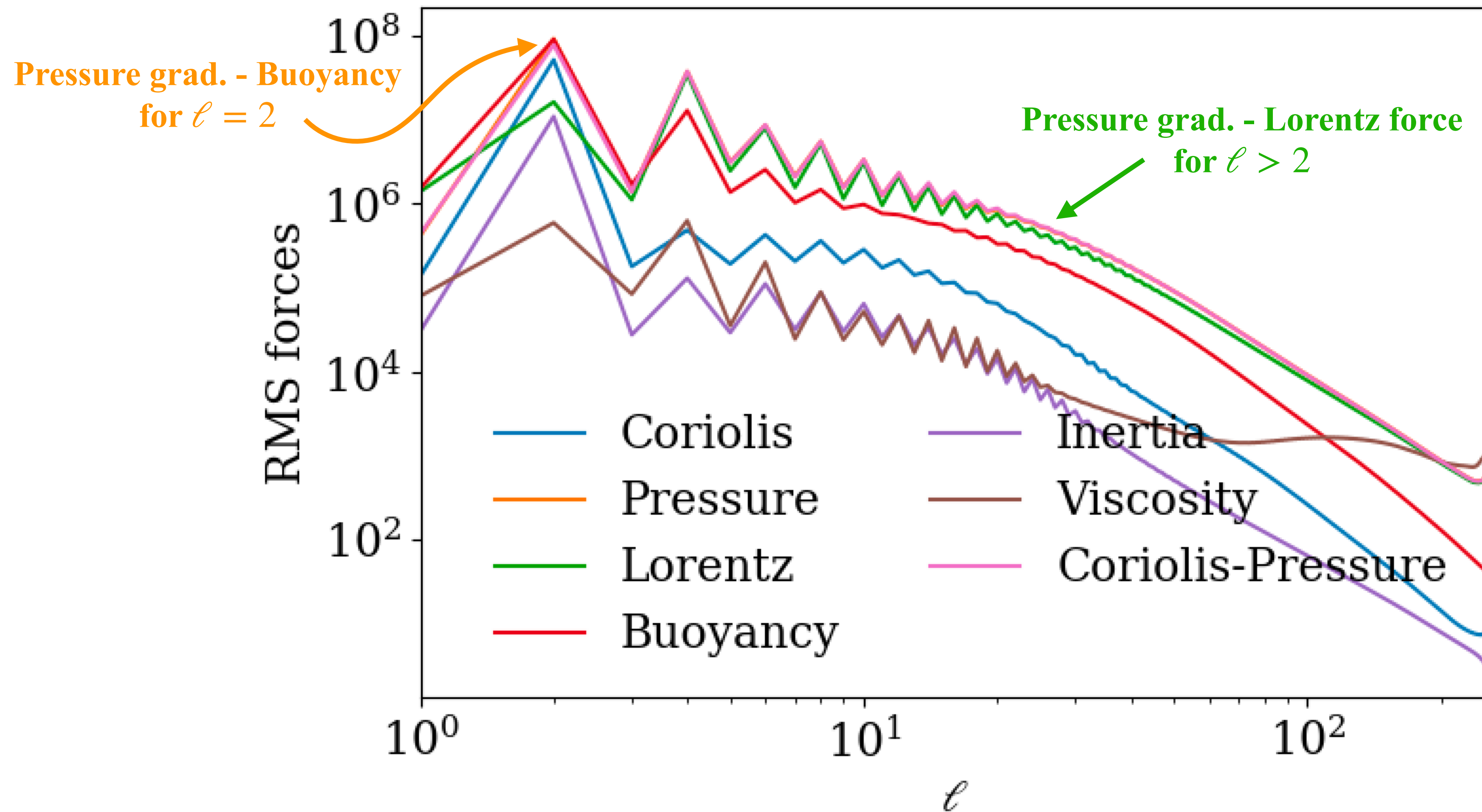




Critical Rossby and magnetic Reynolds numbers



Force balance spectrum

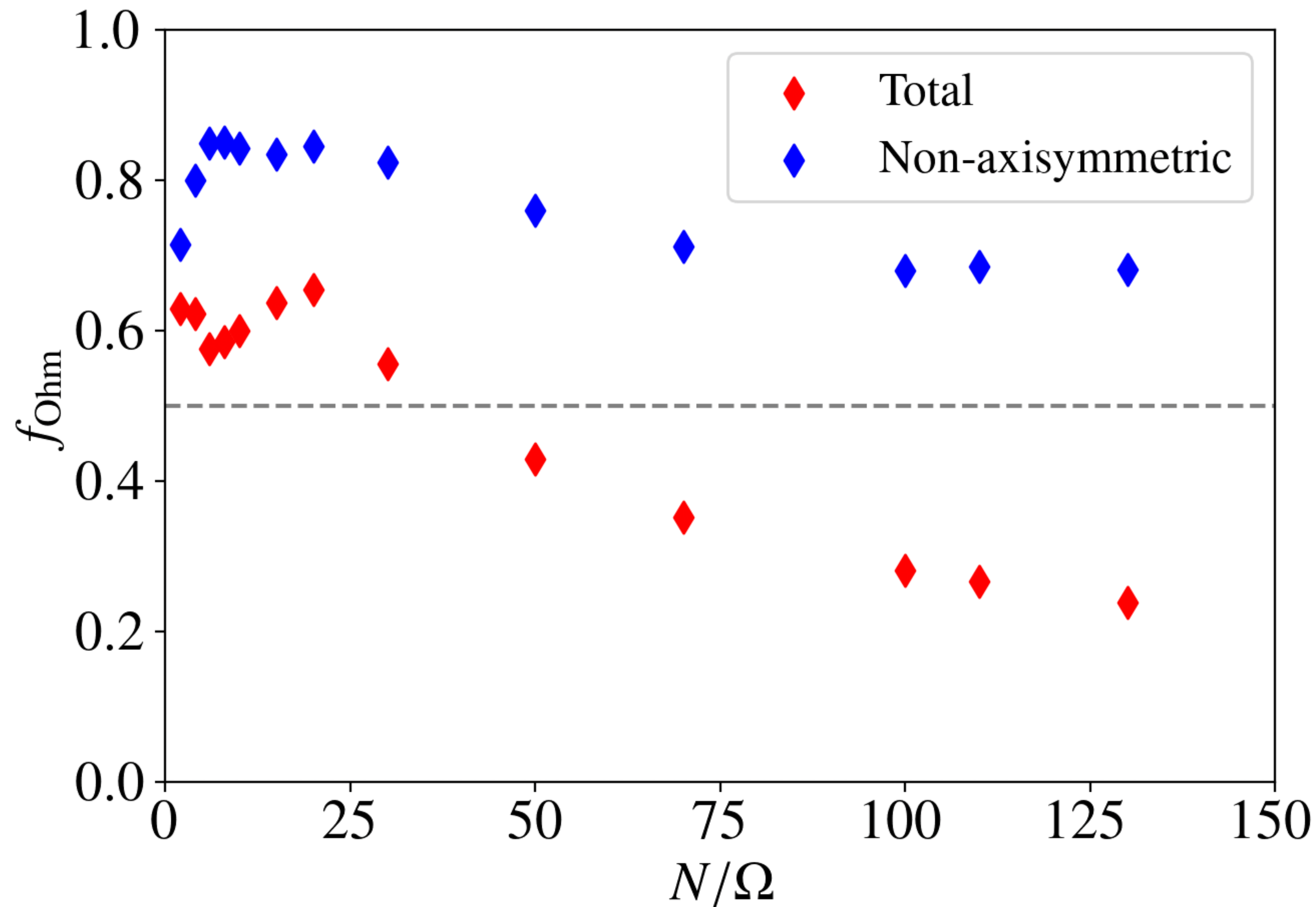


Dissipations

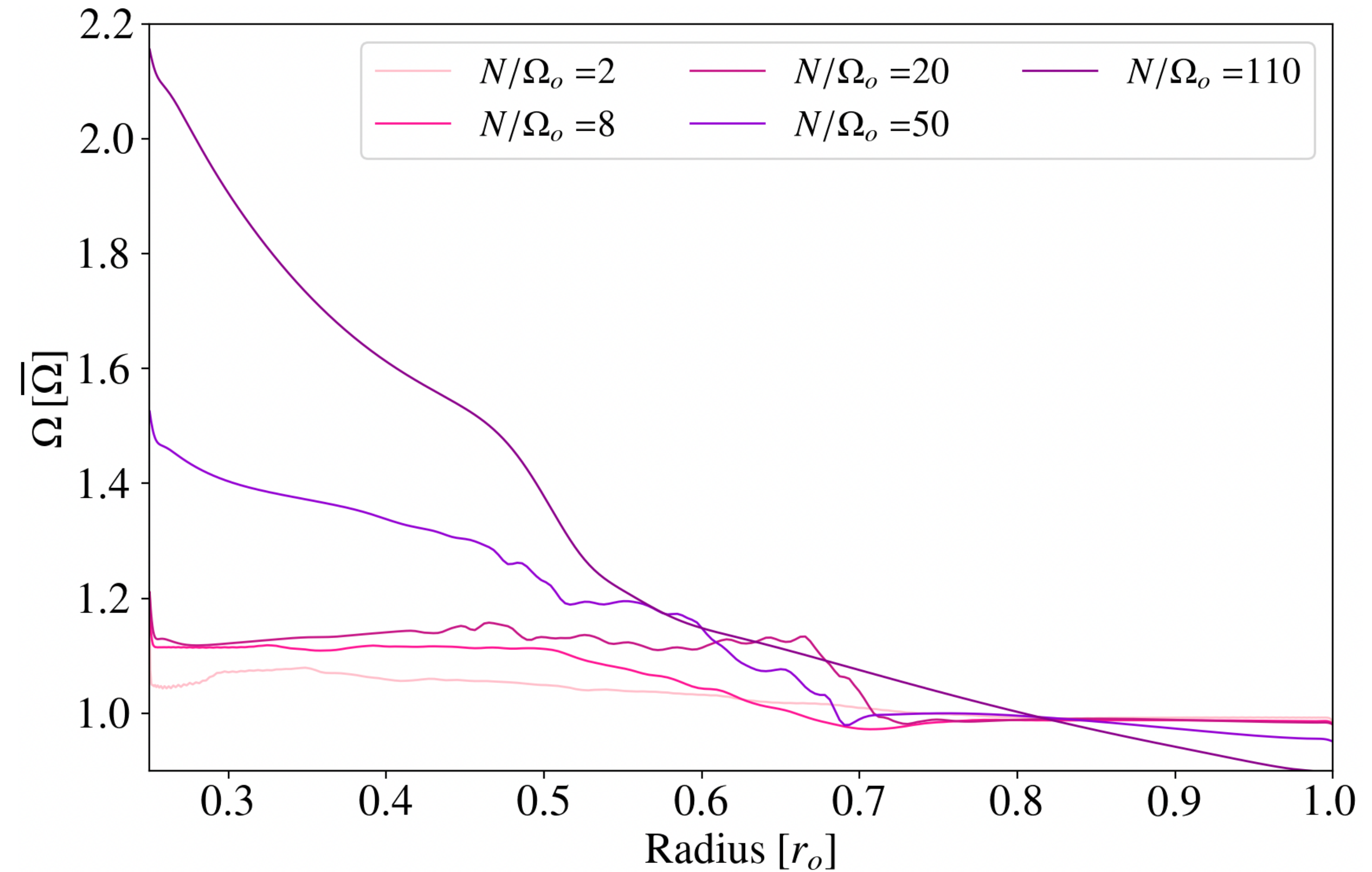
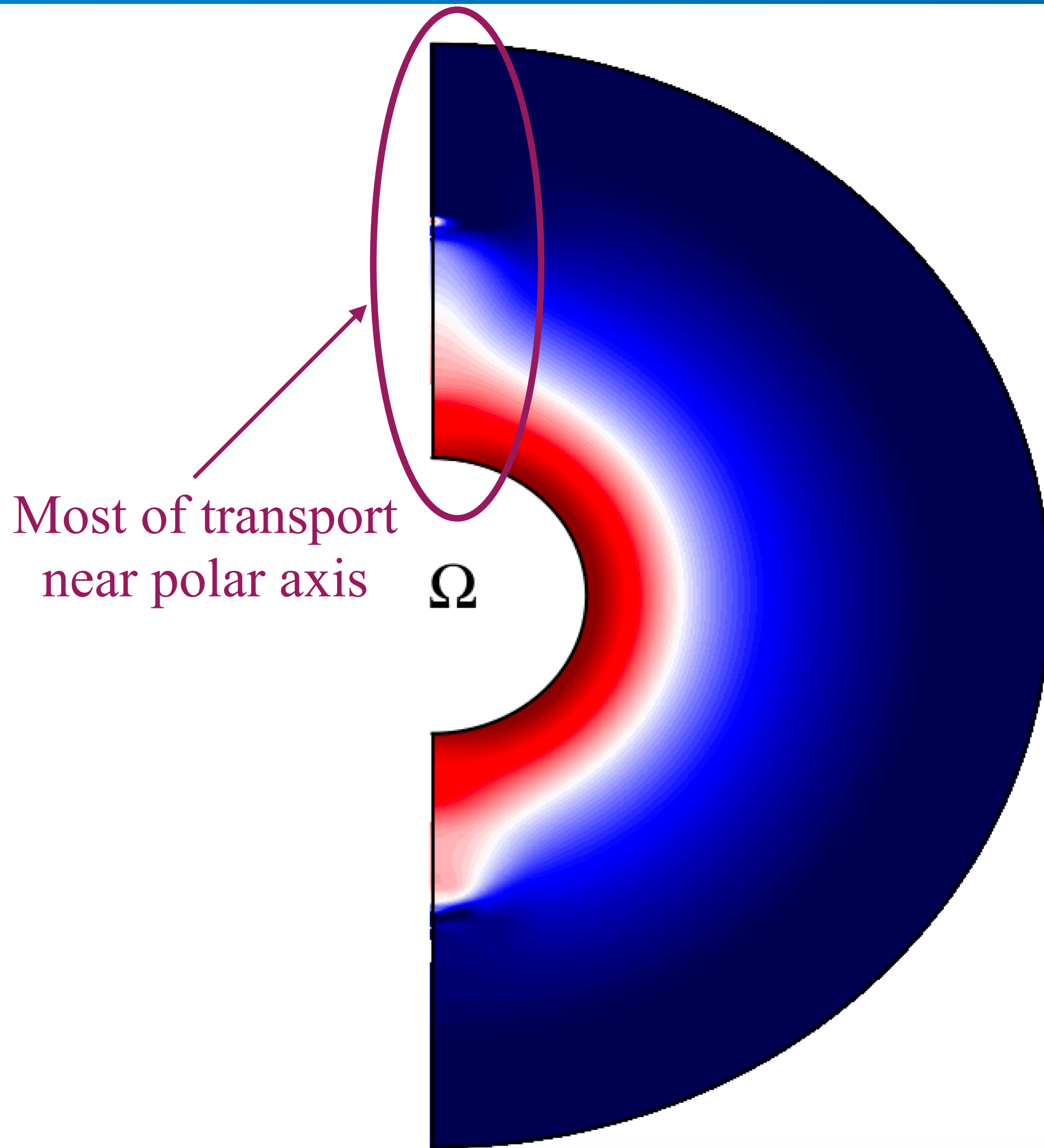
$$f_{\text{Ohm}} \equiv Q^\eta / (Q^\nu + Q^\eta)$$

with $Q^\eta \propto (\nabla \times \mathbf{B})^2$

and $Q^\nu \propto (\nabla \times \mathbf{v})^2$

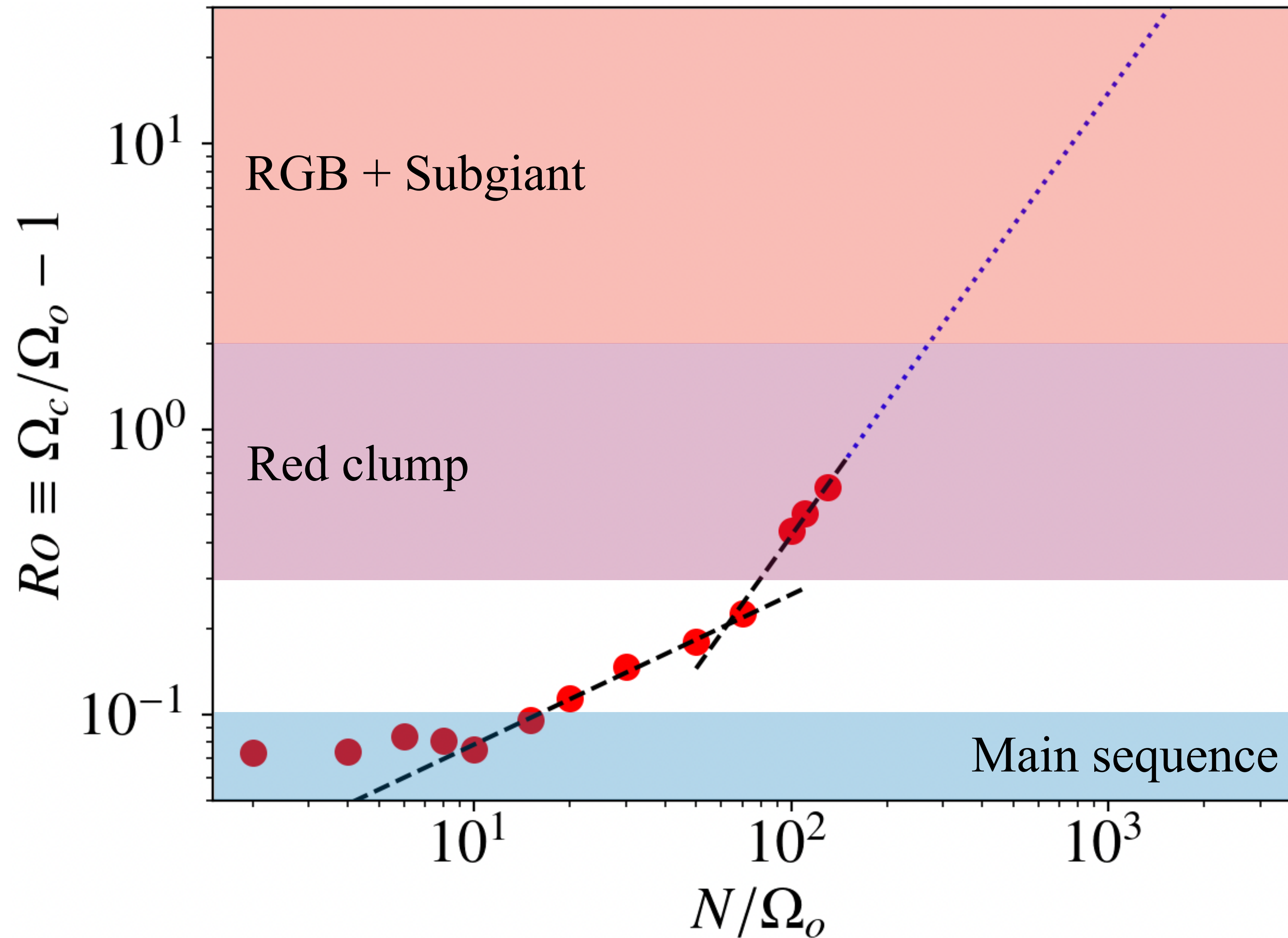


Local vs. global analysis



**3D nature of the dynamo is important
and complicates the application to stellar evolution models**

Ratio core/env. rotation rate



Implementing realistic diffusivity radial dependence

