

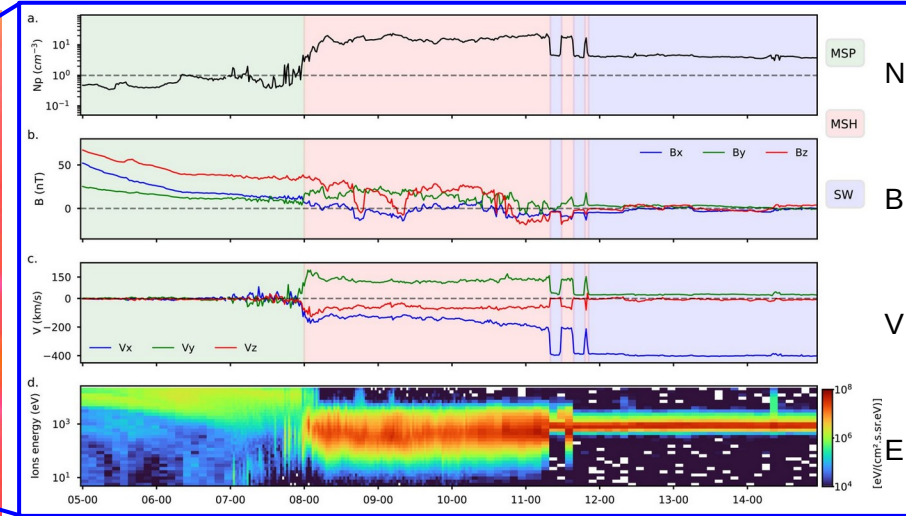
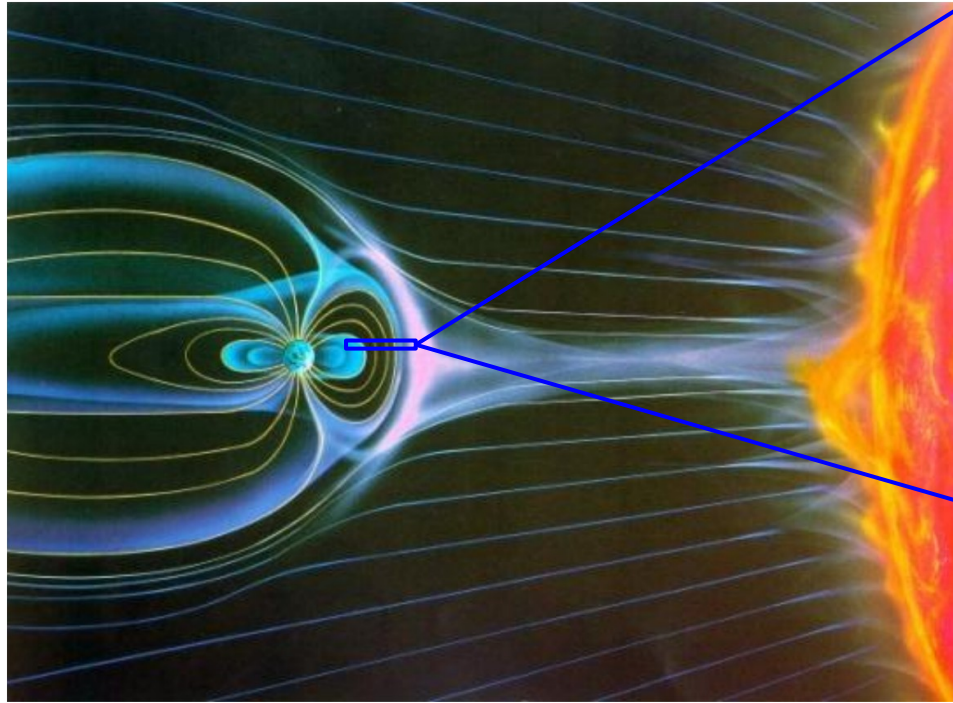
PHARE: Towards mesh and model refinement for space plasmas

Ulysse CAROMEL

Nicolas AUNAI, Philip DEEGAN, Roch SMETS, Ivan GIRAULT



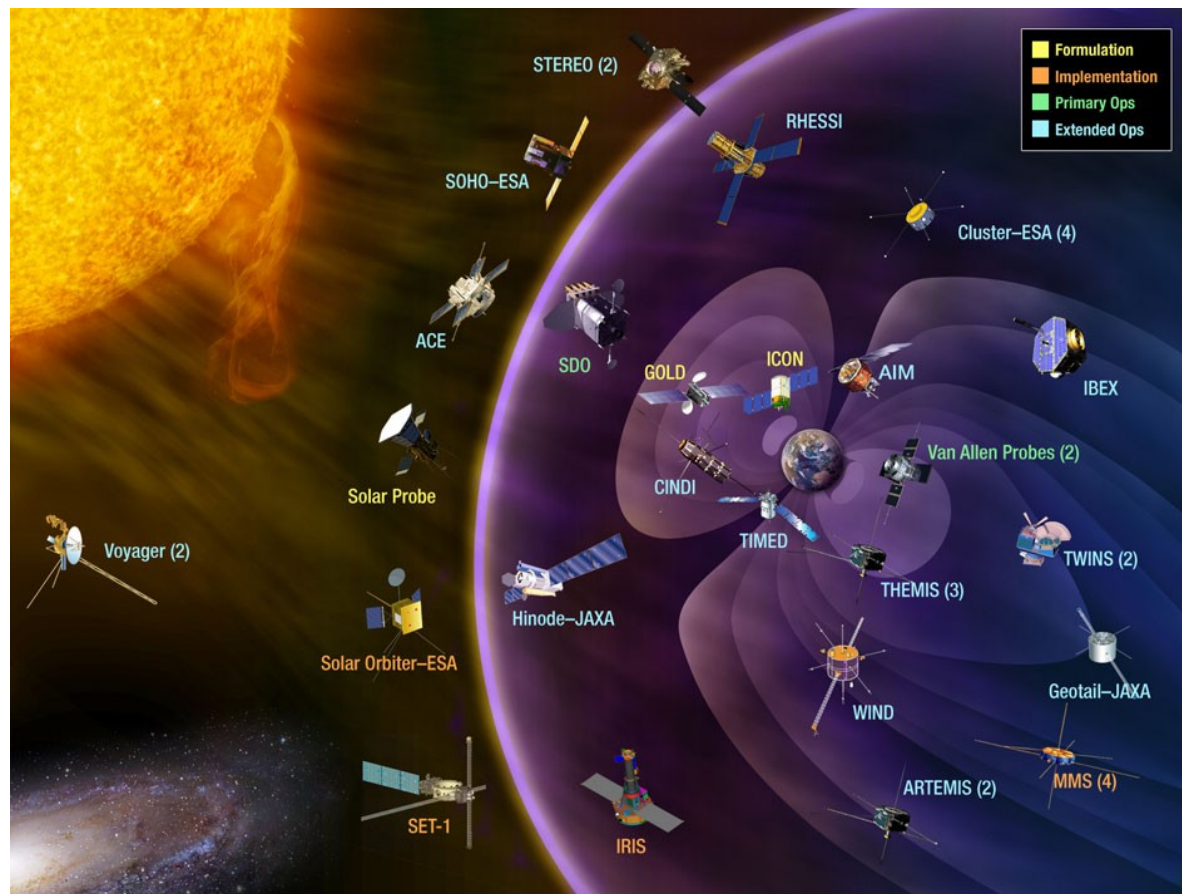
High resolution but localised



[Michotte de Welle et al. 2022]

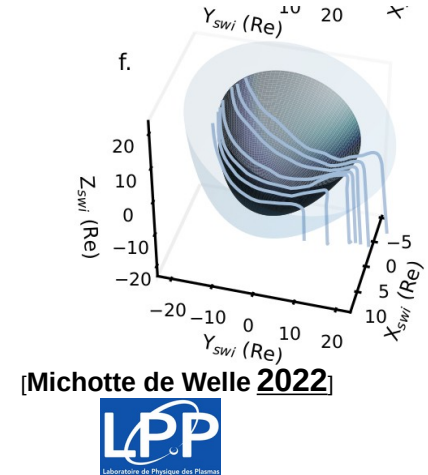
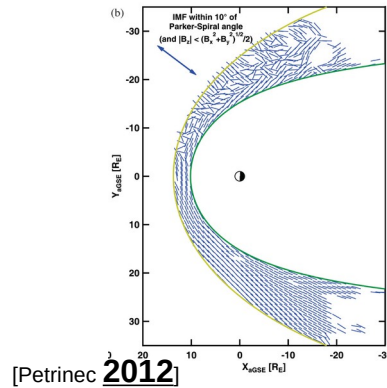
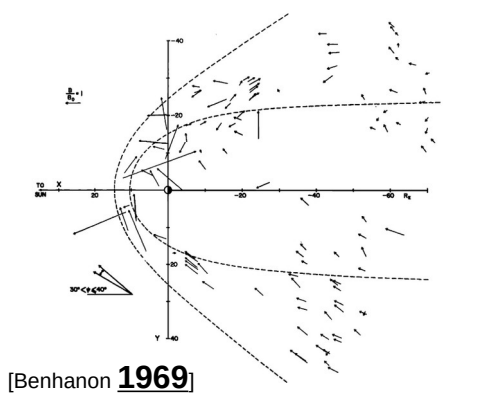
High space/time resolution in but no global vision.

We have decades of good quality data



Statistics and machine learning: Global but not dynamic

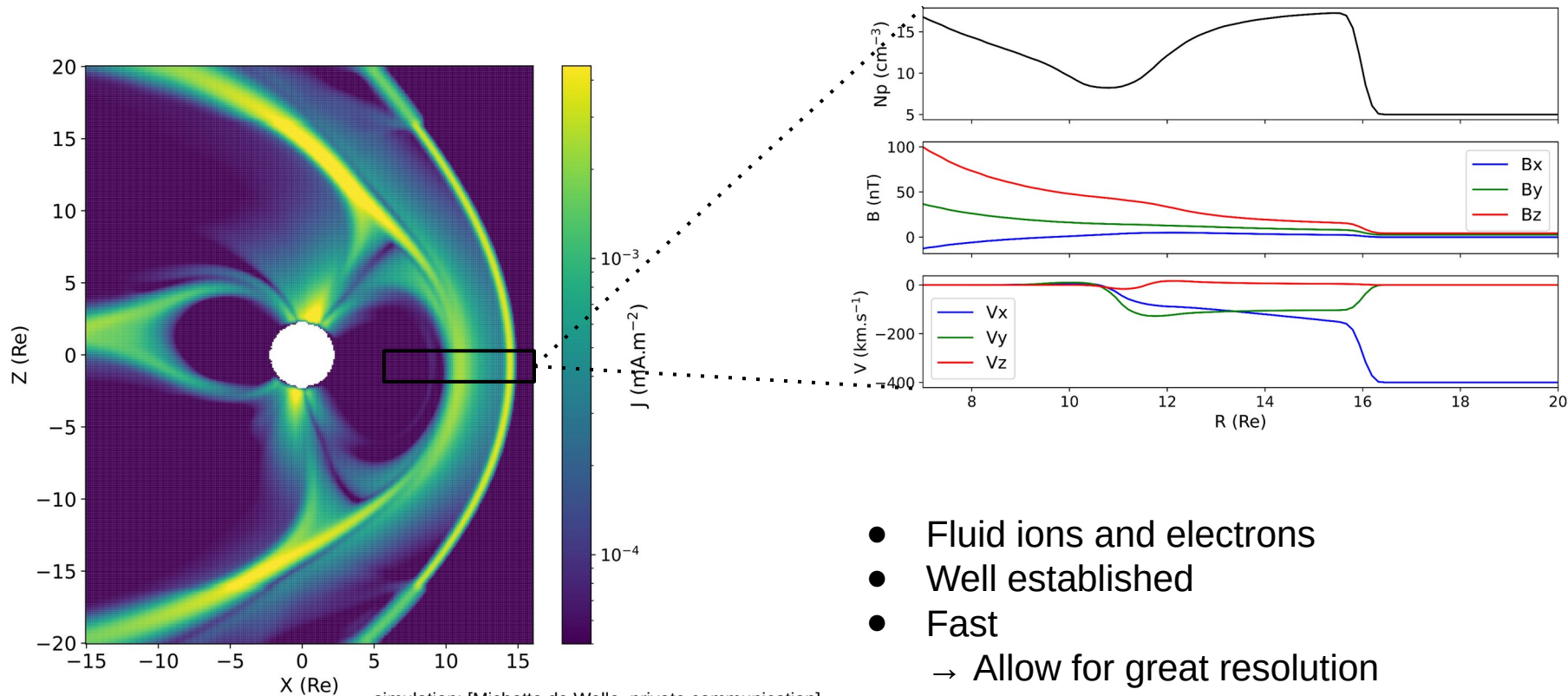
Recent efforts in machine learning has allowed scaling to much larger number of point



Averaged: we lose all small (temporal and spatial) scale processes that drive the global dynamics.

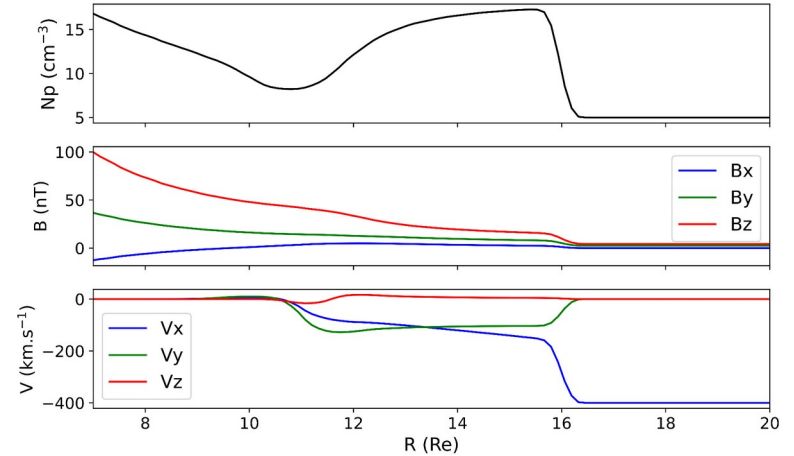
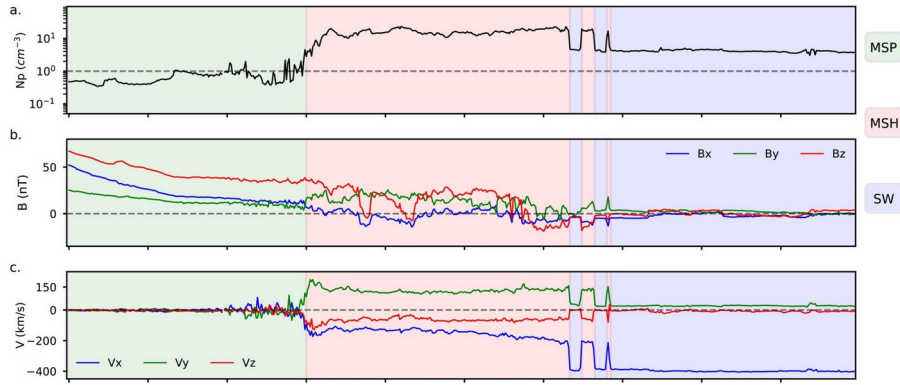
We want simulations to have both the global scale and the dynamics.

State of the art of global simulations: MHD



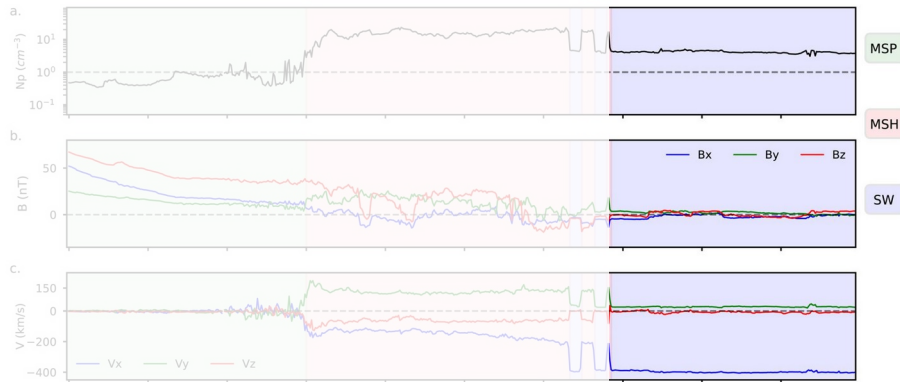
- Fluid ions and electrons
 - Well established
 - Fast
- Allow for great resolution

MHD capture regions and their interfaces

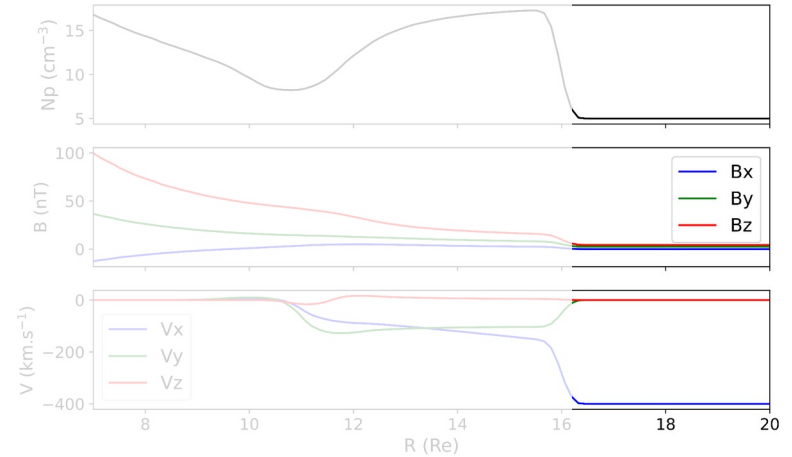


MHD capture regions and their interfaces

Solar wind

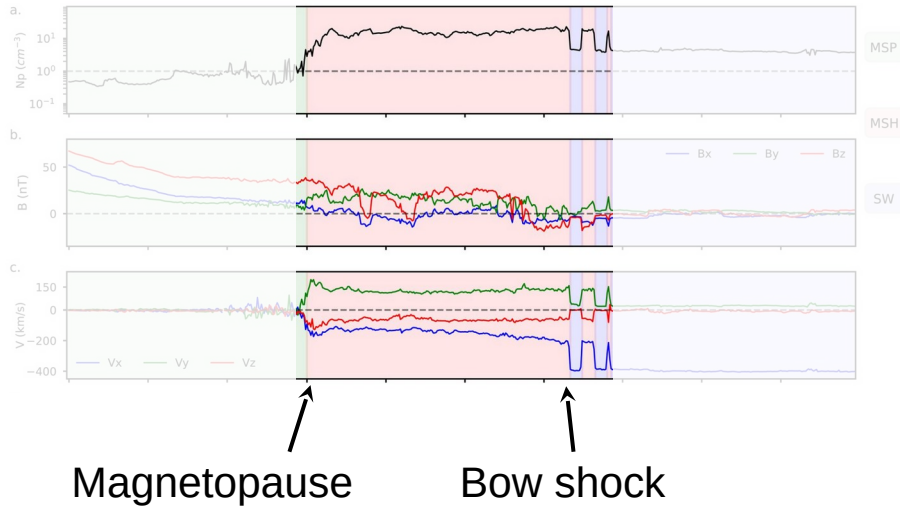


Solar wind

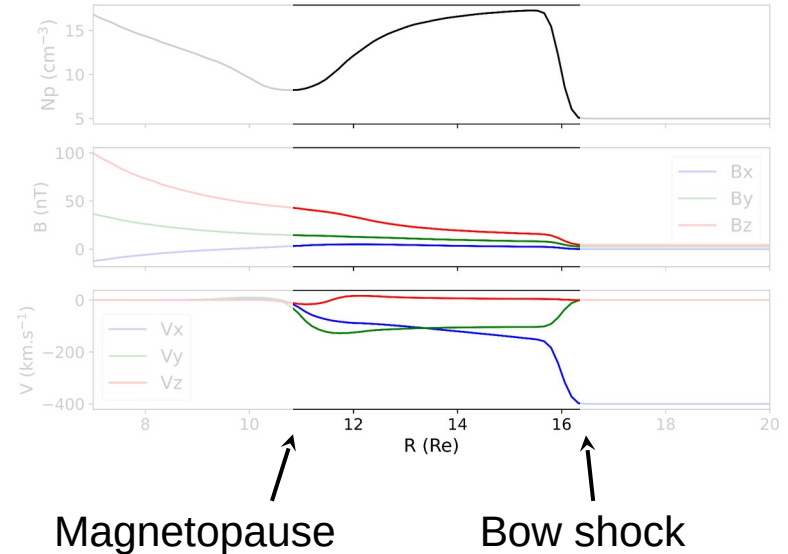


MHD capture regions and their interfaces

Magnetosheath

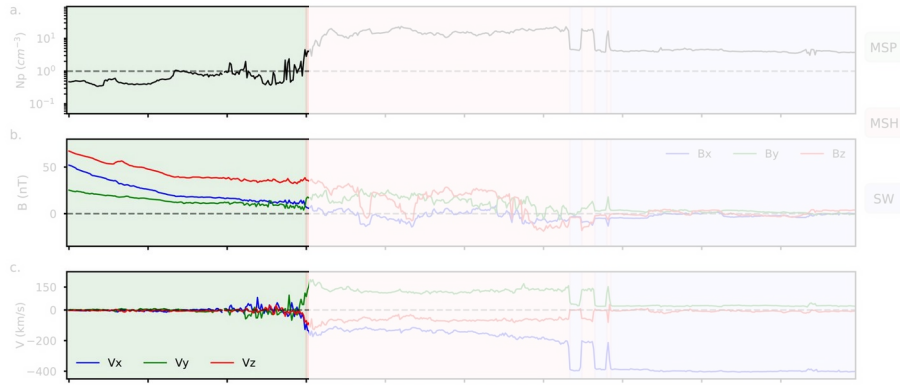


Magnetosheath



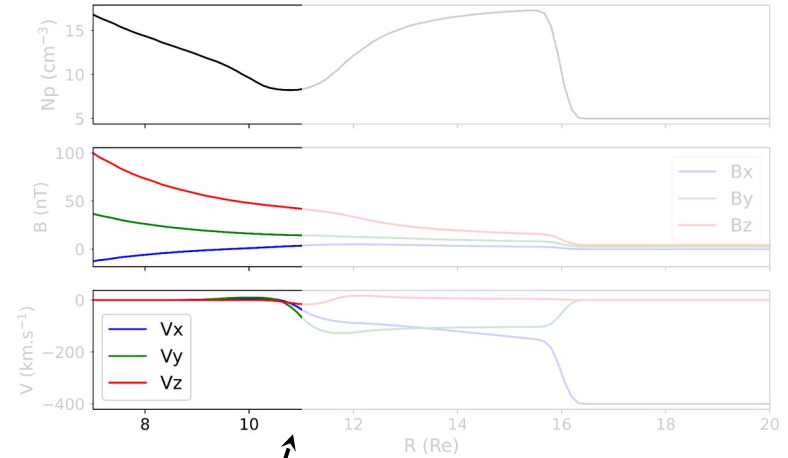
MHD capture regions and their interfaces

Magnetosphere



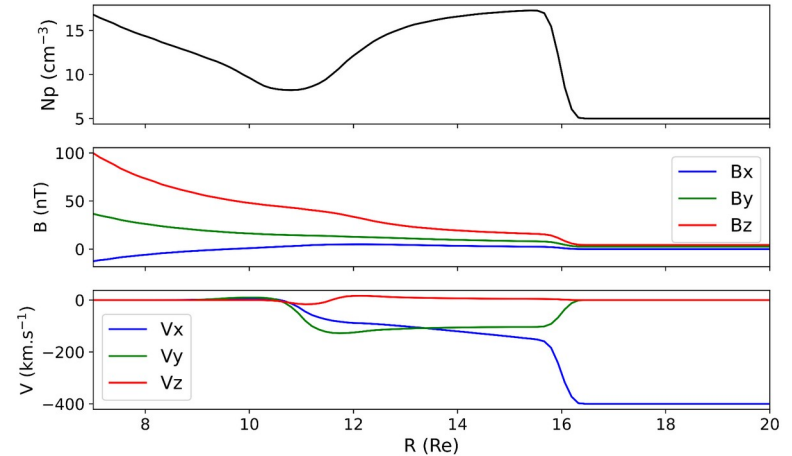
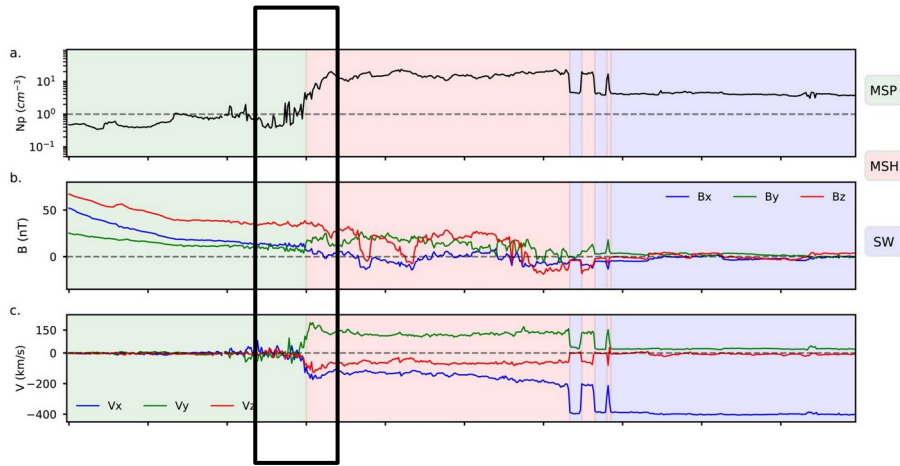
Magnetopause

Magnetosphere



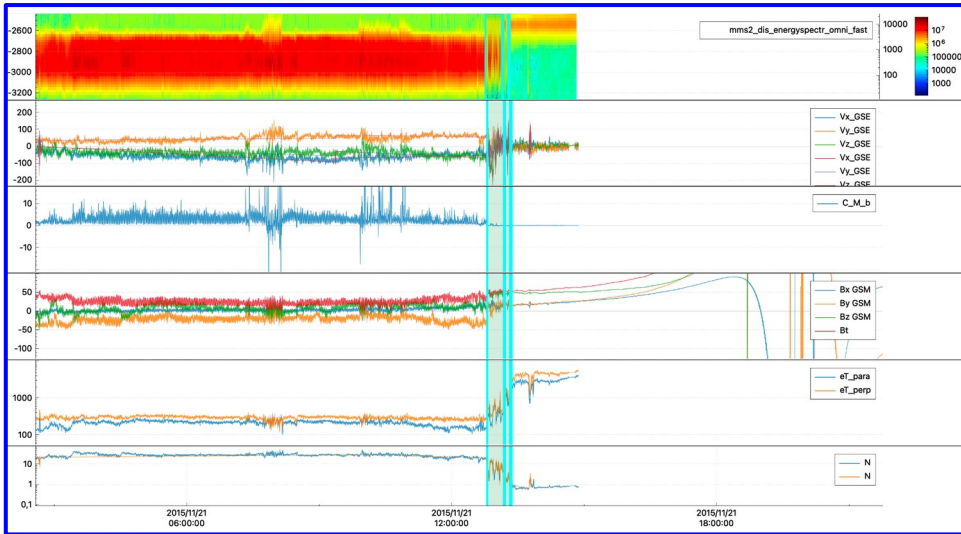
Magnetopause

But misses small-scale processes driving their coupling



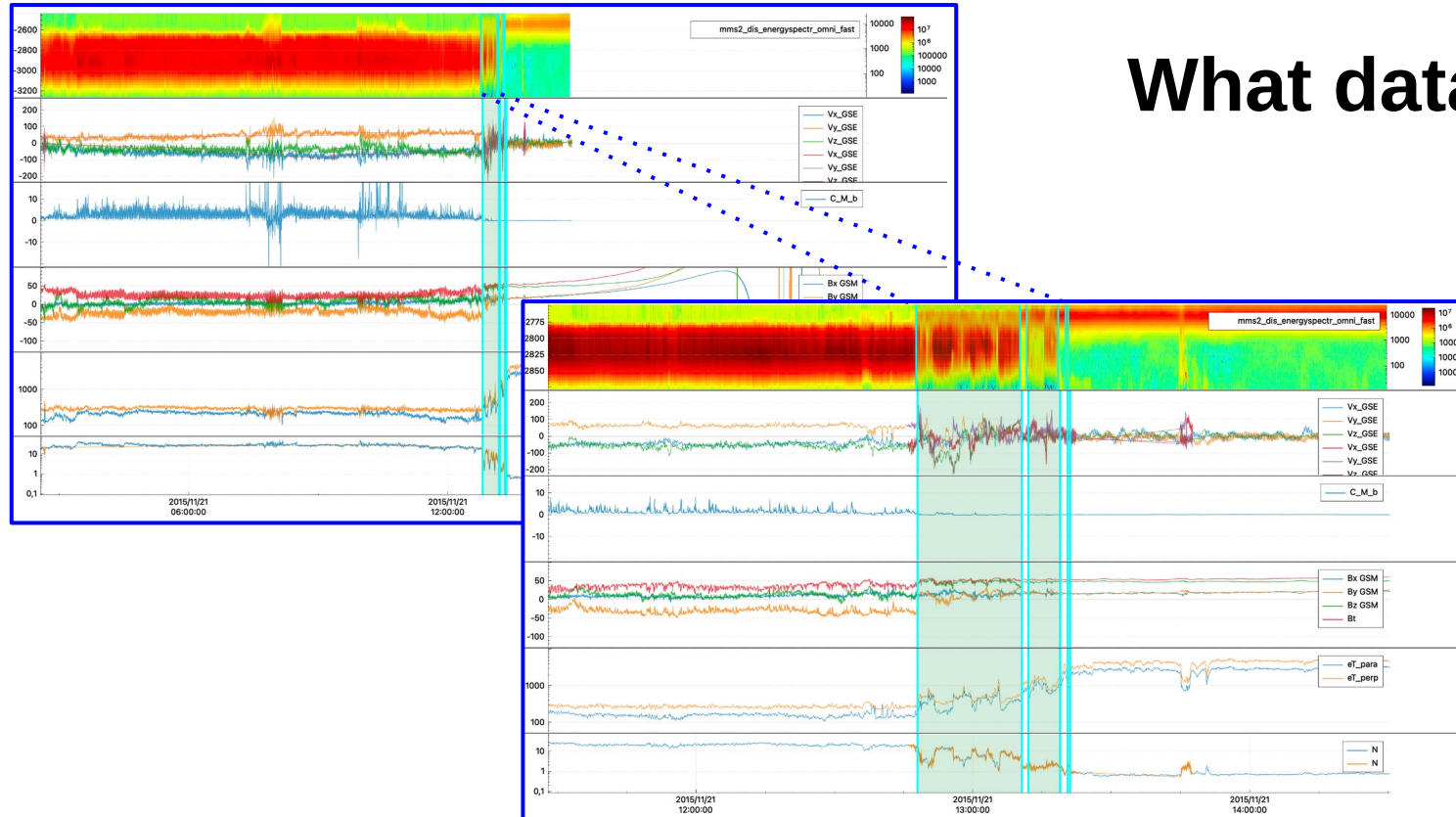
What happens in the active boundaries between the different environments ?

What data shows

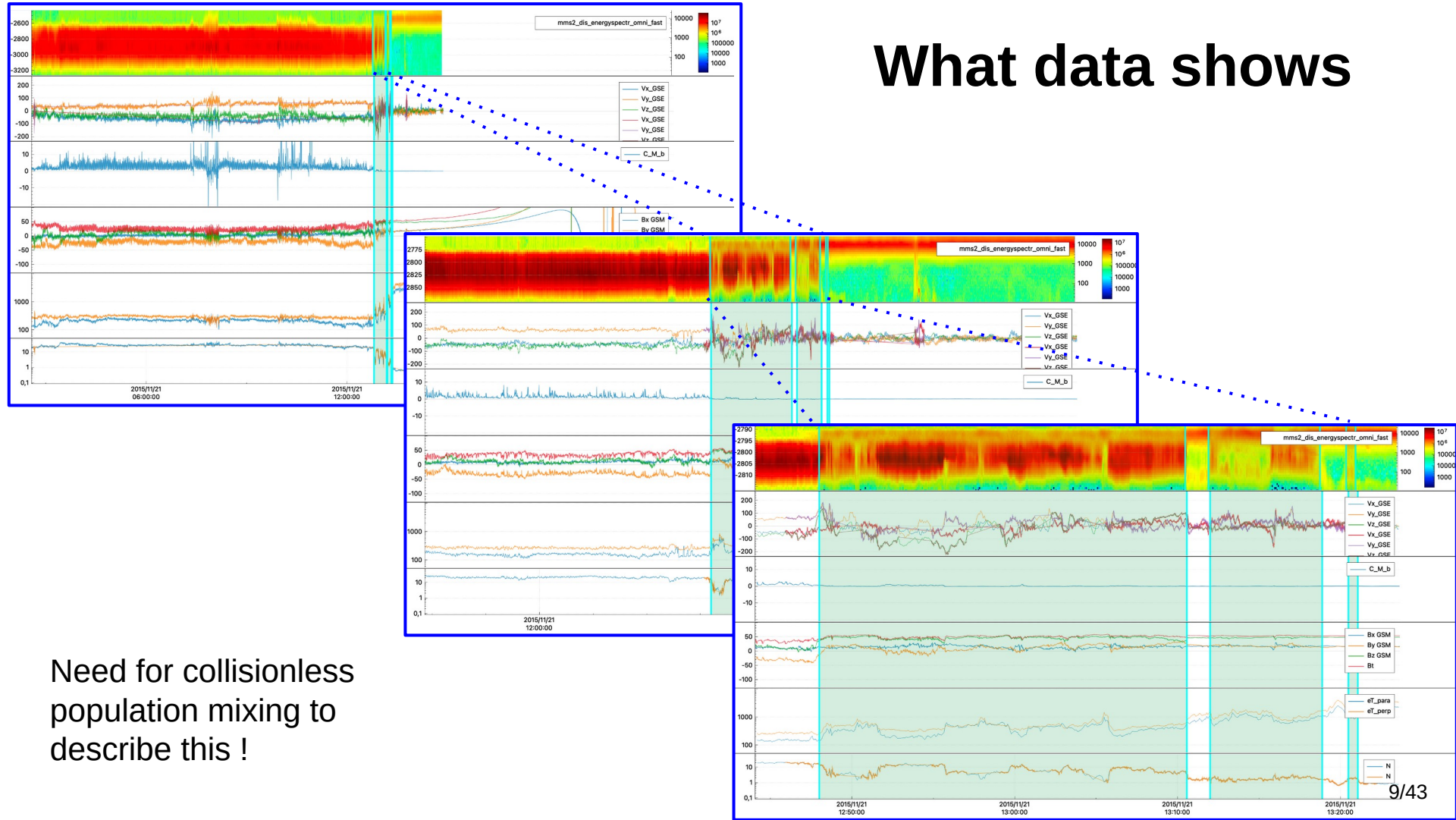


[SciQLop software: Jeandet et al. (2025)]
<https://doi.org/10.5281/zenodo.17176824>

What data shows

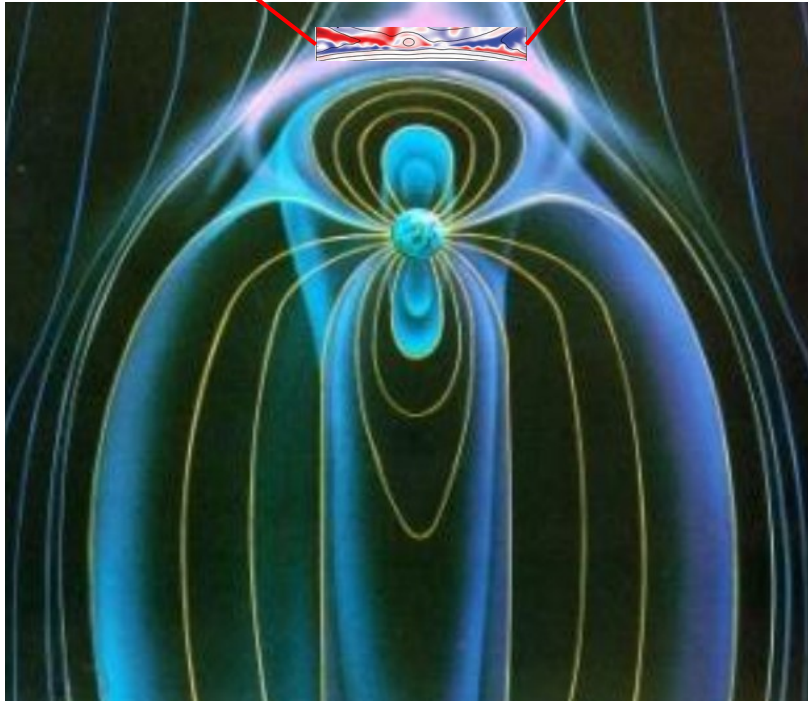
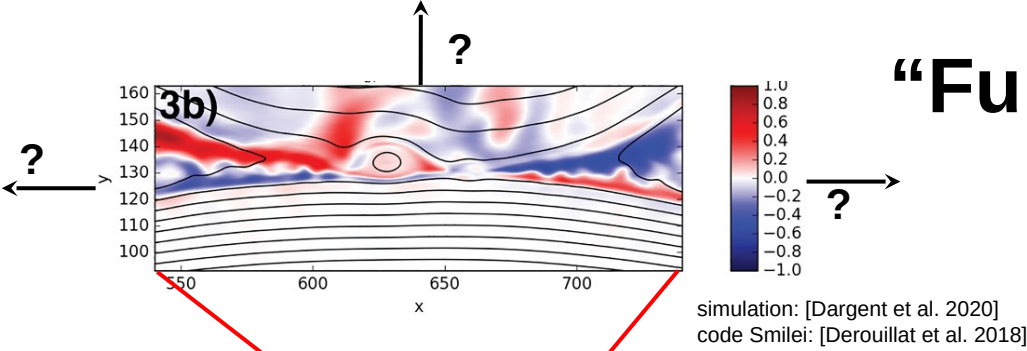


What data shows



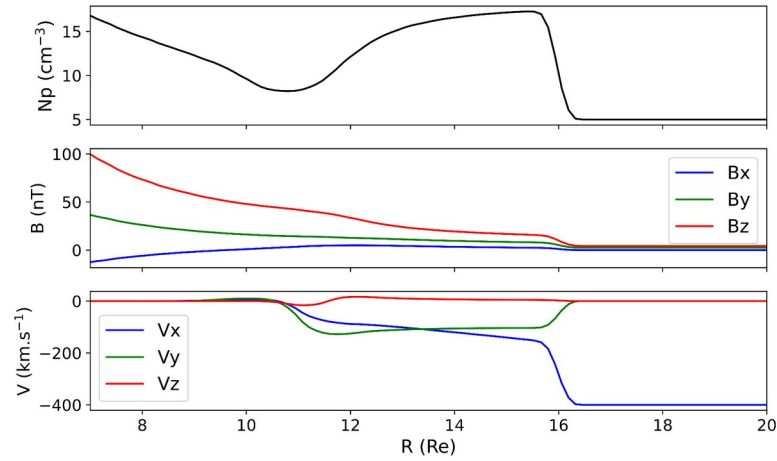
Need for collisionless
population mixing to
describe this !

“Full” PIC: no global scale



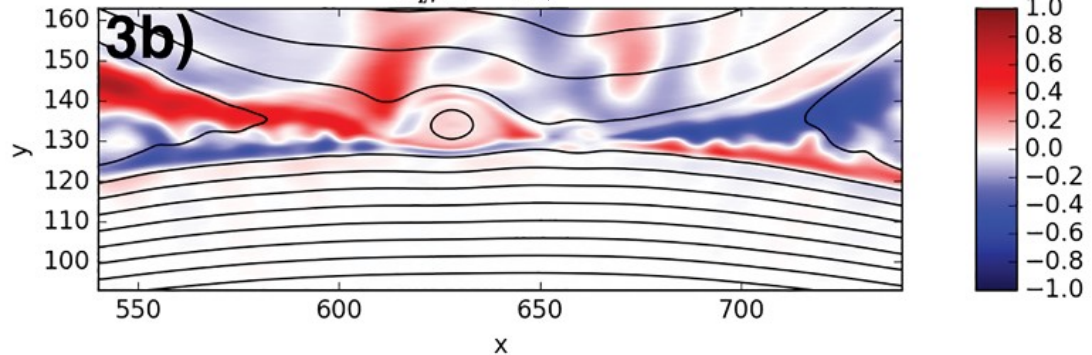
- Kinetic ions and electrons
- Realistic physics
- Very computationally heavy
 - Restrained to a small box
 - Global scale is not attainable

How can we do better ?



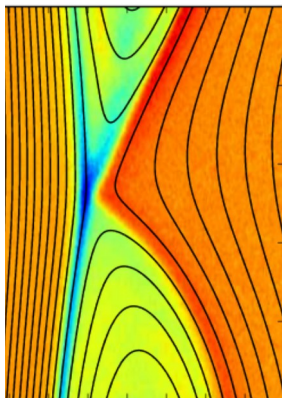
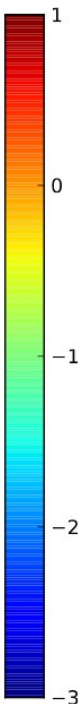
No microphysics !

No global scales !

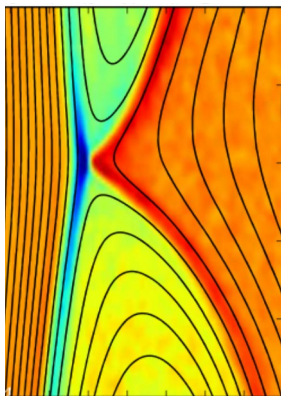


Hybrid PIC: the good compromise ?

J_z



Full PIC



Hybrid PIC

- Kinetic ions, fluid electrons
- Physically realistic if well resolved

→ Global simulations attainable ?

Hard to have both small and large scales

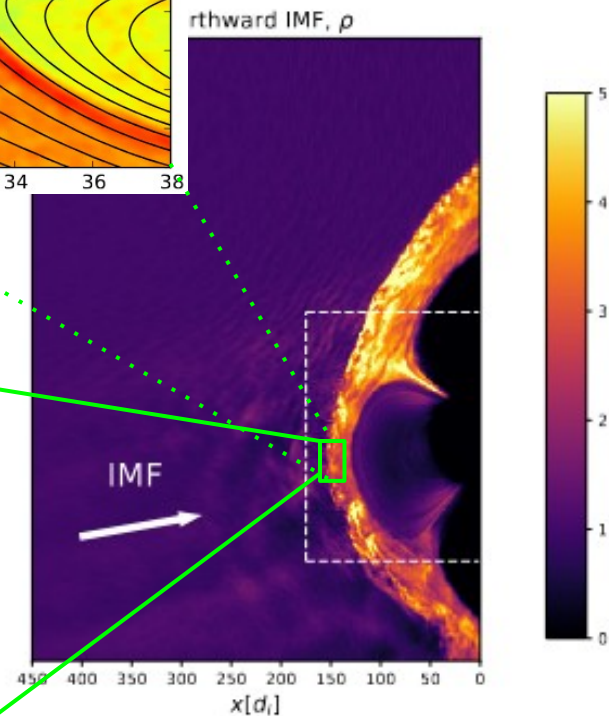
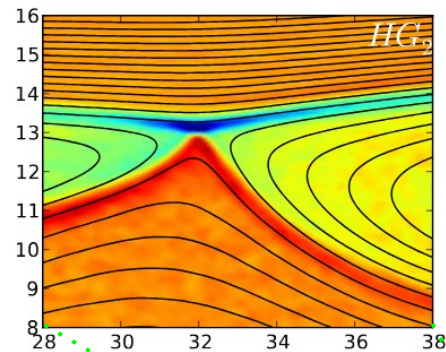
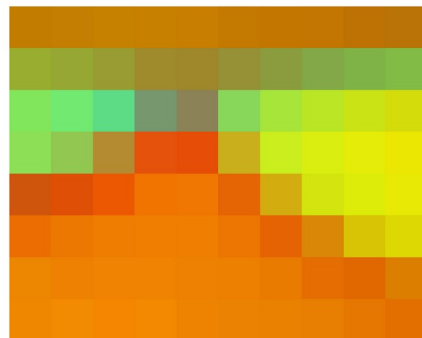
no large scale



$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \times \left(\boxed{-\mathbf{v}_i \times \mathbf{B} + \eta \mathbf{j}} + \boxed{\frac{\mathbf{j} \times \mathbf{B} - \nabla P_e}{en}} \boxed{-\nu \nabla^2 \mathbf{j}} \right) = 0$$



unresolved small scales



We can't just “use bigger computers”

Resolution in 3d hybrid
simulations:

Goal:

$$1 \text{ to } 5\delta_i \longrightarrow 0.05\delta_i \quad 100$$

Per dimension:

$$100^3 = 10^6$$

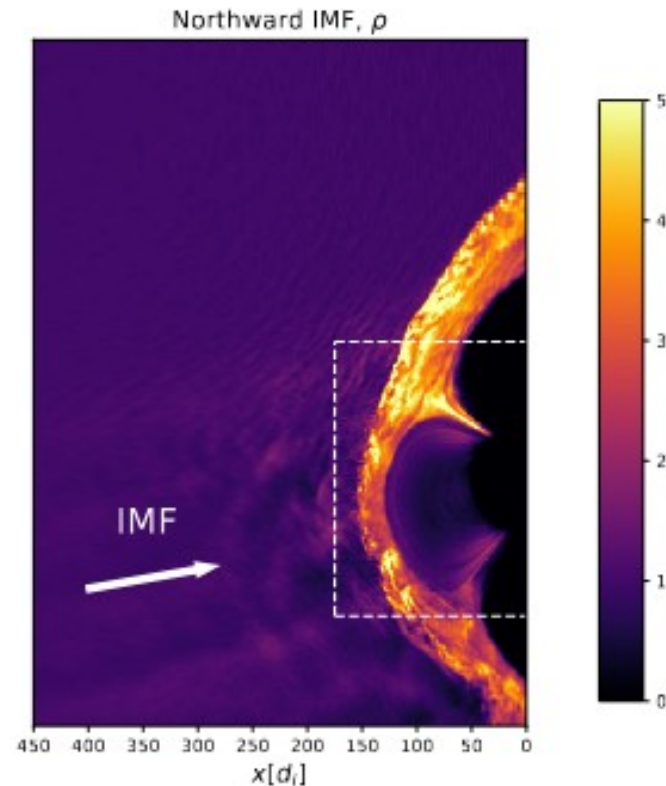
\times

Time step: $\Delta t \sim \Delta x^2$

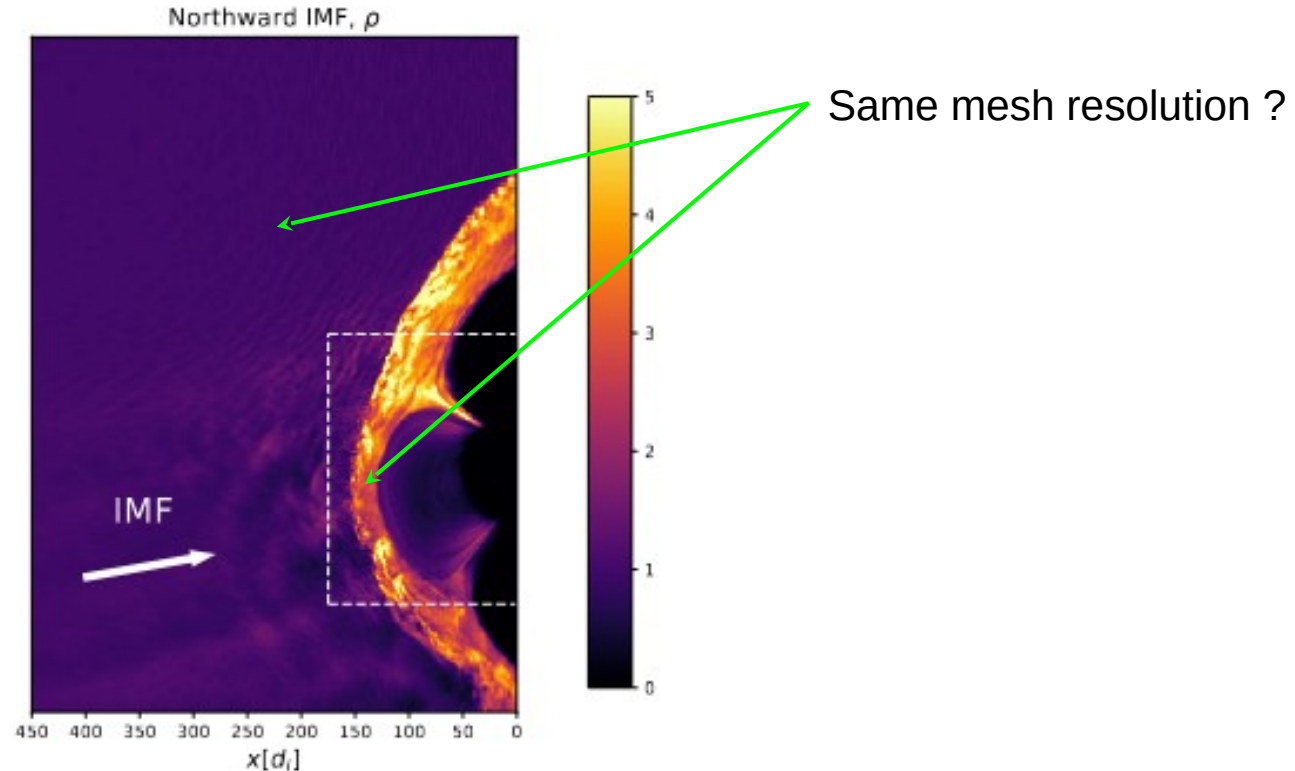
$$100^2 = 10^4$$

$$100^3 \times 100^2 = 10^{10}$$

We would need 10^{10} times more compute
power !

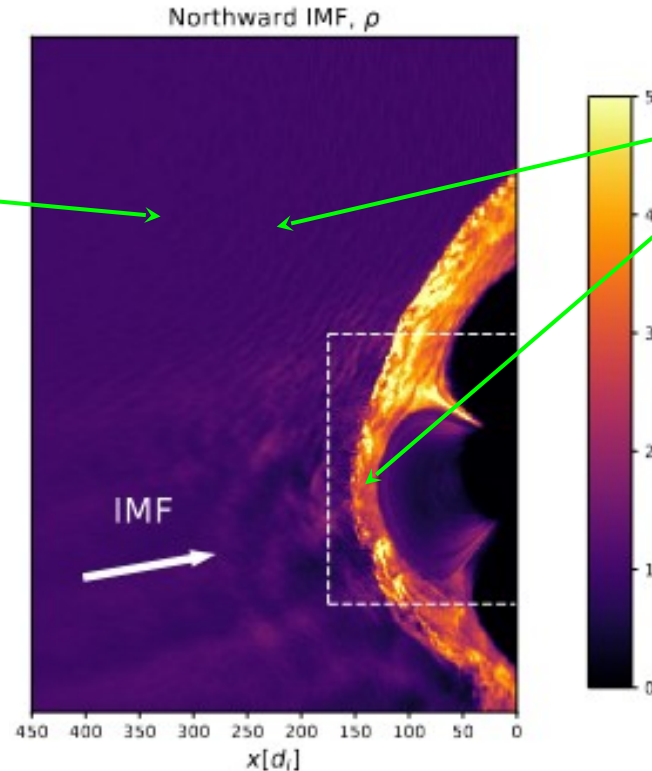


We need a different approach



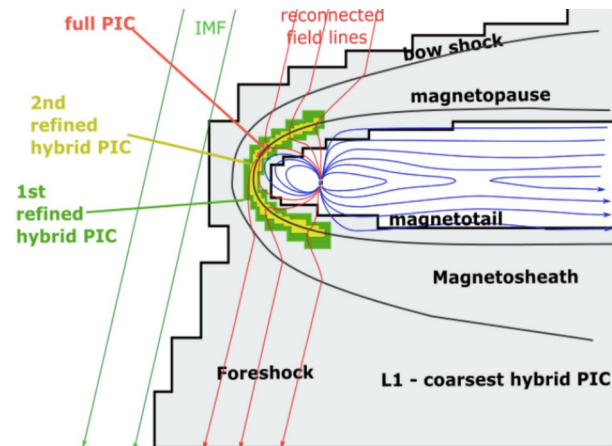
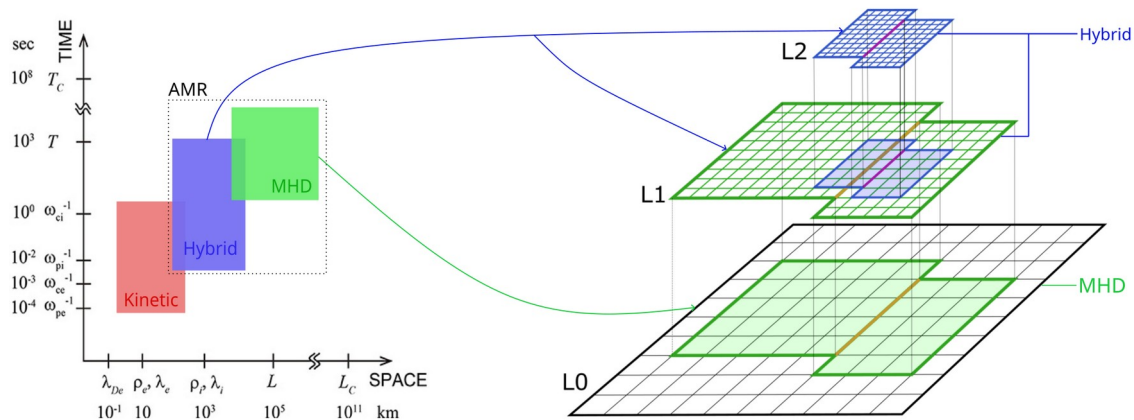
We need a different approach

Do we need hybrid there ?



Same mesh resolution ?

Our goal in PHARE: adaptive mesh and model refinement (AM2R)



PHARE: an open-source, modern C++ plasma simulation code developed at LPP **for the community.**

A code for the community



Open source



High perf,
abstractions



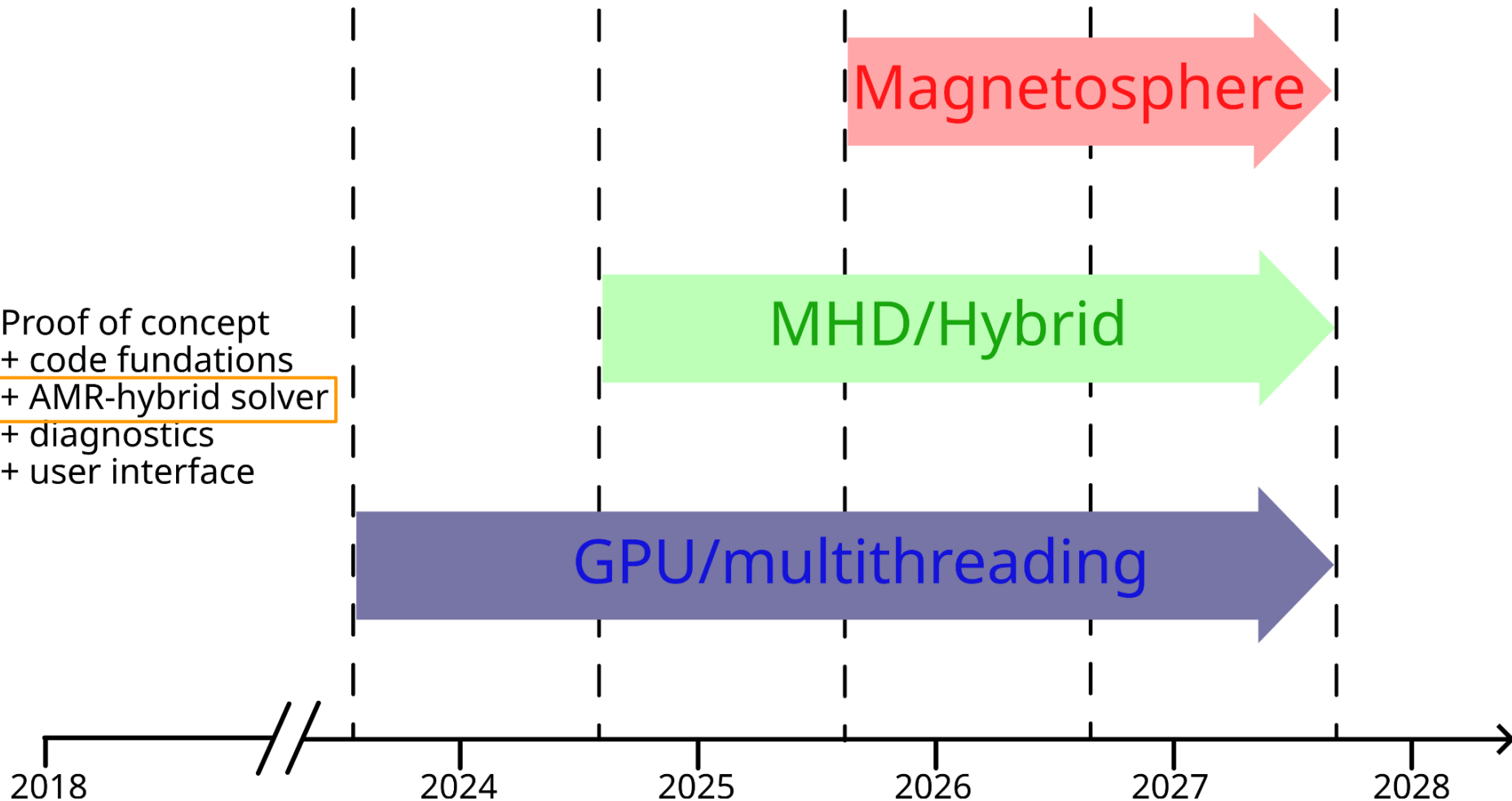
User friendly

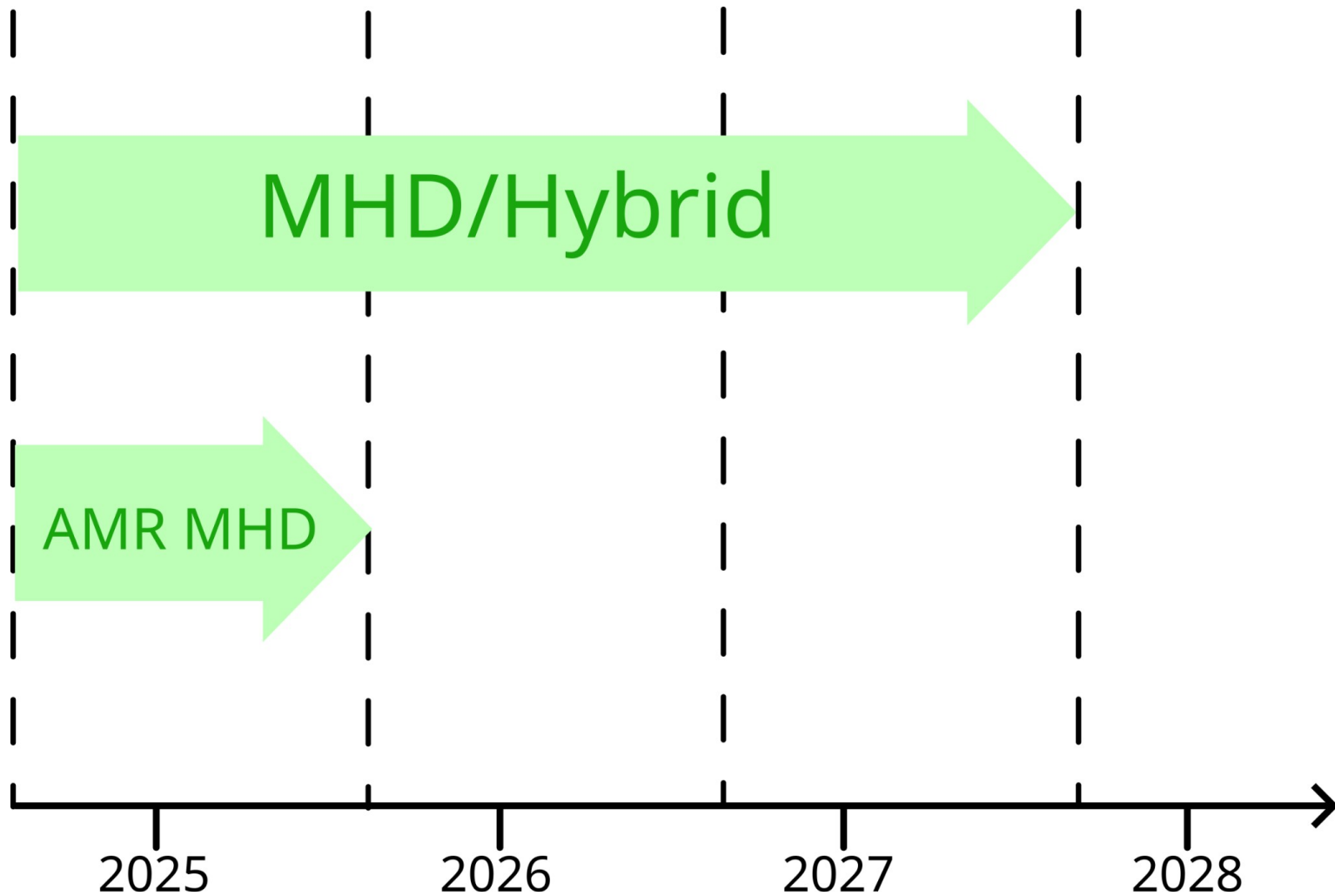


Testability,
Robustness

- 100% open source.
- Modern c++ (c++20) for performance and abstraction while maintaining friendly python user interface.
- Extensive testing running on continuous integrations.
- Relies on well established AMR library called SAMRAI.

Proof of concept
+ code foundations
+ AMR-hybrid solver
+ diagnostics
+ user interface





Step 1: choosing the equations

We advance the following equations:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

Continuity equation

$$\frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot [\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + P^*] = 0$$

Momentum equation

$$\frac{\partial \mathcal{E}}{\partial t} + \nabla \cdot \left[(\mathcal{E} + P^*) \mathbf{v} + \frac{\mathbf{E} \times \mathbf{B}}{\mu_0} \right] = 0$$

Energy equation

with:

$$P^* = P + \frac{\mathbf{B} \cdot \mathbf{B}}{2}$$

The total pressure

Step 1: choosing the equations

We advance the following equations:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 \quad \text{Continuity equation}$$

$$\frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot [\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + P^*] = 0 \quad \text{Momentum equation}$$

$$\frac{\partial \mathcal{E}}{\partial t} + \nabla \cdot \left[(\mathcal{E} + P^*) \mathbf{v} + \frac{\mathbf{E} \times \mathbf{B}}{\mu_0} \right] = 0 \quad \text{Energy equation}$$

Polytropic closure:

$$\mathcal{E} = \frac{P}{\gamma - 1} + \frac{1}{2}(\rho(\mathbf{v} \cdot \mathbf{v}) + \frac{\mathbf{B} \cdot \mathbf{B}}{\mu_0})$$

Step 1: choosing the equations

We advance the following equations:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

Continuity equation

$$\frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot [\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + P^*] = 0$$

Momentum equation

$$\frac{\partial \mathcal{E}}{\partial t} + \nabla \cdot \left[(\mathcal{E} + P^*) \mathbf{v} + \frac{\mathbf{E} \times \mathbf{B}}{\mu_0} \right] = 0$$

Energy equation

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} = 0$$

Faraday's law

with:

$$\mathbf{E} = -\mathbf{v} \times \mathbf{B} \overset{\text{Hall}}{+ \frac{\mathbf{j} \times \mathbf{B}}{en_e}} \overset{\text{Ohmic resistivity}}{+ \eta \mathbf{j}} \overset{\text{Hyper-resistivity}}{- \nu \nabla^2 \mathbf{j}}$$

Ohm's law

$$\mathbf{j} = \frac{\nabla \times \mathbf{B}}{\mu_0}$$

Ampere's law

Step 1: choosing the equations

We advance the following equations:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

Continuity equation

$$\frac{\partial (\rho \mathbf{v})}{\partial t} + \nabla \cdot [\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + P^*] = 0$$

Momentum equation

$$\frac{\partial \mathcal{E}}{\partial t} + \nabla \cdot \left[(\mathcal{E} + P^*) \mathbf{v} + \frac{\mathbf{E} \times \mathbf{B}}{\mu_0} \right] = 0$$

Energy equation

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} = 0$$

Faraday's law

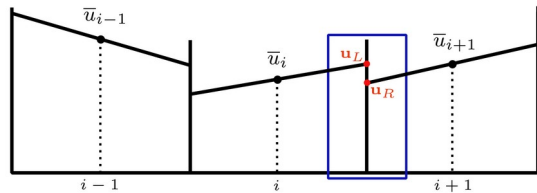
Under the constraint:

$$\nabla \cdot \mathbf{B} = 0$$

Step 2: Numerical implementation

Numerical scheme: Godunov Finite Volumes

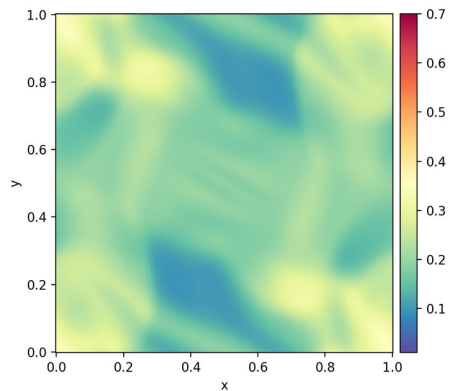
[Godunov
1959]



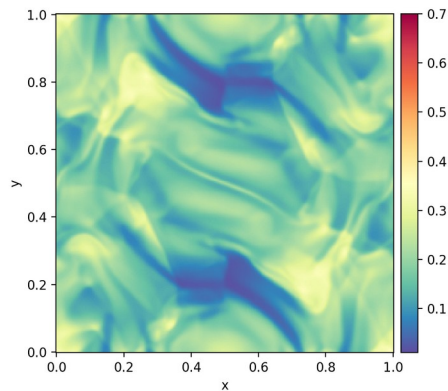
Reconstruction step

Pressure on Orszag-Tang vortex 256x256 at $t=1$.

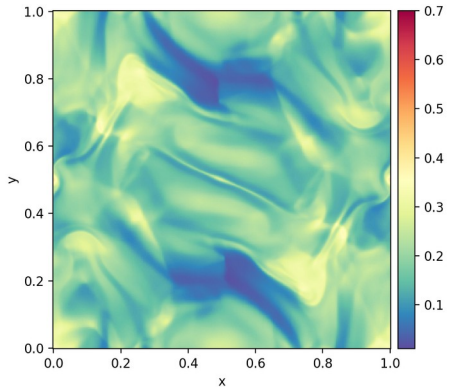
Constant



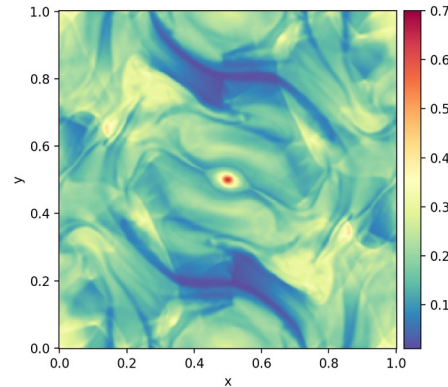
Linear



WENO3



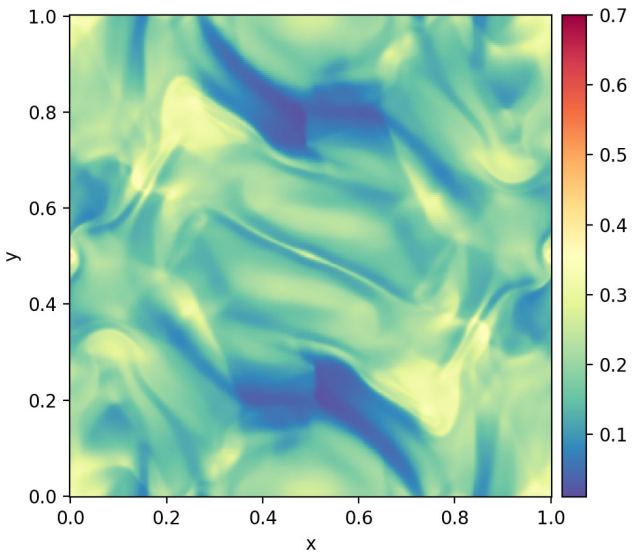
WENOZ



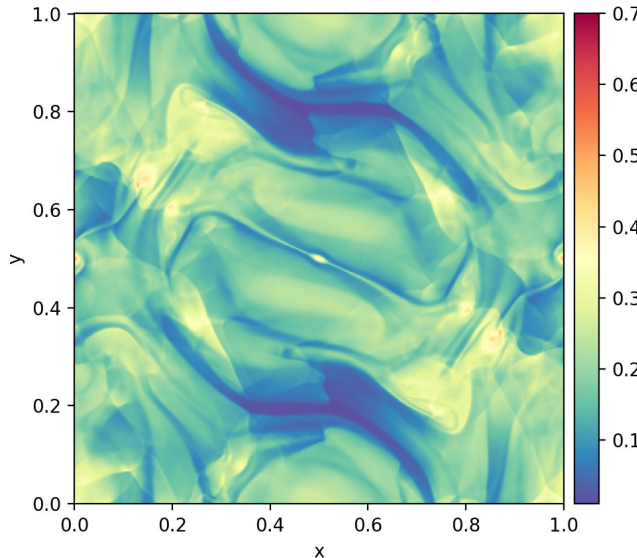
[Orszag & Tang 1979]

Functional Test: Orszag-Tang vortex

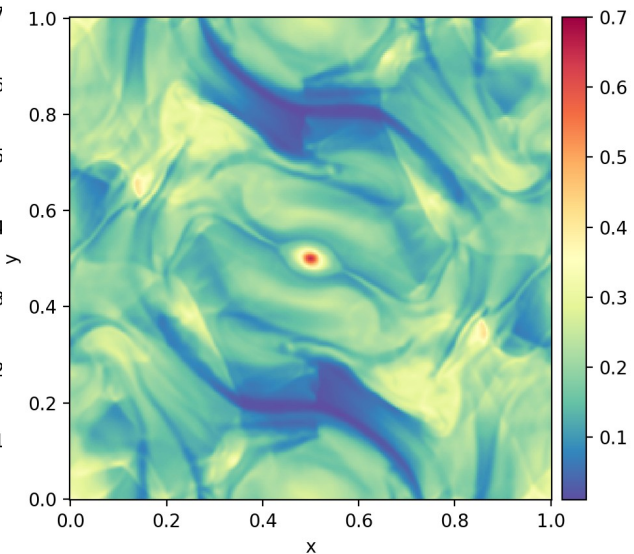
WENO3 256x256



WENO3 1024x1024



WENOZ 256x256

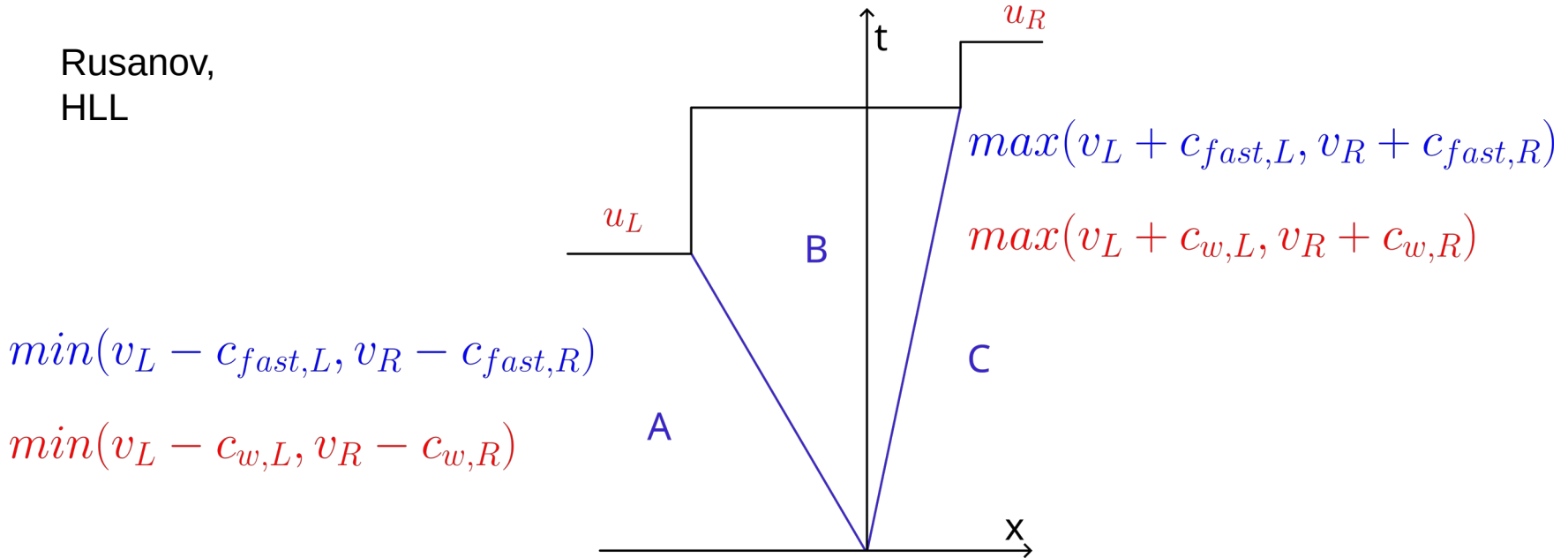


Riemann Solvers

Ideal MHD

Hall MHD

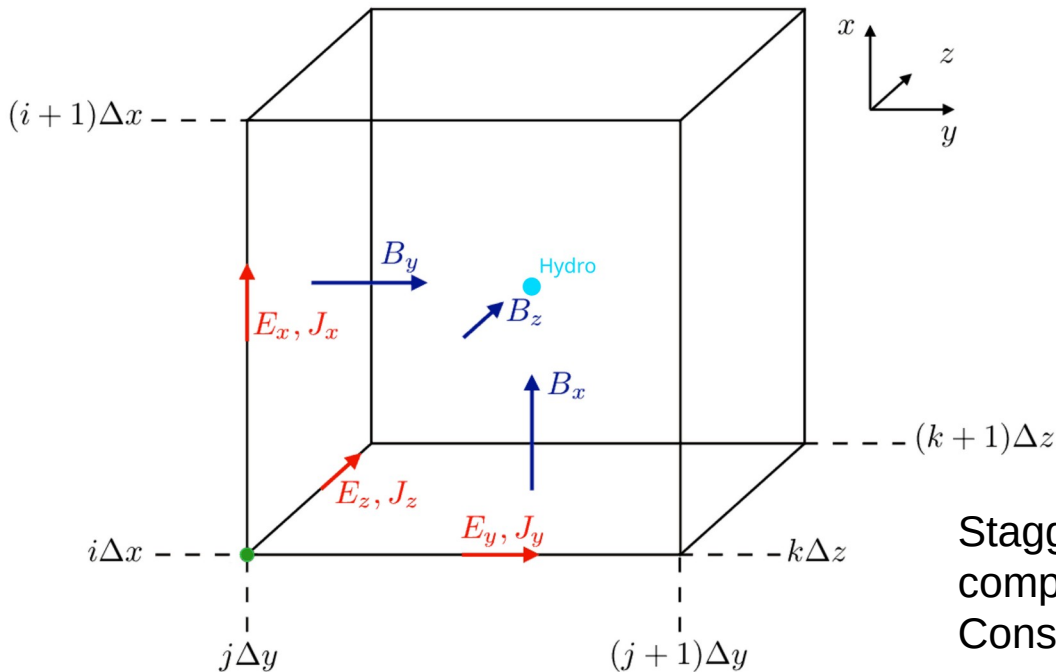
Rusanov,
HLL



c_{fast} = fast magnetosonic speed

c_w = whistler speed

Maintaining divergence free condition



[Yee 1966]

$$\nabla \cdot \mathbf{B} = 0$$

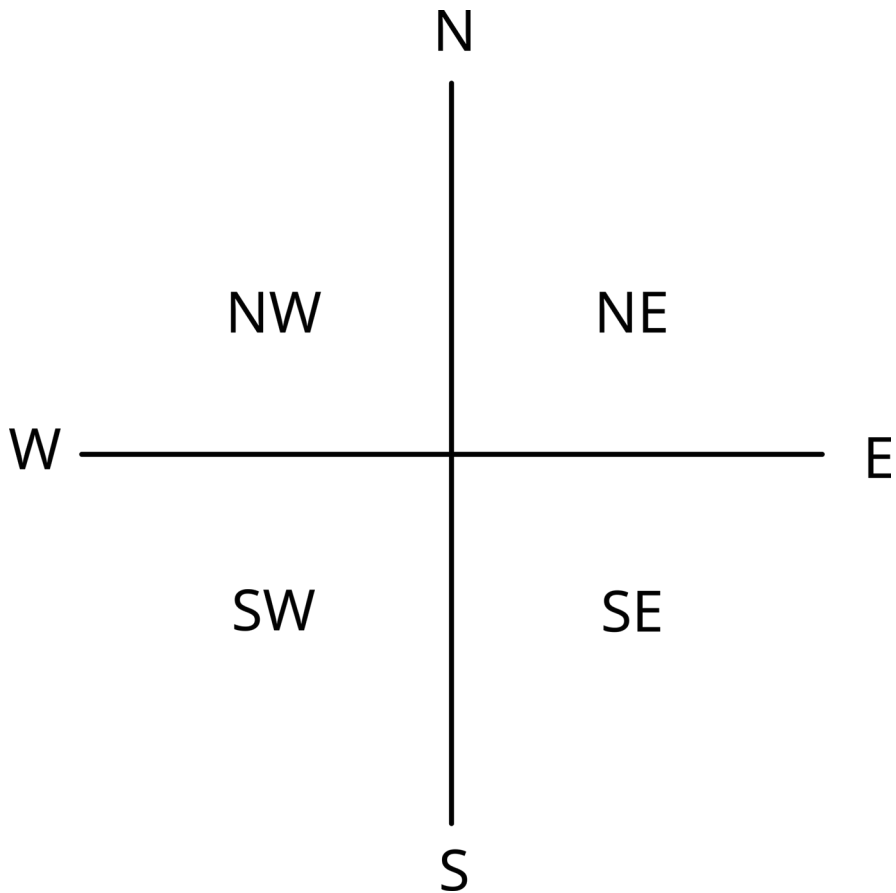
$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} = 0$$

Staggered grid for \mathbf{B} , \mathbf{E} , \mathbf{J} . Special care needed to compute \mathbf{E} , using a class of schemes called Constrained Transport (CT).

In PHARE: CT averaged, Upwind CT.

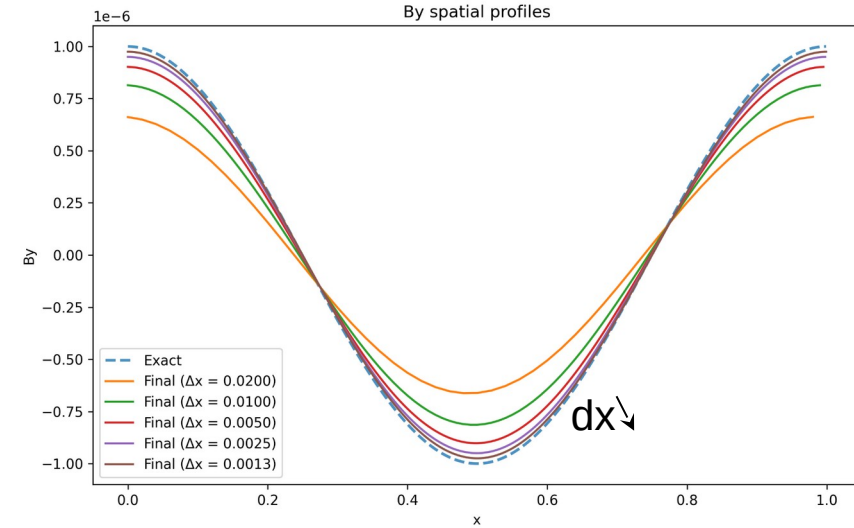
Upwind Constrained Transport (UCT)

[Mignone & Del Zanna
2021]



- Reconstruction on the edge directly
- 2d Riemann problem solve as opposed to arithmetic averaging of 1d riemann solvers

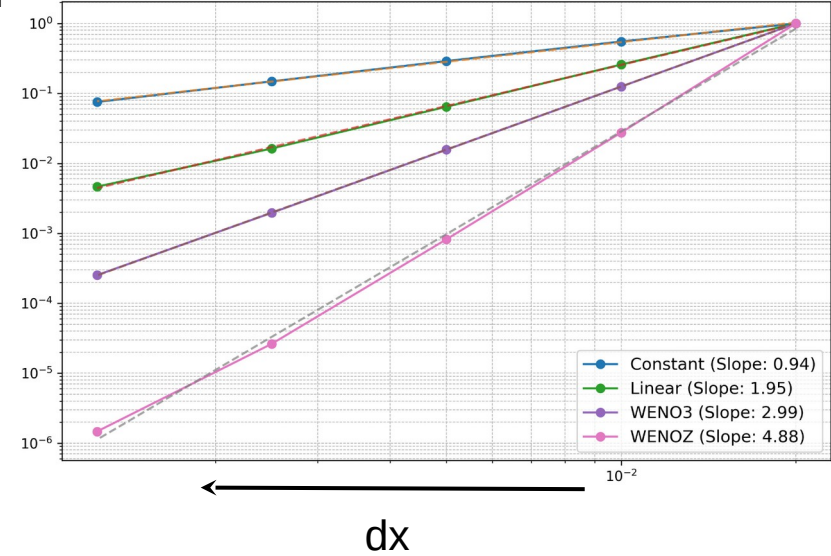
Theoretical Validation: Convergence



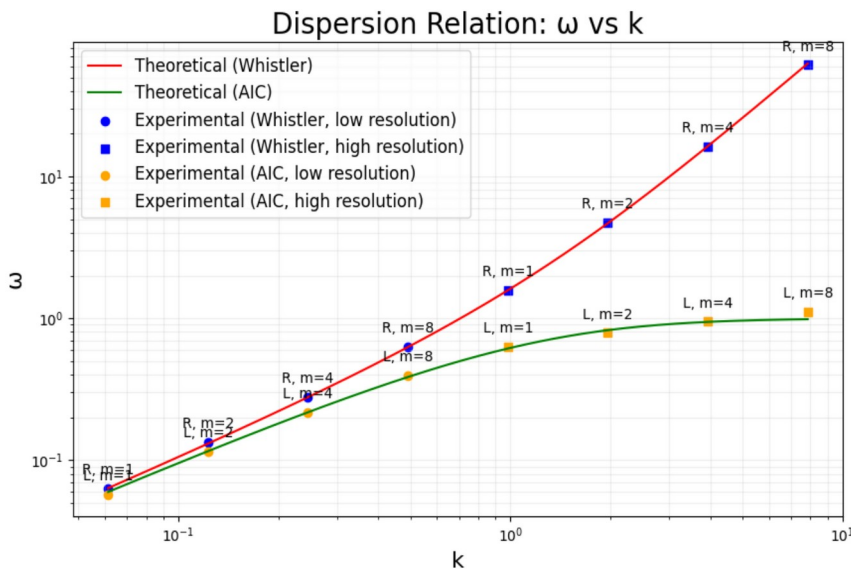
error



Convergence Plot



Theoretical Validation: Hall MHD dispersion



Alfvén Ion Cyclotron:

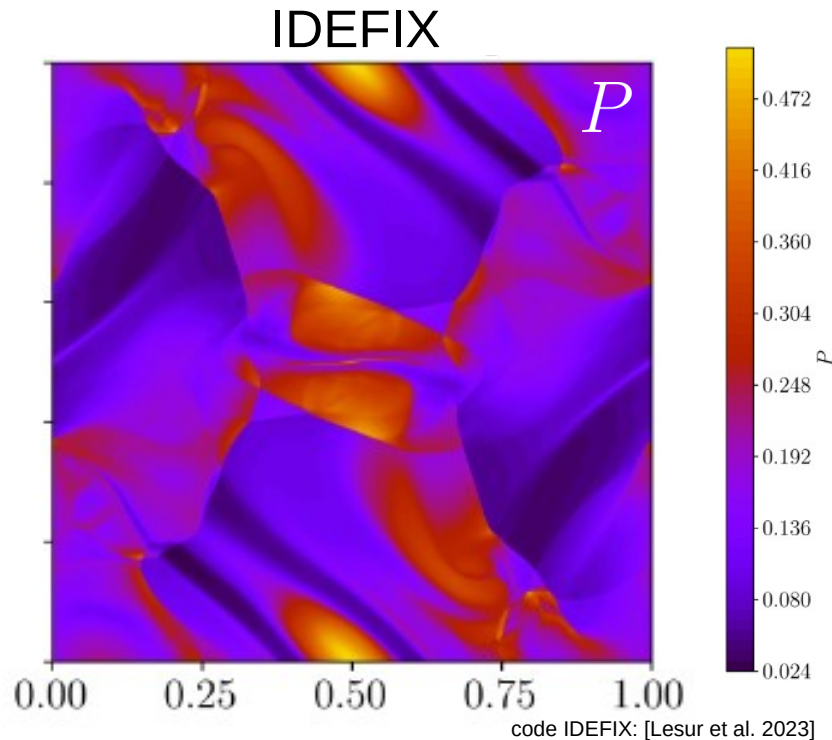
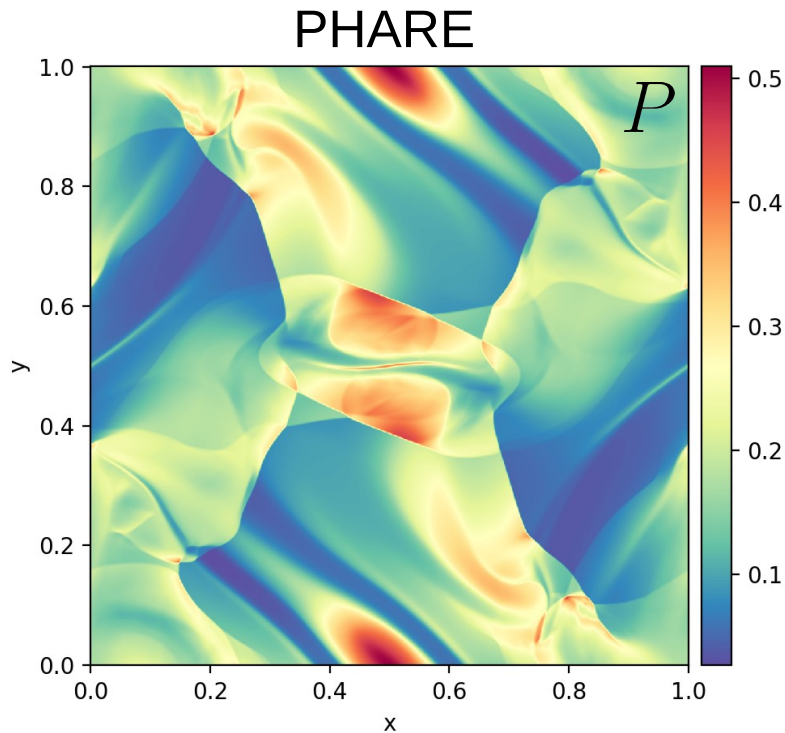
$$\omega_L = \frac{k^2}{2} \left(\sqrt{1 + \frac{4}{k^2}} - 1 \right)$$

Whistler:

$$\omega_R = \frac{k^2}{2} \left(\sqrt{1 + \frac{4}{k^2}} + 1 \right)$$

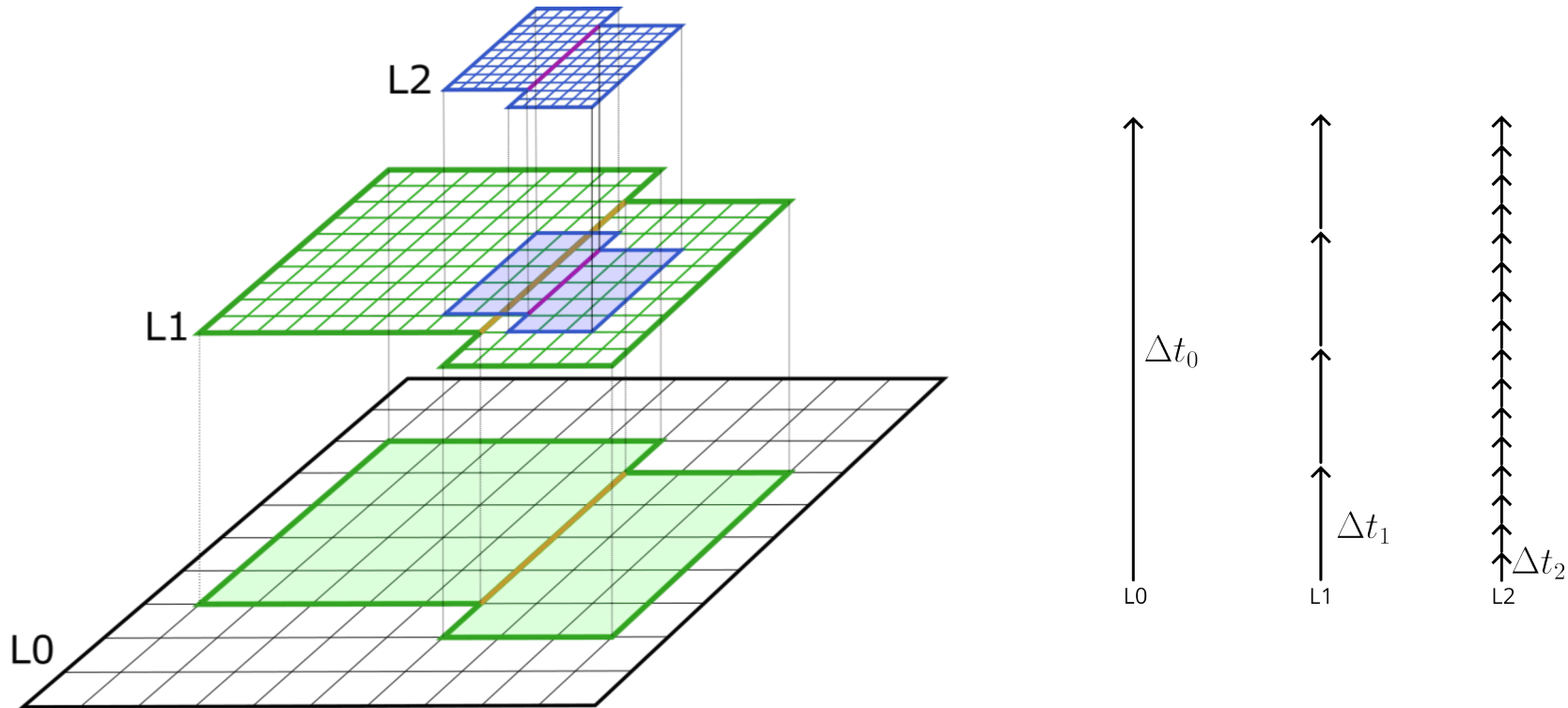
Functional Test: Orszag-Tang vortex

[Orszag & Tang 1979]

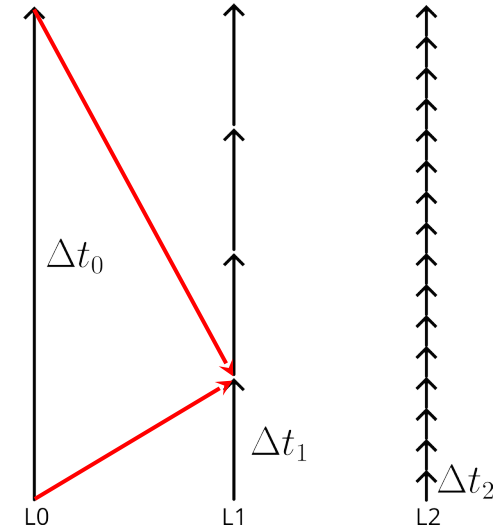
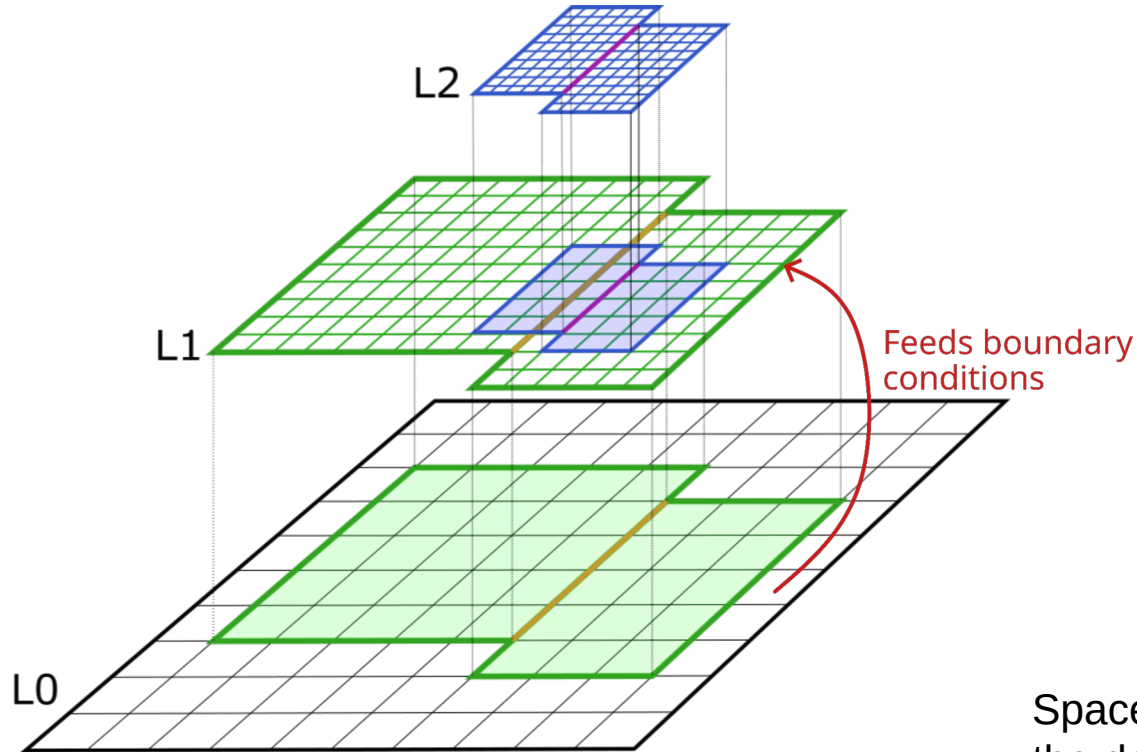


Pressure plot of the Orszag-Tang vortex at $t=0.5$ with a domain of 1, 1024x1024 cells and a second order scheme. PHARE (left) Idefix (right).

Step 3: Adaptive Mesh Refinement

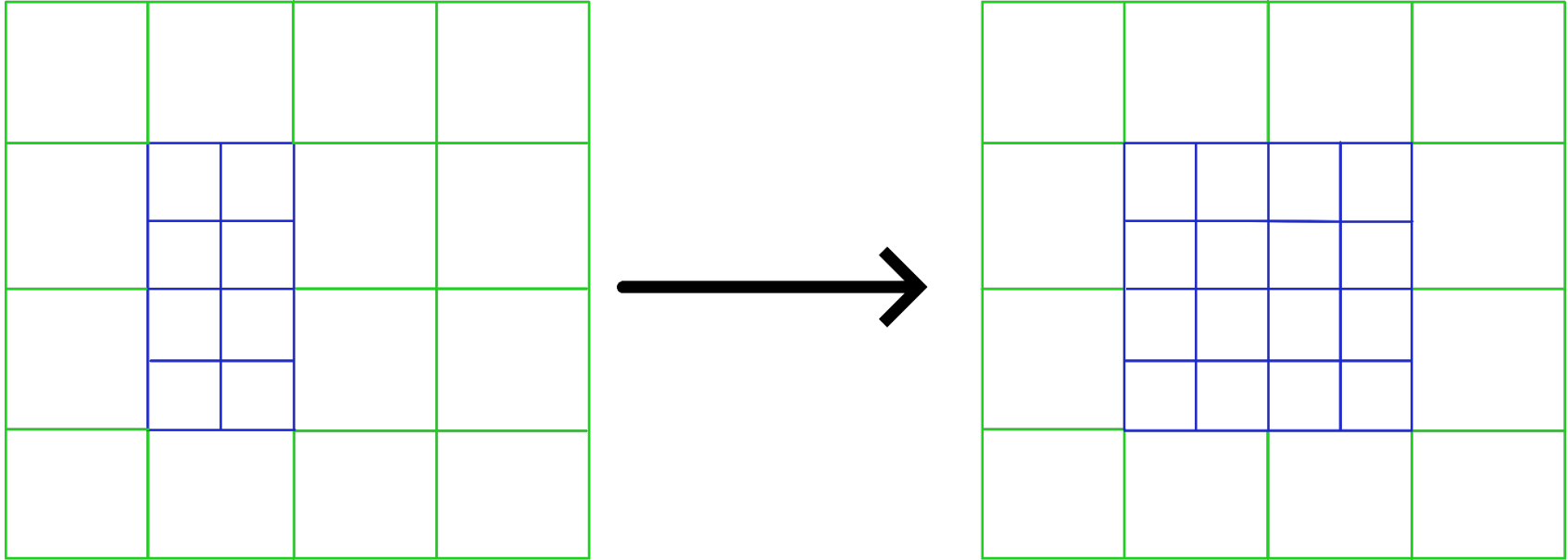


Step 3: Adaptive Mesh Refinement



Space and time interpolation of the data at boundaries.

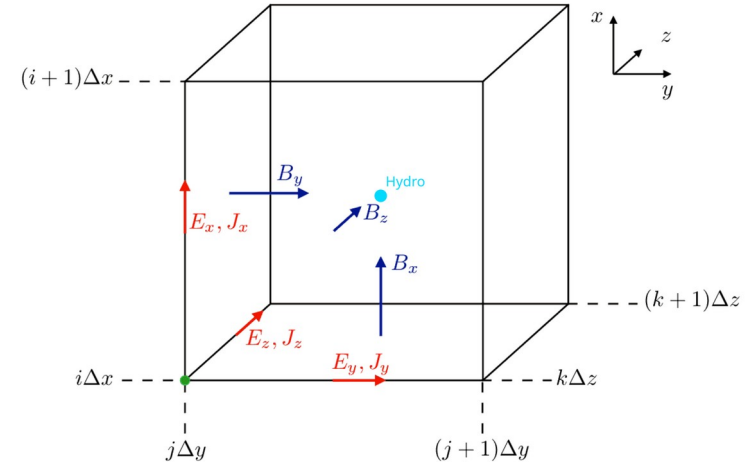
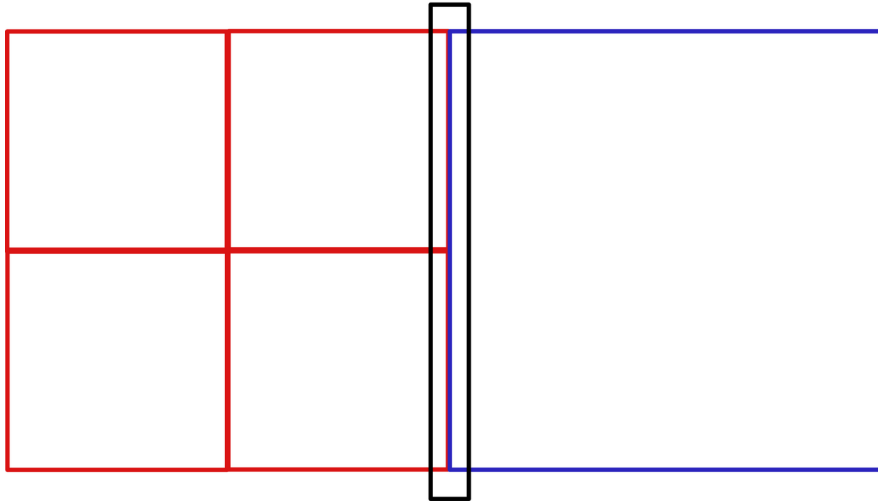
Step 3: Adaptive Mesh Refinement



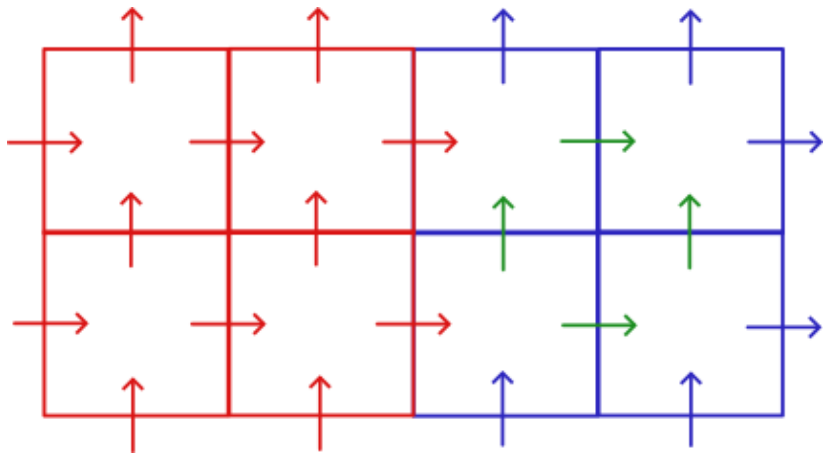
Adaptive mesh: grid can change to follow a criteria on the solution

The conservative refinement problem

coarse-fine boundary



The conservative refinement problem



[Toth and Roe (2002)]

Idea: refine the new fine faces in such a way that:

$$\nabla \cdot \mathbf{B}_{fine} = \frac{1}{r^d} \nabla \cdot \mathbf{B}_{coarse}$$

with r the refinement ratio and d the dimension.

$$u = B_x, v = B_y, w = B_z$$

$$u^{0,j,k} = \frac{1}{2}(u^{+,j,k} + u^{-,j,k}) + U_{xx} + k(\Delta z)^2 V_{xyz} + j(\Delta y)^2 W_{xyz}$$

$$v^{i,0,k} = \frac{1}{2}(v^{i,+,k} + v^{i,-,k}) + V_{yy} + i(\Delta x)^2 W_{xyz} + k(\Delta z)^2 U_{xyz}$$

$$w^{i,j,0} = \frac{1}{2}(v^{i,j,+} + v^{i,j,-}) + W_{zz} + j(\Delta y)^2 U_{xyz} + i(\Delta x)^2 V_{xyz}$$

with: $U_{xx} = \frac{1}{8} \sum_{ijk=\pm} ijv^{i,j,k} + ikw^{i,j,k}$

$$V_{yy} = \frac{1}{8} \sum_{ijk=\pm} iju^{i,j,k} + jkw^{i,j,k}$$

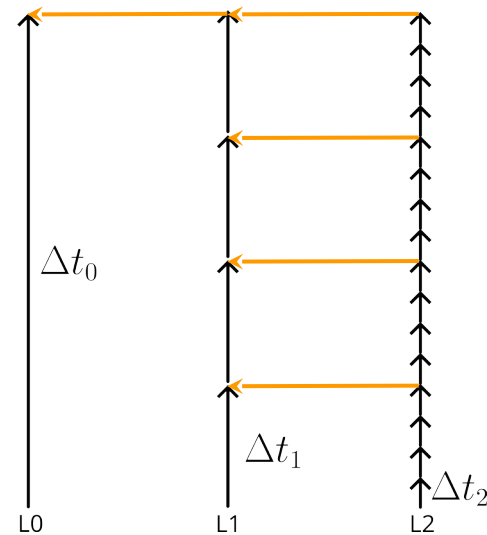
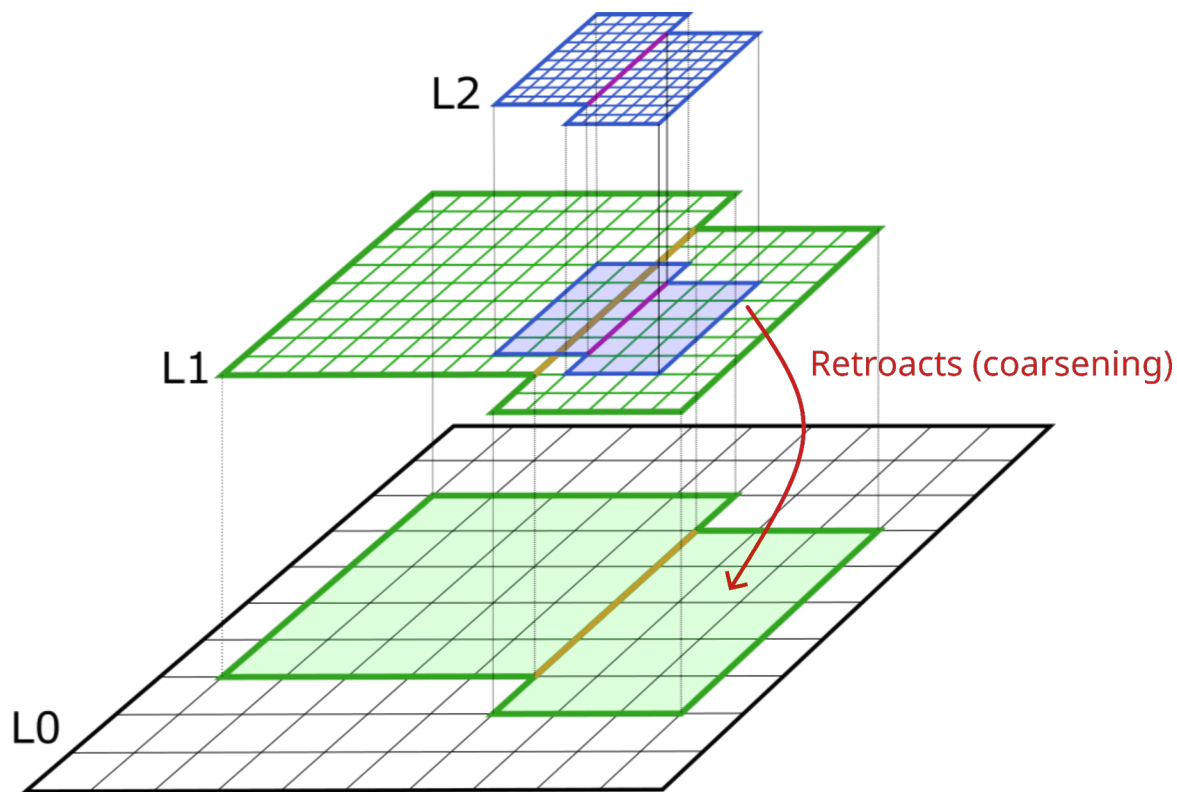
$$W_{zz} = \frac{1}{8} \sum_{ijk=\pm} iku^{i,j,k} + jkv^{i,j,k}$$

and: $U_{xyz} = \frac{1}{8} \sum_{ijk=\pm} \frac{ijk u^{i,j,k}}{(\Delta y)^2 + (\Delta z)^2}$

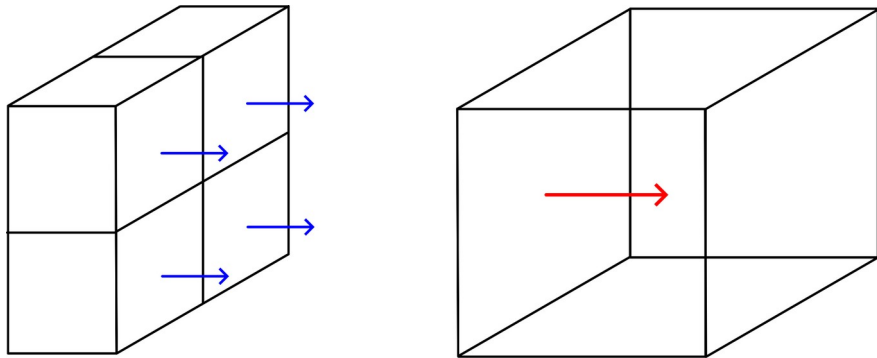
$$V_{xyz} = \frac{1}{8} \sum_{ijk=\pm} \frac{ijk v^{i,j,k}}{(\Delta x)^2 + (\Delta z)^2}$$

$$W_{xyz} = \frac{1}{8} \sum_{ijk=\pm} \frac{ijk w^{i,j,k}}{(\Delta x)^2 + (\Delta y)^2}$$

Step 3: Adaptive Mesh Refinement



Conserving flux at coarse fine interfaces



Need to express the whole time integration step with a single flux. With multistep integrators, use Butcher flux:

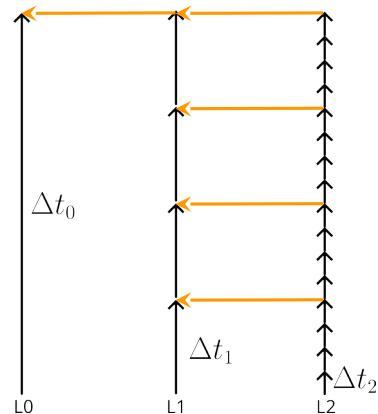
$$\mathbf{U}^{n+1} = \mathbf{U}^n + \sum_{step} \gamma(\omega_{step}, \omega_{step-1}, \dots, \omega_0) \mathbf{F}^{step}$$

$$\mathbf{F}_{fine}^{total} = \frac{1}{c_{sub-steps} \times c_{shared \text{ fine faces}}} \sum_{sub-steps} \sum_{shared \text{ fine faces}} \mathbf{F}_{fine}$$

← Space-time average of the fluxes

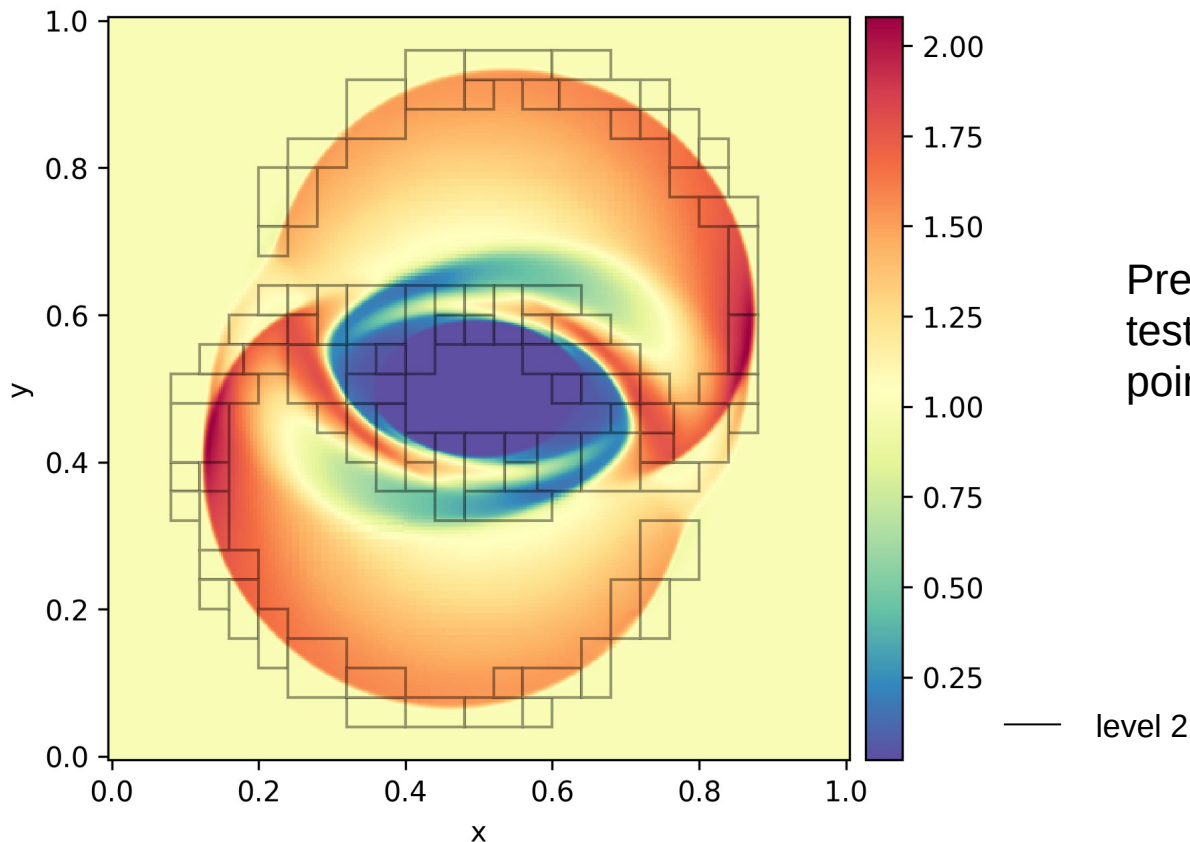
$$\mathbf{U}^{n+1} = \mathbf{U}^{n+1} - \frac{\Delta t_{coarse}}{\Delta l} \left(\mathbf{F}_{coarse} - \mathbf{F}_{fine}^{total} \right)$$

← Correction step



AMR Validation: MHD Rotor test

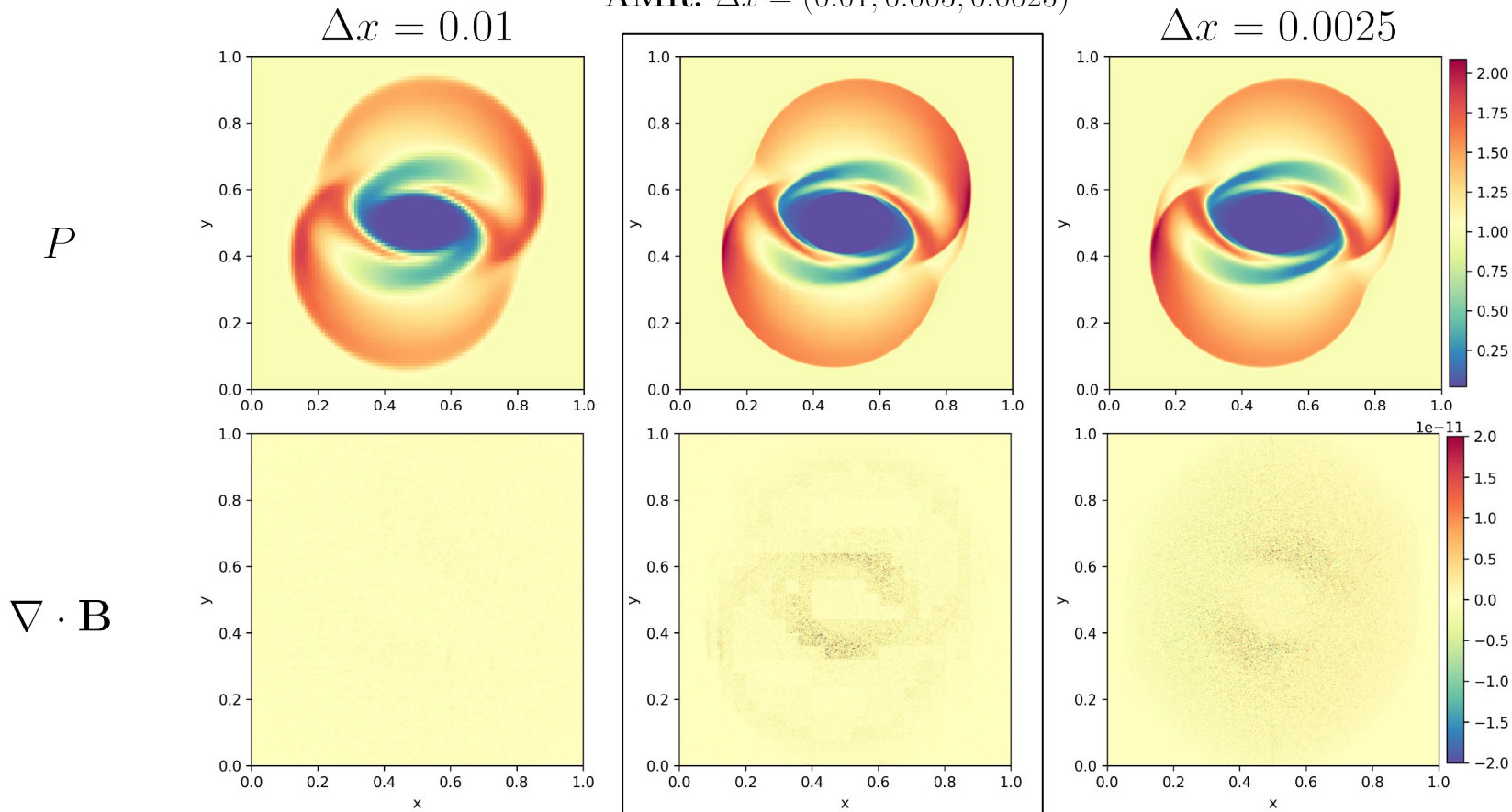
[Balsara & Spicer
1999]



Pressure for the MHD rotor
test at $t=0.15$. 100x100 grid
points, 3 AMR level.

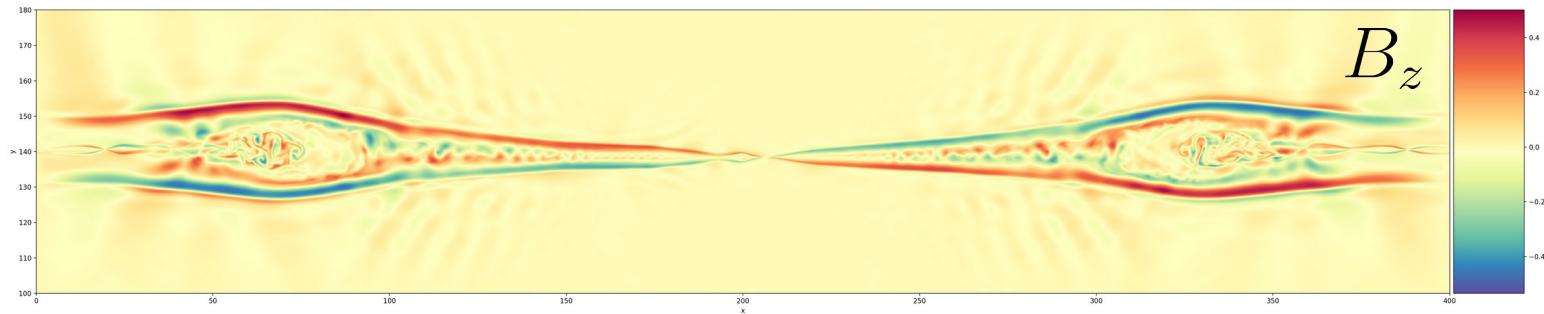
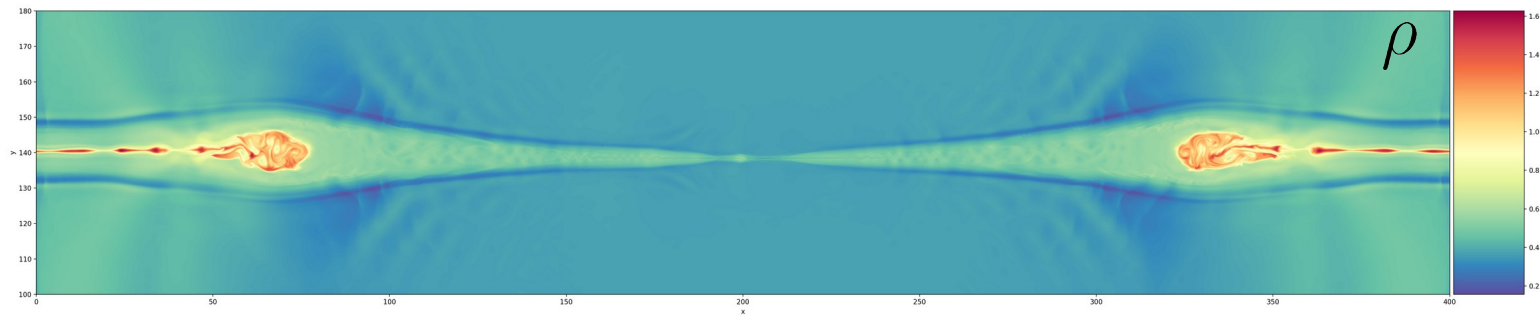
AMR Validation: MHD Rotor test

AMR: $\Delta x = (0.01, 0.005, 0.0025)$

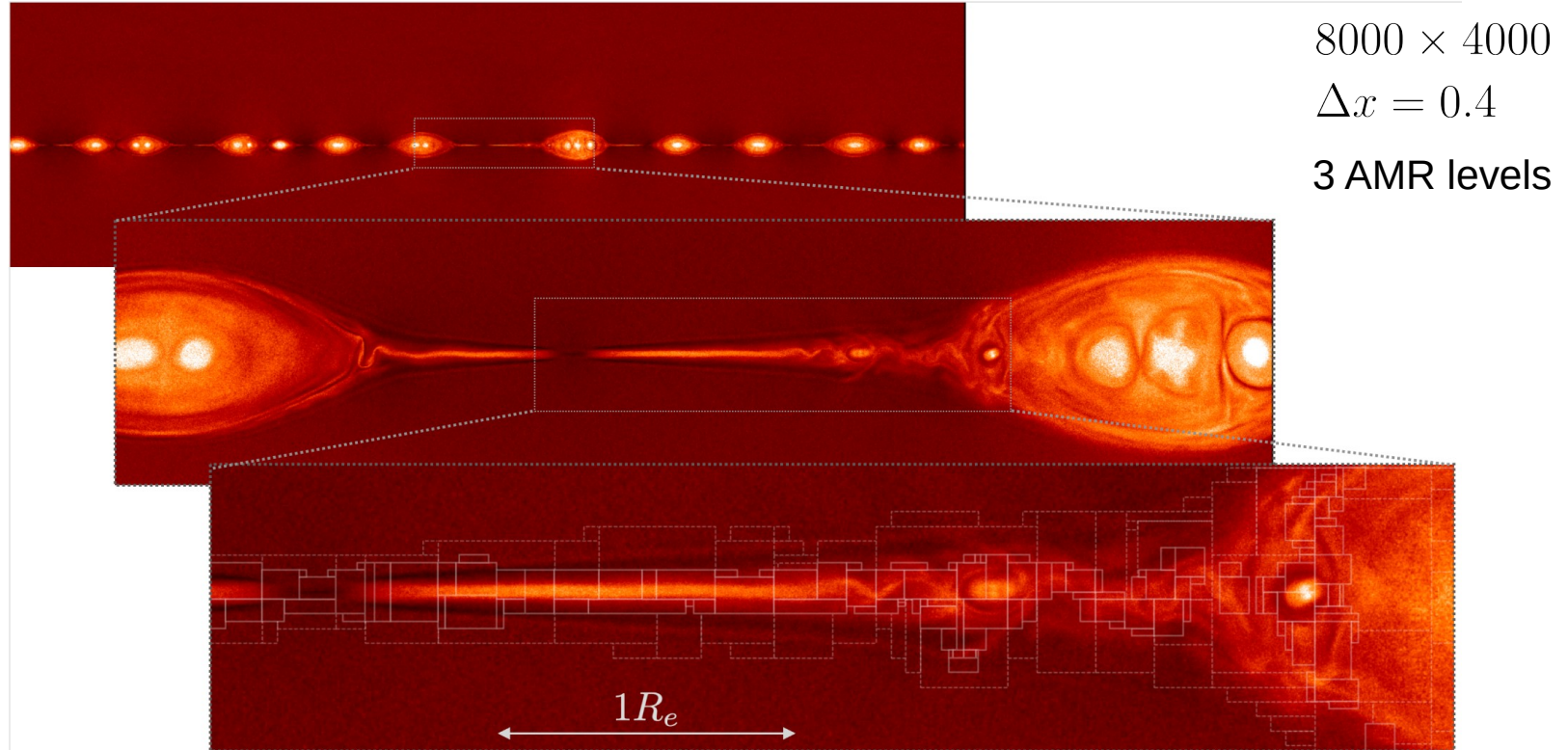


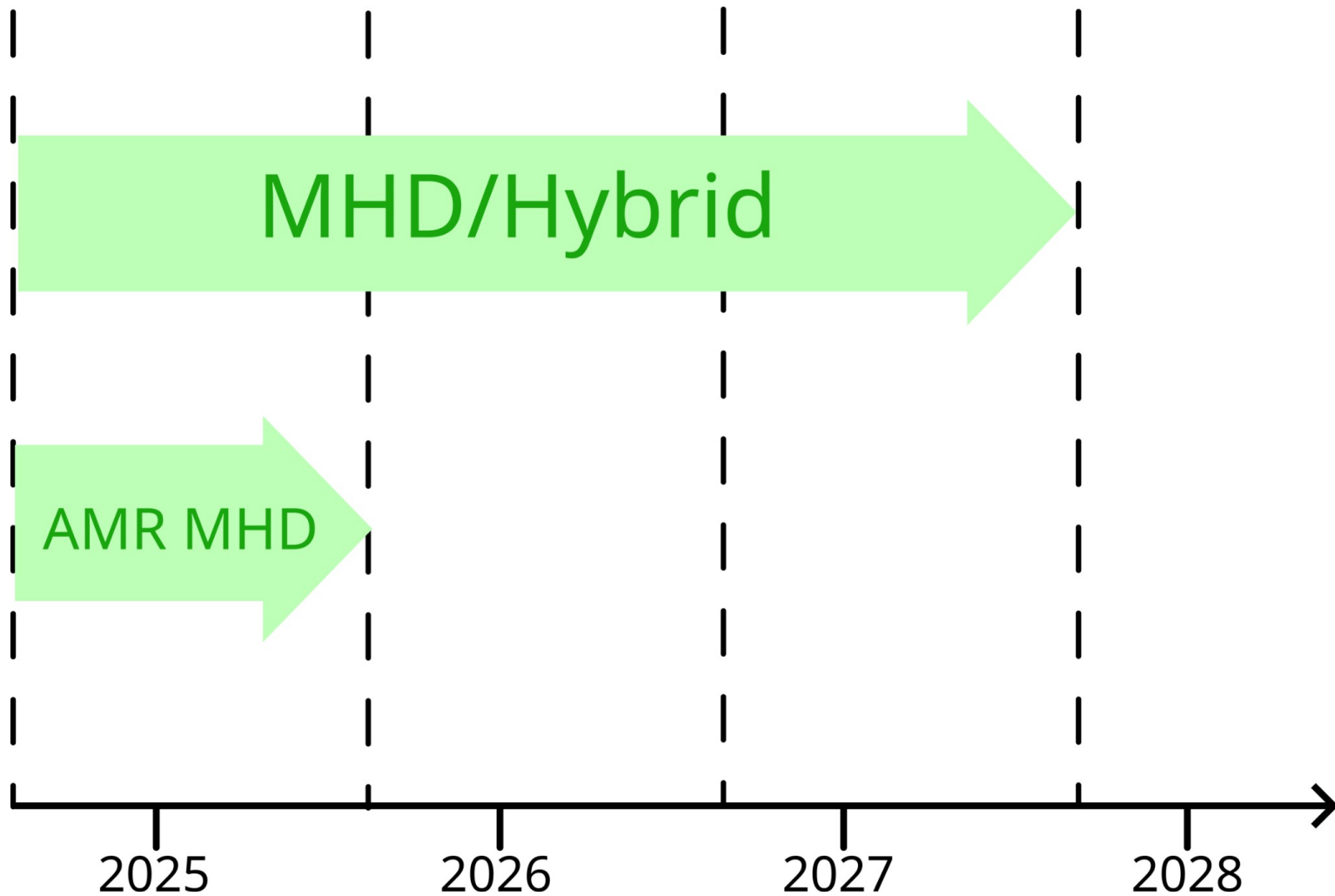
1000×500
 $\Delta x = 0.4$
3 AMR levels

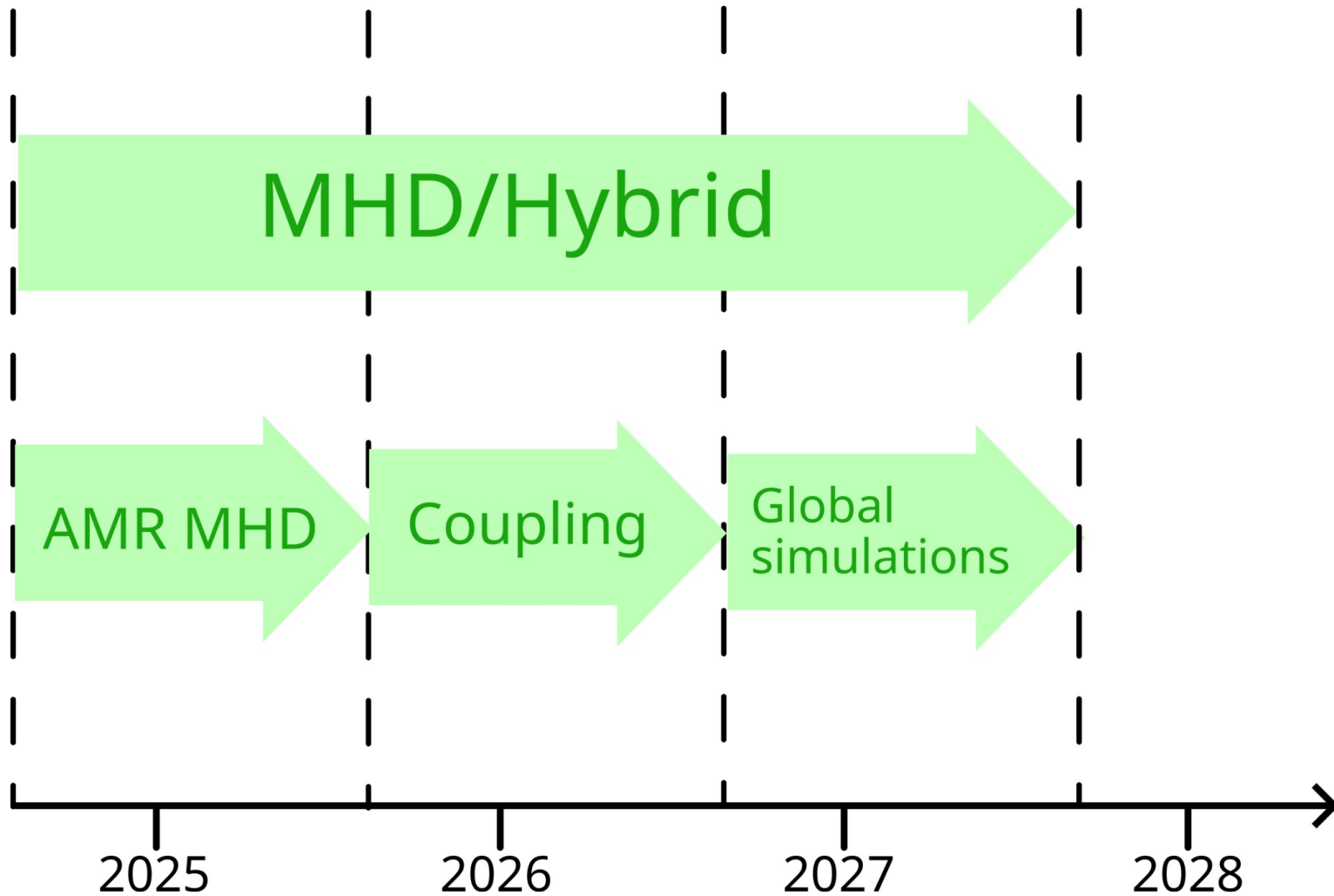
Symmetric current sheet [Harris 1962]



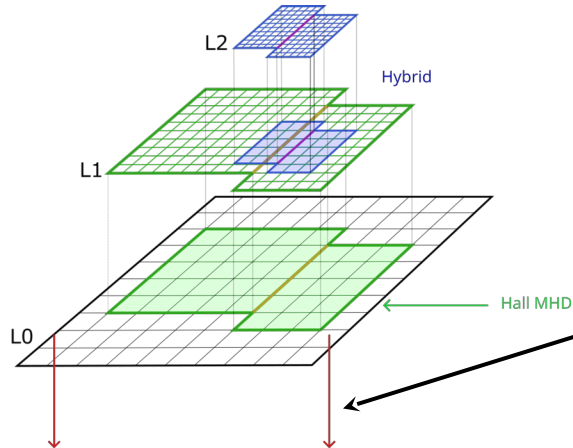
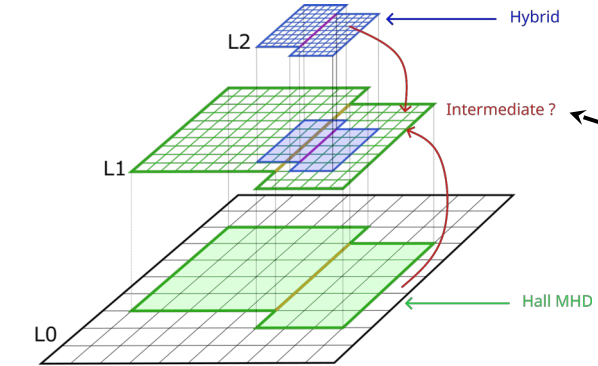
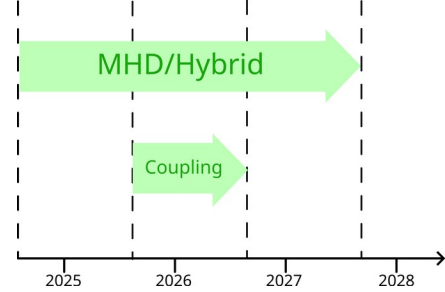
Symmetric current sheet with the hybrid solver







Lots to experiment with



- Higher moment schemes ?
- Test particles in MHD ?
- Guiding center schemes ?
- Ideal MHD ?
- With asymptotic preserving schemes ?

Many new things coming to the code in the coming year

- Magnetospheric boundary conditions
- 3d visualisation
- Improved performances
- MHD/Hybrid coupling
 - The main ingredients for our global simulations !

