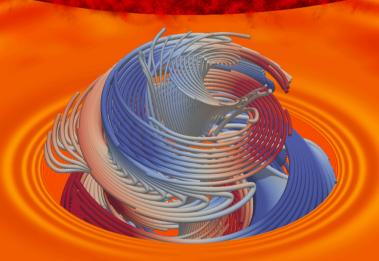
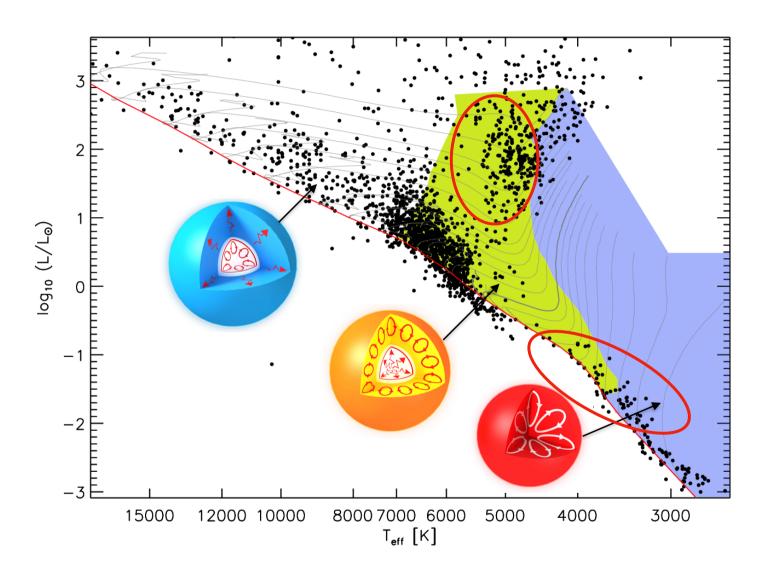
Modelling some aspects of stellar magnetism with MagIC

Laurène Jouve IRAP Toulouse, France



Grenoble, December 2025

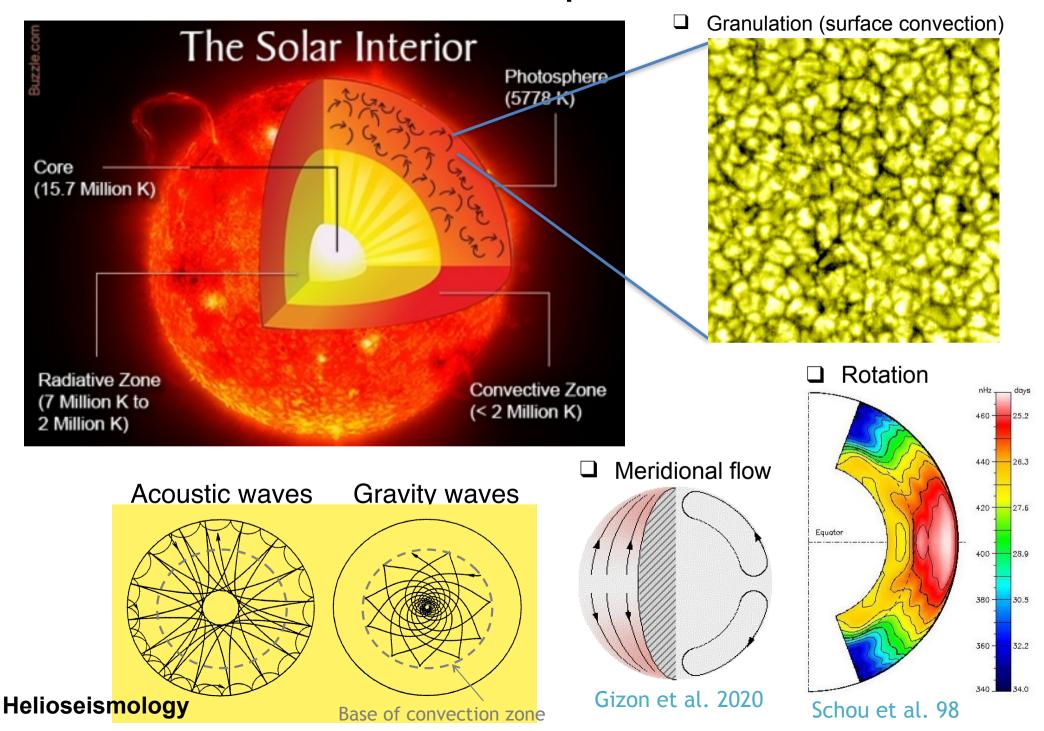
Internal structure of stars



Focus in this presentation on:

- √ Convective envelopes of cool stars
- ✓ Radiative cores of red giant stars

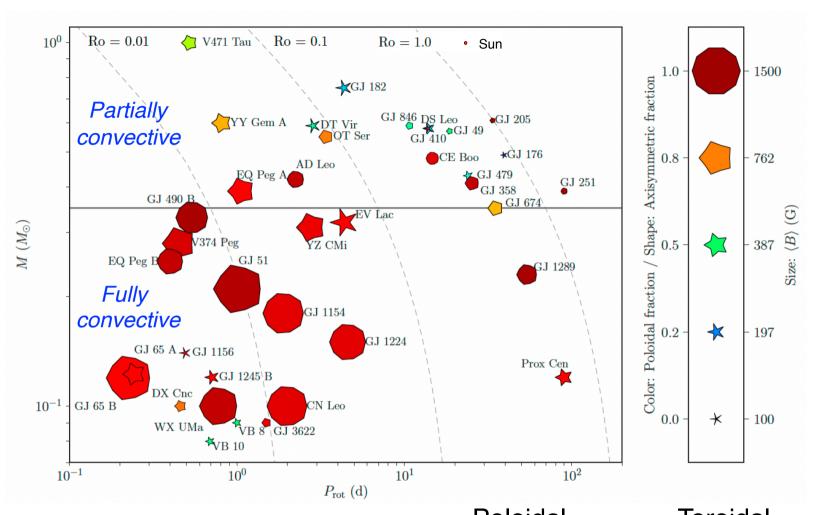
Solar interior and plasma flows



Magnetic fields in cool stars (M dwarfs)

Morin, Donati et al. (2008-2011)
Moutou et al. (2017)
Kochukhov &
Shulyak (2019)
Klein et al. (2021)
Kochukhov (2021)
Zaire PhD thesis (2022)

Rossby number $Ro = P_{rot}/\tau_c$



- ☐ Mostly multipolar $M_{\odot} > 0.35$ for and/or Ro < 0.1
- Mostly dipolar for M_☉ < 0.35 and/or Ro > 0.1

Poloidal Toroidal

What controls the magnetic topology of cool stars?

Anelastic MHD simulations

Anelastic equations: filtering out sound waves to avoid untractably small time steps: very efficient and accurate for stellar interiors!

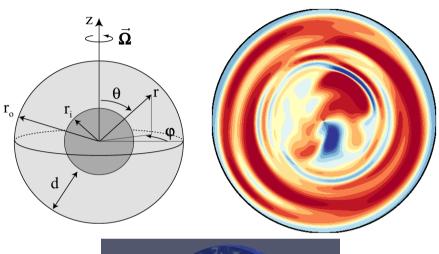
$$\begin{split} \left(\frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \vec{\nabla} \vec{u}\right) &= -\vec{\nabla} \left(\frac{p'}{\tilde{\rho}}\right) - \frac{2}{E} \vec{e_z} \times \vec{u} + \frac{Ra}{Pr} \tilde{g} \, s' \, \vec{e_r} + \frac{Ra_\xi}{Sc} \, \tilde{g} \, \xi' \, \vec{e_r} + \frac{1}{Pm \, E \, \tilde{\rho}} \left(\vec{\nabla} \times \vec{B}\right) \times \vec{B} + \frac{1}{\tilde{\rho}} \, \vec{\nabla} \cdot \mathbf{S}, \\ \vec{\nabla} \cdot \tilde{\rho} \vec{u} &= 0, \\ \vec{\rho} \left(\frac{\partial \xi'}{\partial t} + \vec{u} \cdot \vec{\nabla} \xi'\right) &= \frac{1}{Sc} \, \vec{\nabla} \cdot \left(\kappa_\xi(r) \tilde{\rho} \vec{\nabla} \xi'\right) \\ \frac{\partial \vec{B}}{\partial t} &= \vec{\nabla} \times \left(\vec{u} \times \vec{B}\right) - \frac{1}{Pm} \, \vec{\nabla} \times \left(\lambda(r) \, \vec{\nabla} \times \vec{B}\right) \\ \tilde{\rho} \tilde{T} \left(\frac{\partial s'}{\partial t} + \vec{u} \cdot \vec{\nabla} s'\right) &= \frac{1}{Pr} \, \vec{\nabla} \cdot \left(\kappa(r) \tilde{\rho} \tilde{T} \vec{\nabla} s'\right) + \frac{Pr \, Di}{Ra} \Phi_{\nu} + \frac{Pr \, Di}{Pm^2 \, E \, Ra} \lambda(r) \left(\vec{\nabla} \times \vec{B}\right)^2, \end{split}$$

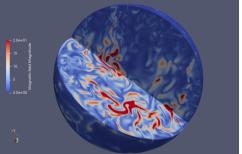
Dimensionless control parameters

$$E=rac{
u}{\Omega d^2}, \hspace{1cm} Ra=rac{lpha_o g_o T_o d^3 \Delta s}{c_p \kappa_o
u_o}, \hspace{1cm} Pr=rac{
u_o}{\kappa_o}, \hspace{1cm} Pm=rac{
u_o}{\lambda_i}.$$
 Viscosity/Coriolis Buoyancy/Dissipation Viscosity/Thermal diff Viscosity/Mag diff

The MagIC code

Geometry and numerical method



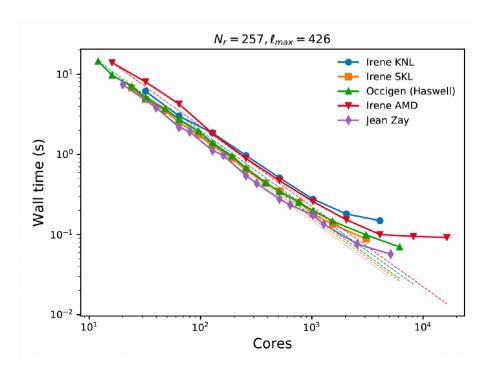


- □ Parallelisation and scaling
 - Hybrid MPI/OpenMP parallelisation
 - Runs on regional and national centers
 - Scales up to several thousands of CPUs
 - No GPU version yet



Wicht 2002, Gastine & Wicht 2012

- Boussinesq/anelastic MHD equations
- Pseudo-spectral code: spherical harmonics in θ, ϕ (use of SHTns) Chebyshev or FD in r
- Several IMEX time stepping methods (here CNAB2 method)
- Full sphere setup now
- Freely available on Github: https://magic-sph.github.io/



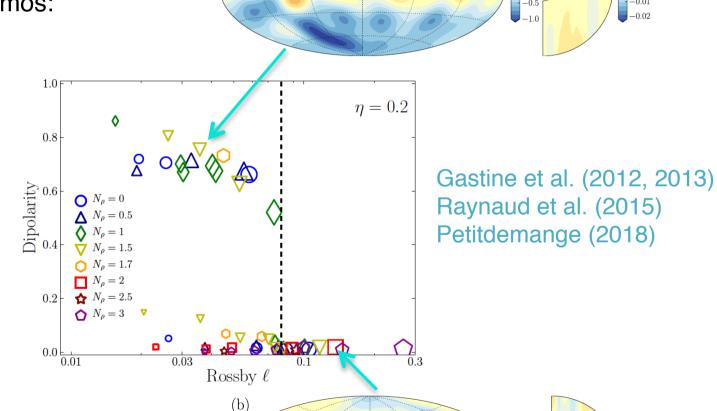
Magnetic topology in cool stars: influence of the Rossby number

(a)

☐ Change in Rossby
Ro=inertia/Coriolis
(also seen in planetary dynamos:

Christensen & Aubert 2006)

- Small Ro: Ordering role of Coriolis=dipolar (no role of shear)
- Large Ro:Inertia becomesdominant=multipolar(important role of shear)



0.5

0.4

0.01

0.01

Strong stratification leads to multipolar fields ? $N_{\rho} = \ln(\rho_i/\rho_o)$

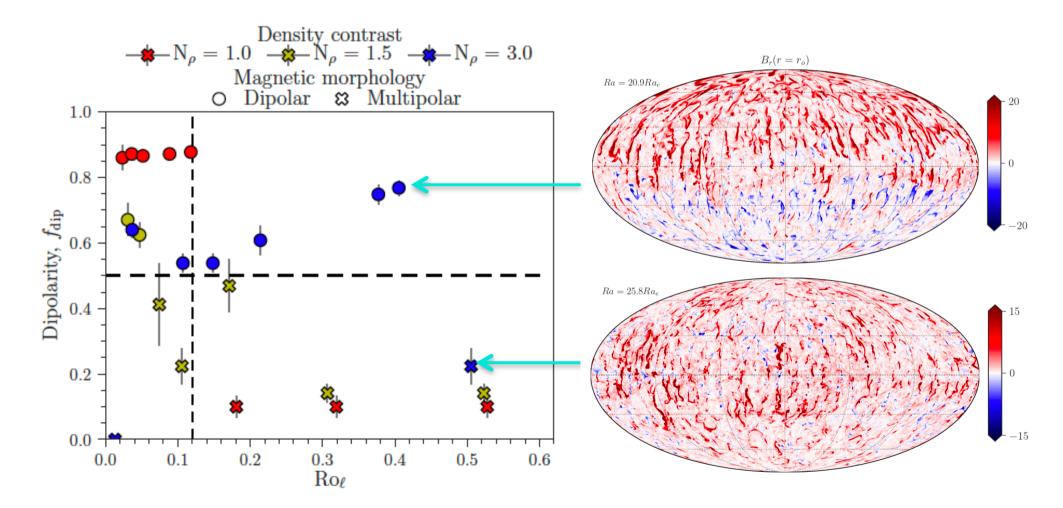
Yadav et al. 2015, 2016

Magnetic topology: influence of the Rossby number?

Zaire, Jouve et al. 2022

MagIC Code (pseudo-spectral 3D MHD) Wicht 2002, Gastine & Wicht 2012

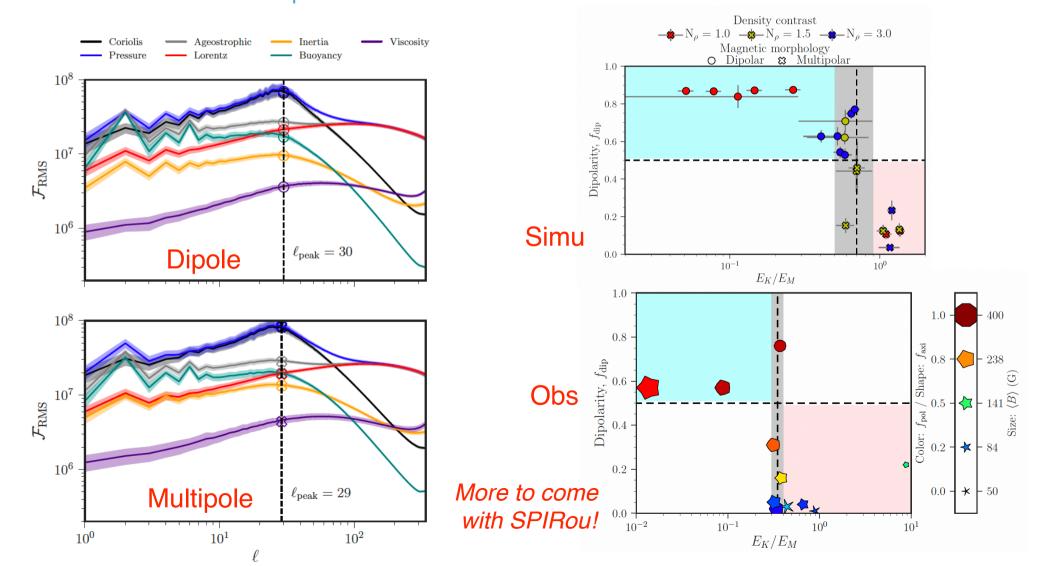
☐ With more turbulent simulations with stronger role of Lorentz force on dynamics (Pm=5), dipole seems to survive at high Ro and high Nrho.



Magnetic topology: influence of the Rossby number?

- Need to look at new force balance (Schwaiger et al. 2019)
- ☐ The ratio of inertia to Lorentz forces (instead of inertia to Coriolis) seems to be a better indicator of dipolarity => ratio of kinetic to magnetic energies as a proxy

First seen in Boussinesq calculations of Menu et al. 2020 and then Tassin et al. 2021

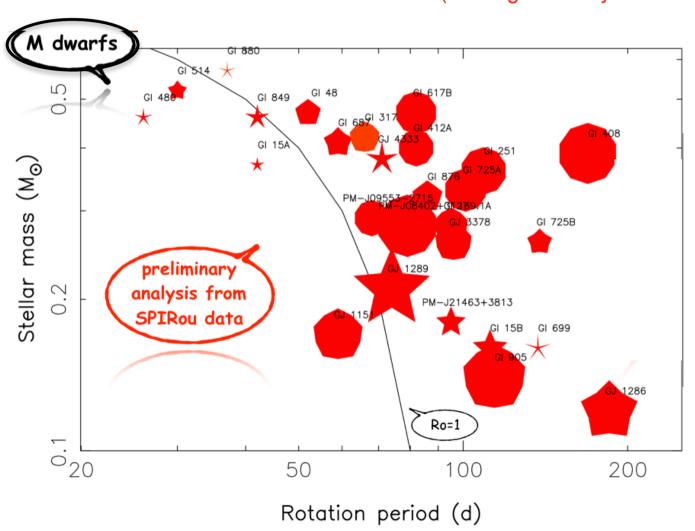


Magnetic topology: influence of the Rossby number?

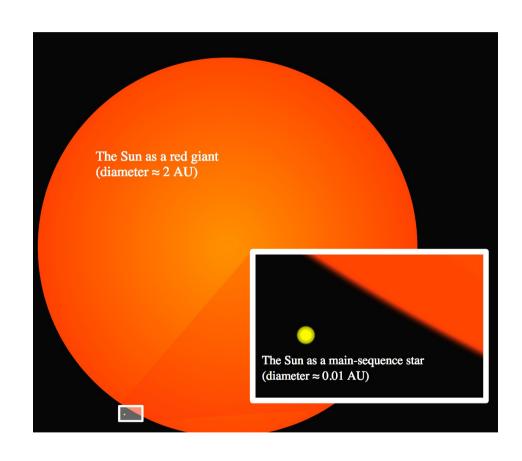
New SPIRou results

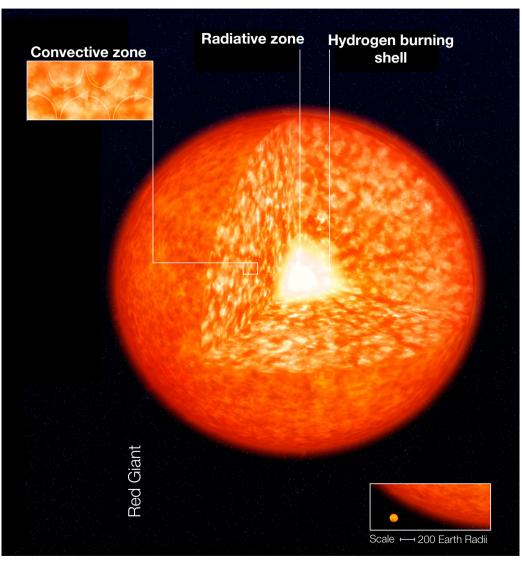
Donati et al. (not published yet) Lehmann et al. 2024, See et al. 2025

- 50 weakly active M dwarfs monitored for 7 semesters (2019-2022)
- Some polarity reversals (cycles?)
- Large-scale fields even for slow rotators (i.e. large Rossby numbers)



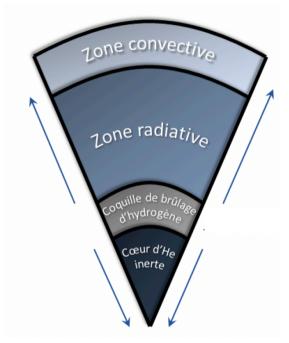
Red giant stars

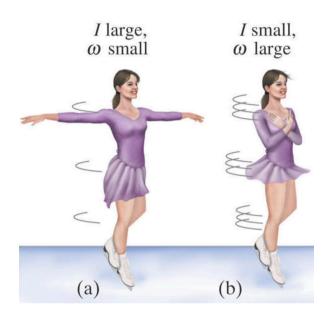


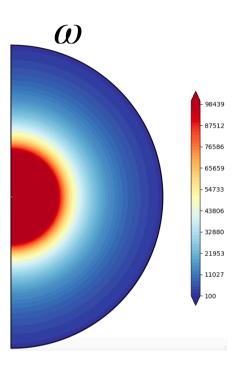


What do we know about them: their core rotates much slower than predicted

Core contraction => spin-up







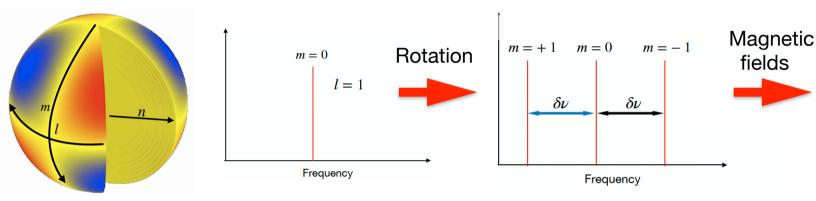
- But strong constrains by asteroseismology:
 - For young sub-giants: evidence for quasi-solid body rotation (Deheuvels et al. 2020)
 - For more evolved sub-giants: Core spin-up and envelope spin-down with $\Omega_c \approx 5-15~\Omega_{env}$ (e.g. Deheuvels et al. 2014)
 - For red giants: Ω_c remains constant despite contraction (Gehan et al. 2018)

=> An efficient process at play to redistribute angular momentum at least during young subgiant and red giant phases: without magnetic fields, all models have failed so far

What do we know about them: strong magnetic fields are present in their cores!

☐ Recent discovery of fields of 100kG and more in the core of red giants

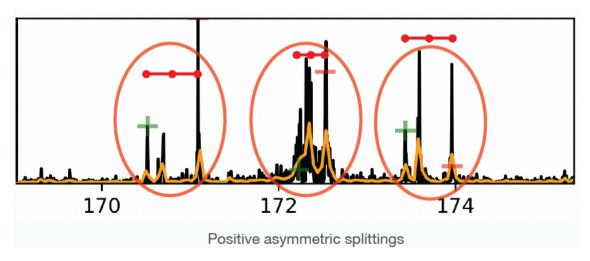
ANR Beaming (PI S. Deheuvels): Li et al. 2022, Nature, Deheuvels et al. 2023, Li et al. 2023



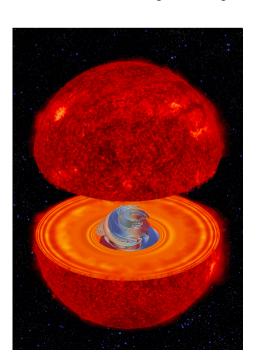
m = +1 m = 0 m = -1Frequency

Symmetric multiplet Positive or negative asymmetry

Observed magnetic splitting

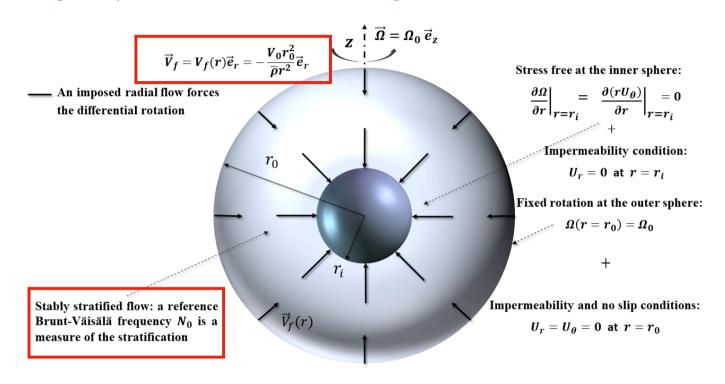


Can magnetic fields help at solving the AM problem in red giants?



2D Boussinesq/anelastic MHD simulations of a contracting RZ

☐ Fluid (+ mag field) contained in a contracting zone with stable stratification.

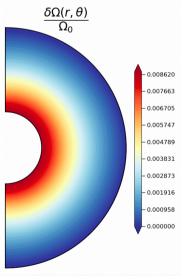


Gouhier et al. 2022, 2023

- $lue{}$ Without magnetic fields, a radial differential rotation is established when N_0 is large (as expected)
- With an initial magnetic field?

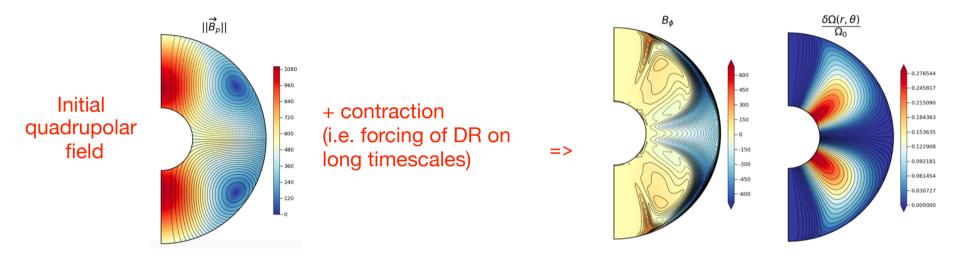
Alfvén time-scale:
$$au_{Ap} = \frac{\sqrt{4\pi\rho}\,r_o}{B_p}$$

In stars (and in the simulations), the Alfvén time-scale is small compared to the contraction time-scale $\tau_c=r_o/V_0$

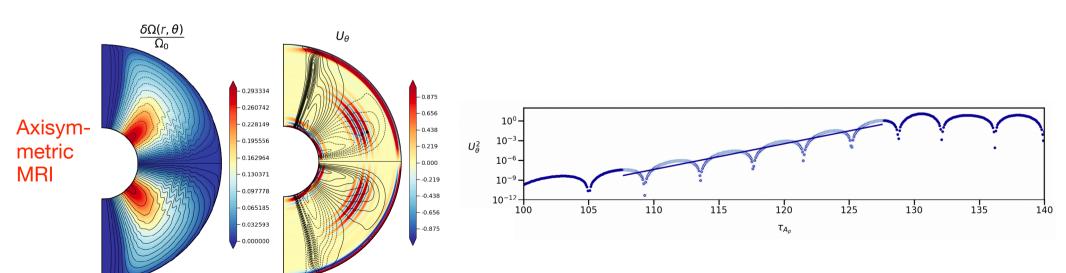


Magnetic tension + magnetorotational instability

Initial poloidal field, production of toroidal field, shaping of differential rotation by magnetic tension



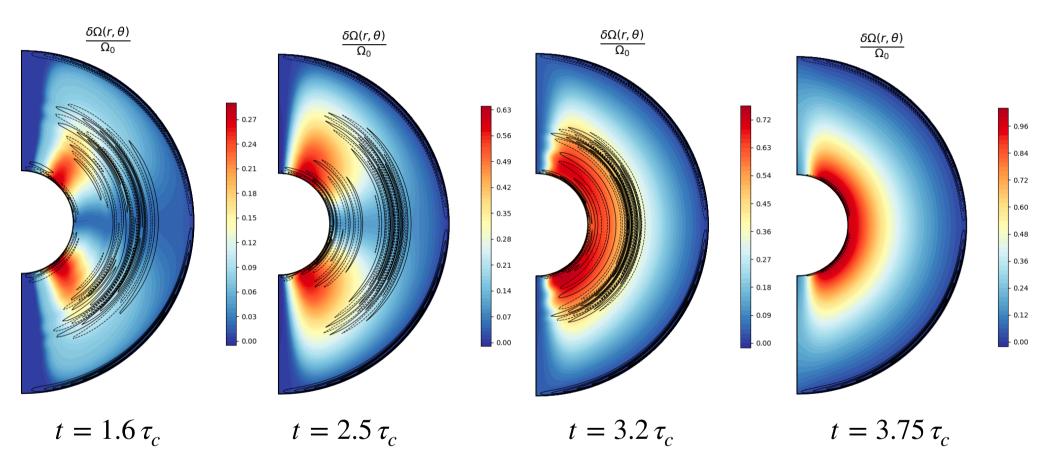
 $lue{}$ An instability sets in if field not too strong: in strongly stratified cases, MRI may be triggered on latitudinal gradients of Ω , with elongated structures in horizontal directions



Development of the instability

■ Non-linear development of the instability:

Meridional flow with strong v_{θ} + dissipation drastically reconfigure flow and field



- ☐ Scenario for AM evolution in red giant stars:
 - \Box 1. Large-scale magnetic field imposes a low level of differential rotation during 1 or 2 τ_c => young sub-giants with quasi-solid body rotation
 - □ 2. MRI destroys the large-scale field: differential rotation builds up freely => measured diff. rot on more evolved sub giants
 - □ 3. Another efficient transport spins-down the core of red giants: non-axi instabilities?

Dynamo action based on magneto-rotational instability

 Ω_*

9375

8750 8125

6250

5625

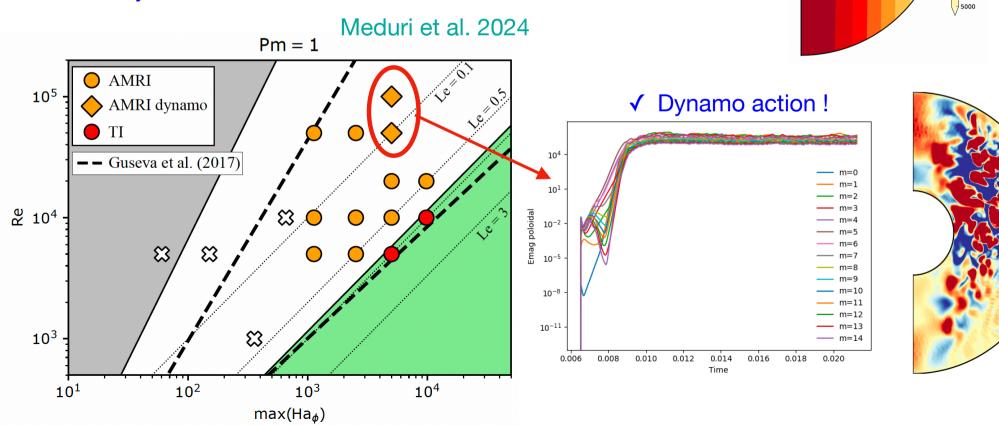
(c)

☐ Local shearing box calculations showed MRI dynamo in accretion disk context

(review by Rincon et al. 2019)

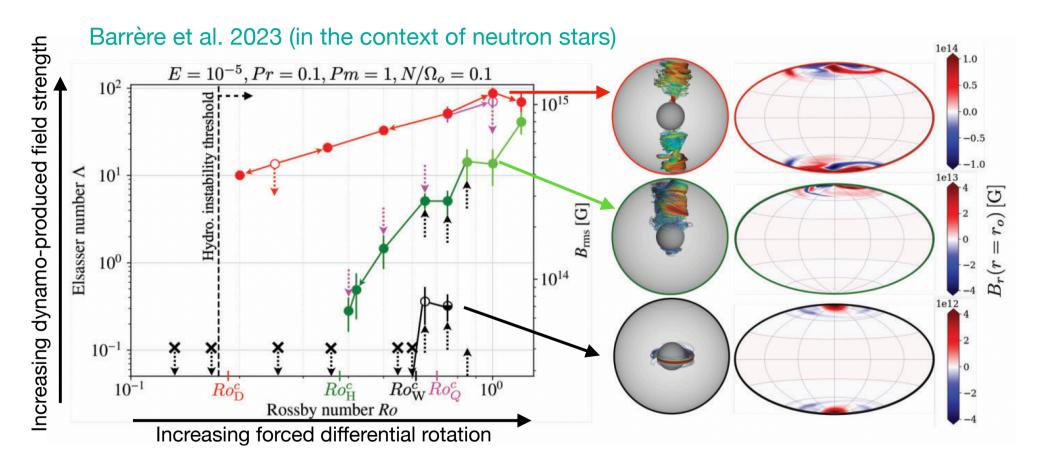
Global geometry: Volumetric forcing of differential rotation (Meduri et al. 2024) or Taylor-Couette flows (Guseva et al. 2017) Reynolds number $Re=\Omega_0 d^2/\nu$

- Purely toroidal field $Ha_{\phi} = B_{\phi}d/\sqrt{4\pi\rho\nu\eta}$
- A variety of behaviours



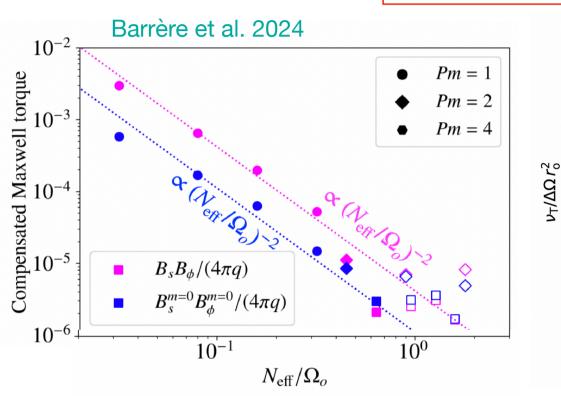
Dynamo action based on Tayler instability

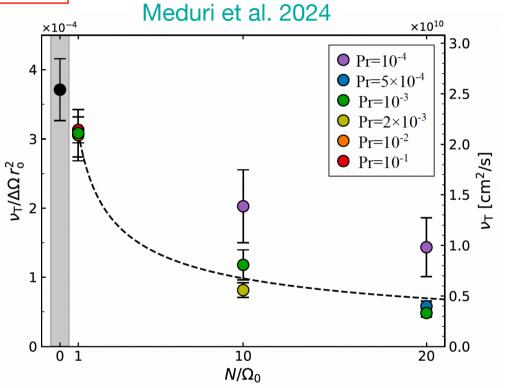
Previous debates and calculations on TI dynamos: Spruit 2002, Braithwaite 2006, Zahn et al. 2007 New developments: Petitdemange et al. 2023, Barrère et al. 2023, 2024



- ☐ Subcritical dynamos with 2 branches: hemispherical and dipolar
- ☐ Found by first triggering a kinematic dynamo which brings magnetic energy above the TI threshold
- ☐ Strong stabilising effect of stratification

Angular momentum transport by MRI and TI dynamos (effect of stable stratification)





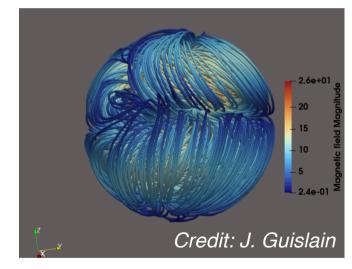
- ☐ Maxwell stresses dominate the transport in both cases
- \square Scaling with stratification different (possibly $N/\Omega^{-1/2}$ for MRI and N/Ω^{-2} for TI): possible comparisons to theoretical predictions (Spruit 1999, 2002, Fuller et al. 2019)
- ☐ Application to angular momentum (and chemical mixing) in stars?

Conclusions/open questions

- Dynamo models of fully convective stars:
 - Change of geometry with Rossby number (or with internal structure?)
 - Studying force balances in detail is necessary
 - How do dipoles resist strong stratifications? Need more realistic parameters and full sphere simulations (Marti et al. 2014, Brown et al. 2020)?

Stellar radiative zones:

- Strong magnetic fields exist in stellar radiative zones (origin in RGS? Thesis of J. Guislain)
- MRI unstable fields if latitudinal differential rotation
- Possibly strong implications for angular momentum



- Radiative zone dynamo? (talks by A. Reboul-Salze, P. Barrère)
- Angular momentum transport by non-axisymmetric instabilities? Local simulations with SNOOPY (Jouve et al. in prep)
- ☐ More to come with SPIRou, PLATO, SolO, PSP,... and more and more efficient codes